



Last lecture (6)

- Drift motion
- Ionospheric conductivities
- Geomagnetic field
- Magnetosphere size (standoff distance)

Today's lecture (7)

- Particle motion in the magnetosphere
- Other magnetospheres



Today

Activity	Date	Time	Room	Subject	Literature
L1	29/8	13-15	E52	Course description, Introduction, The Sun 1, Plasma physics 1	CGF Ch 1, 5, (p 110-113)
L2	1/9	15-17	L52	The Sun 2, Plasma physics 2	CGF Ch 5 (p 114-121), 6.3
L3	5/9	13-15	E51	Solar wind, The ionosphere and atmosphere 1, Plasma physics 3	CGF Ch 6.1, 2.1-2.6, 3.1-3.2, 3.5, LL Ch III, Extra material
T1	8/9	15-17	D41	Mini-group work 1	
L4	12/9	13-15	E35	The ionosphere 2, Plasma physics 4	CGF Ch 3.4, 3.7, 3.8
L5	14/9	10-12	V32	The Earth's magnetosphere 1, Plasma physics 5	CGF 4.1-4.3, LL Ch I, II, IV.A
T2	15/9	15-17	E51	Mini-group work 2	
L6	19/9	13-15	M33	The Earth's magnetosphere 2, Other magnetospheres	CGF Ch 4.6-4.9, LL Ch V.
T3	22/9	15-17	E51	Mini-group work 3	
L7	26/9	13-15	E31	Aurora, Measurement methods in space plasmas and data analysis 1	CGF Ch 4.5, 10, LL Ch VI, Extra material
L8	28/9	10-12	L52	Space weather and geomagnetic storms	CGF Ch 4.4, LL Ch IV.B-C, VII.A-C
T4	29/9	15-17	M31	Mini-group work 4	
L9	3/10	13-15	E52	Interstellar and intergalactic plasma, Cosmic radiation,	CGF Ch 7-9
T5	6/10	15-17	E31	Mini-group work 5	
L10	10/10	13-15	E52	Swedish and international space physics research.	
T6	13/10	15-17	E31	Round-up, old exams.	
Written examination	26/10	8-13	F2		

Mini-groupwork 3

a)

$$\frac{\partial n_e}{\partial t} = q - \alpha n_e^2$$

$$\frac{dn_e(t)}{dt} = 0 \Rightarrow$$

$$\alpha = \frac{q}{n_e^2}$$

$$q(z) = a_i I(z) n_n(z) = a_i I_0 e^{-H a_n n_0 e^{-\frac{z}{H}}} n_0 e^{-\frac{z}{H}}$$

$$H = 7.2 \text{ km (O}_2\text{)}$$

$$q = 3.9 \cdot 10^8 \text{ m}^{-3}\text{s}^{-1}$$

$$n_e(120 \text{ km}) = 7 \cdot 10^4 \text{ cm}^{-3} = 7 \cdot 10^{10} \text{ m}^{-3}$$

Thus

$$\alpha = 9.8 \cdot 10^{-14} \text{ m}^3\text{s}^{-1}$$

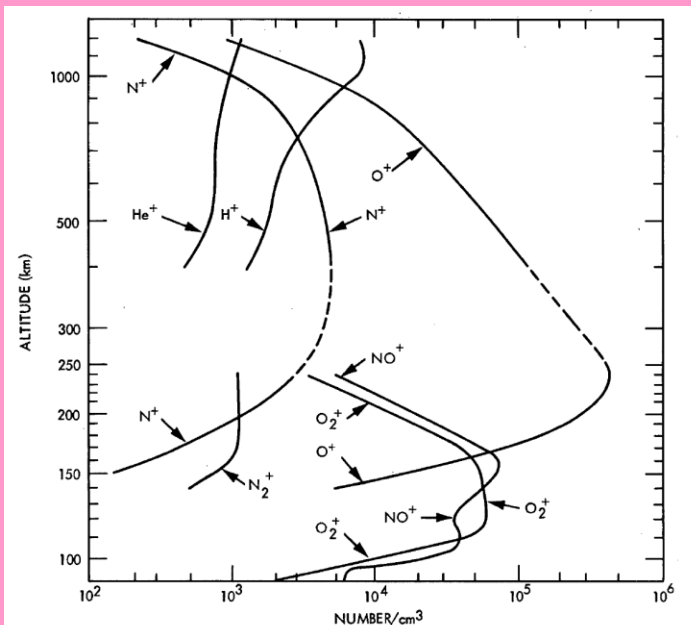
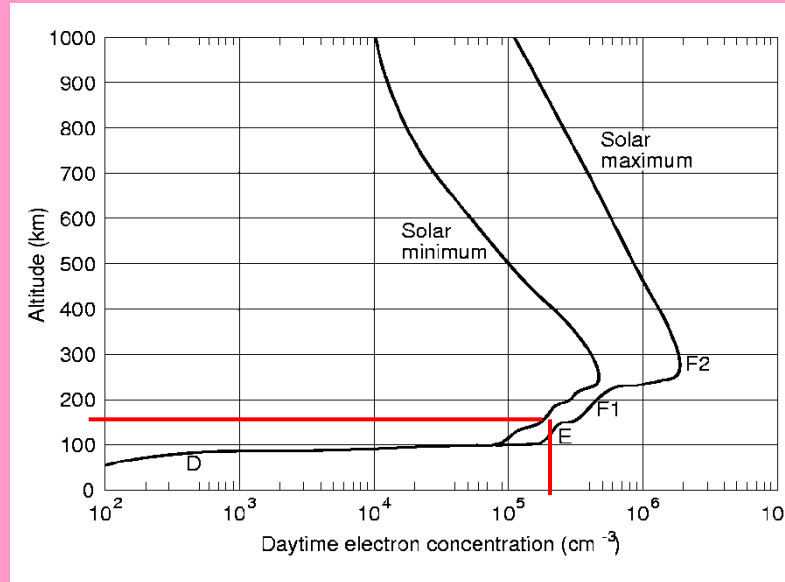
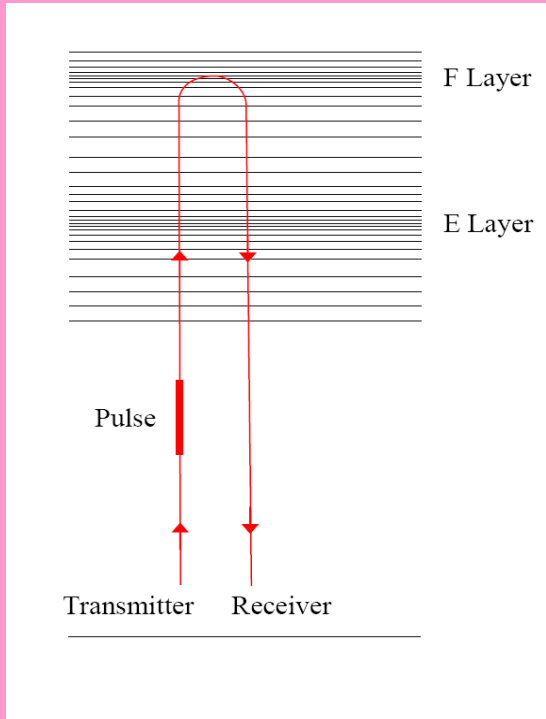


Figure 8 - Ionic composition of solar minimum daytime winter ionosphere (ref. 42).

Mini-groupwork 3

b)



$$f_p = \frac{1}{2\pi} \sqrt{\frac{n_e e^2}{\epsilon_0 m_e}} \approx 9\sqrt{n_e}$$

$$f_p = 5 \cdot 10^6 = 9\sqrt{n_e}$$

⇒

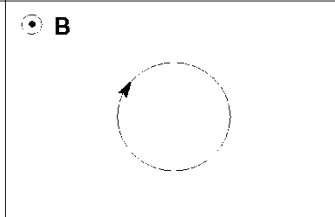
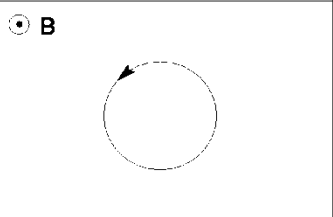
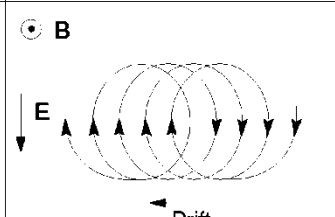
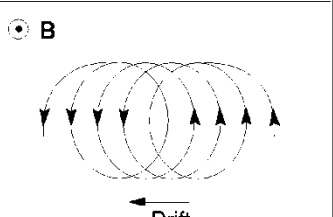
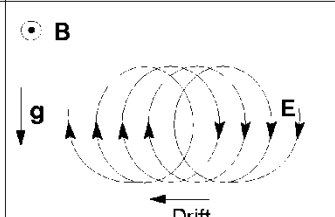
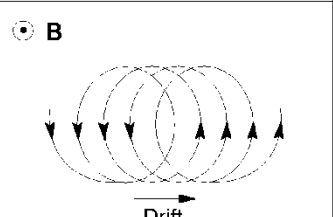
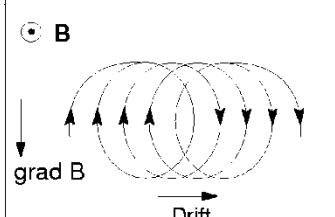
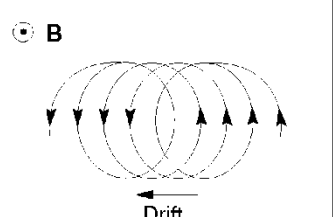
$$n_e = \left(\frac{5 \cdot 10^6}{9} \right)^2 = 3 \cdot 10^{11} \text{ m}^{-3}$$

$$h = 150 \text{ km}$$

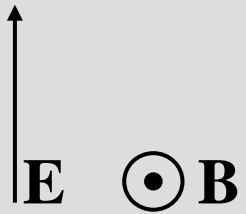
$$t = \frac{2h}{c} = \frac{300 \cdot 10^3}{3 \cdot 10^8} = 10^{-3} \text{ s}$$

Drift motion

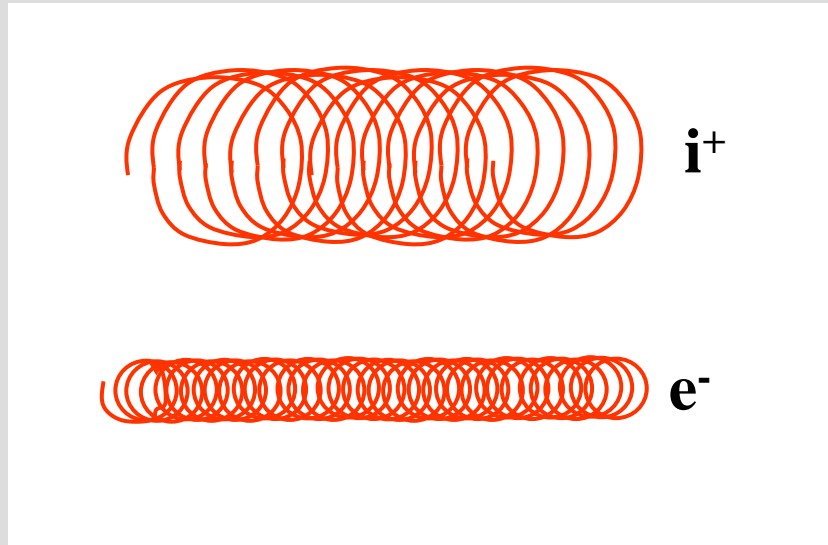
$$\mathbf{u}_{drift} = \frac{\mathbf{F} \times \mathbf{B}}{qB^2}$$

	Positive particles	Negative particles
Homogeneous magnetic field No disturbing force $\mathbf{F} = 0$		
Homogeneous magnetic field Homogeneous electric field $\mathbf{F} = q\mathbf{E}$		
Homogeneous magnetic field Gravitation $\mathbf{F} = m\mathbf{g}$		
Inhomogeneous magnetic field $\mathbf{F} = -\mu \text{ grad } B$		

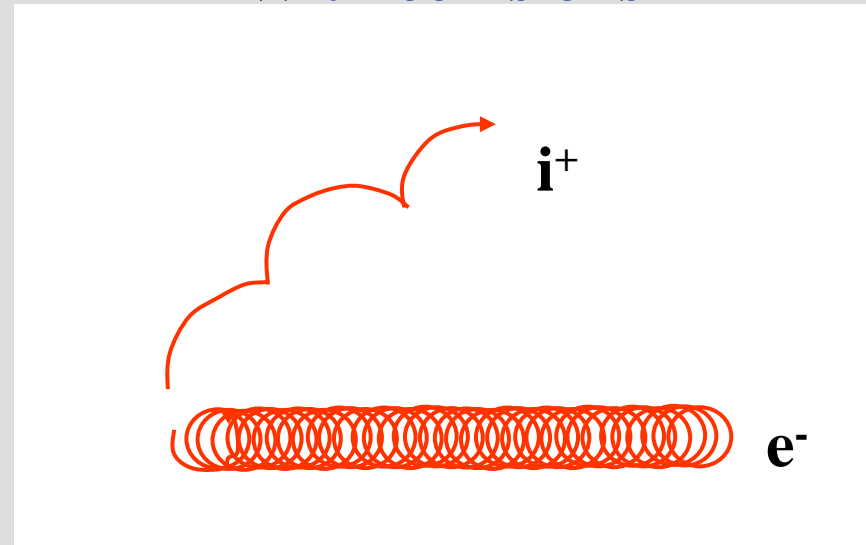
ExB-drift



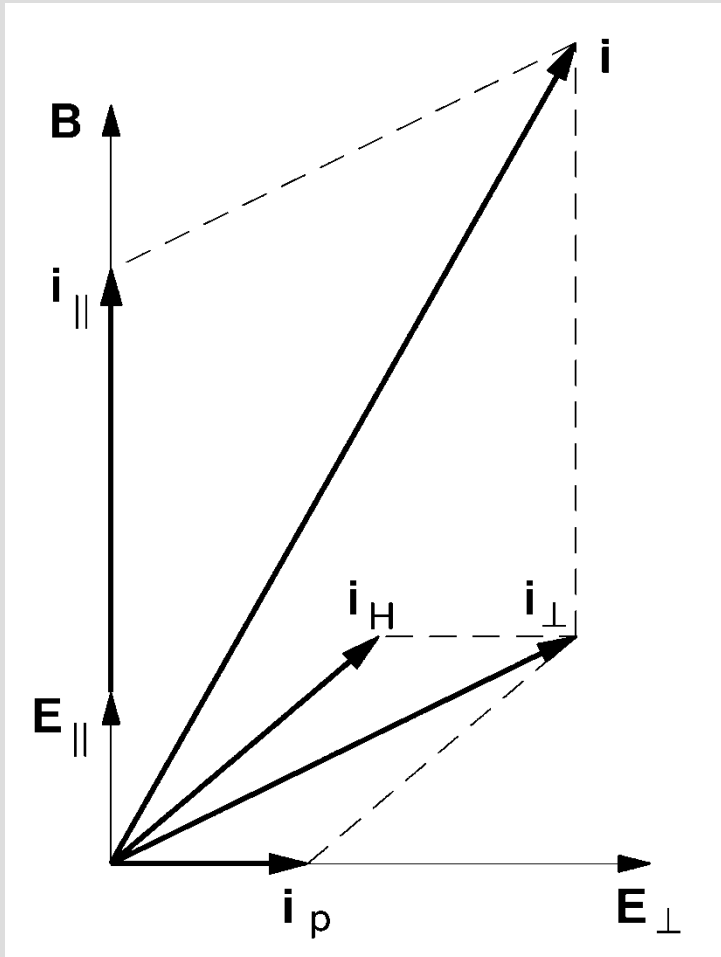
Without collisions



With collisions

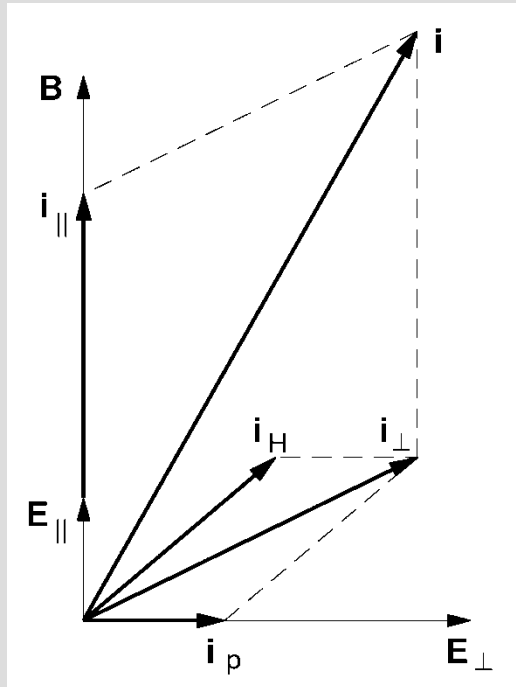


Electric conductivity in a magnetized plasma



- i_{\parallel} = parallel current
- i_p = Pedersen current
- i_H = Hall current

Electric conductivity in a magnetized plasma II



$$\sigma_P = \sigma_e \frac{1}{1 + \omega_{ge}^2 \tau_e^2} + \sigma_i \frac{1}{1 + \omega_{gi}^2 \tau_i^2}$$

$$\sigma_H = \sigma_e \frac{\omega_{ge} \tau_e}{1 + \omega_{ge}^2 \tau_e^2} - \sigma_i \frac{\omega_{gi} \tau_i}{1 + \omega_{gi}^2 \tau_i^2}$$

$$\sigma_{||} = \sigma_e + \sigma_i$$

$$\sigma_e = e^2 n \tau_e / m_e$$

$$\sigma_i = e^2 n \tau_i / m_i$$

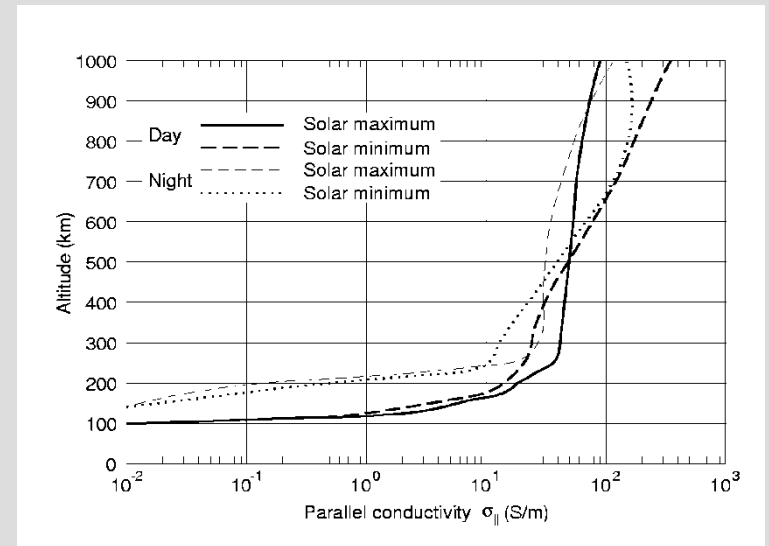
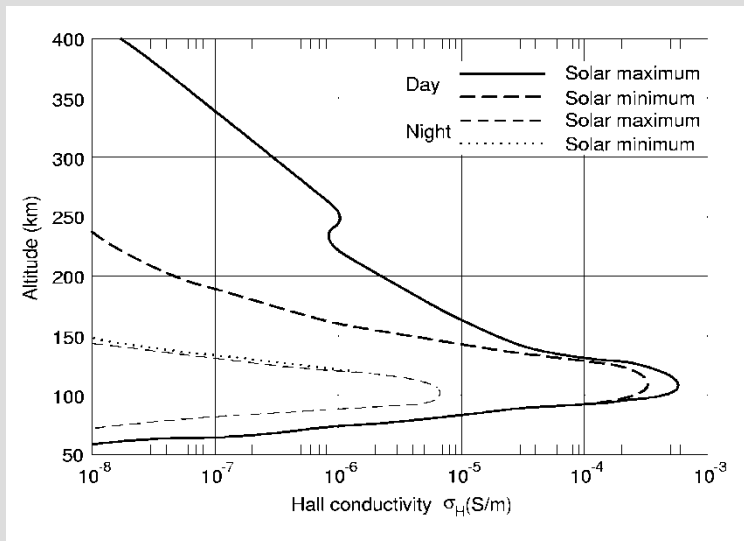
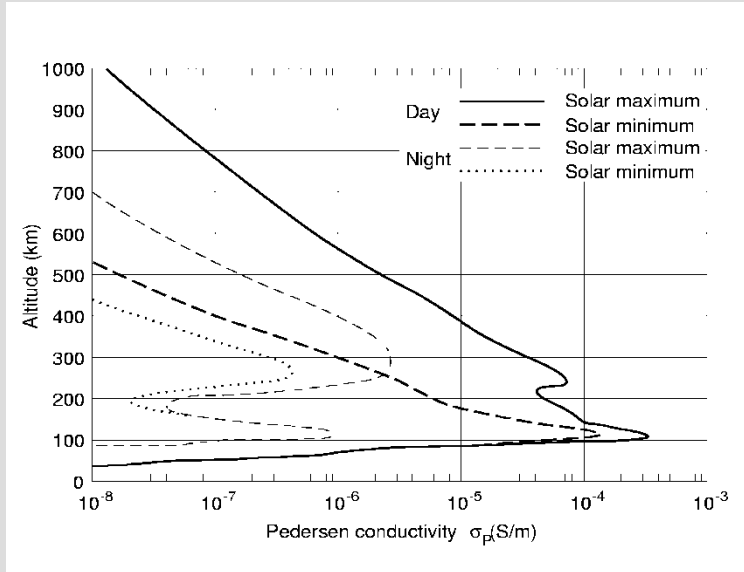
$$i_{||} = \sigma_{||} E_{||}$$

$$\left. \begin{aligned} i_P &= \sigma_P E_{\perp} \\ i_H &= \sigma_H E_{\perp} \end{aligned} \right\}$$

or

$$\mathbf{i} = \mathbf{i}_{||} + \mathbf{i}_{\perp} = \sigma_{||} \mathbf{E}_{||} + \sigma_P \mathbf{E}_{\perp} + \sigma_H \frac{\mathbf{B} \times \mathbf{E}_{\perp}}{B}$$

Ionospheric conductivities



Geomagnetic field

Approximated by a dipole close to Earth.

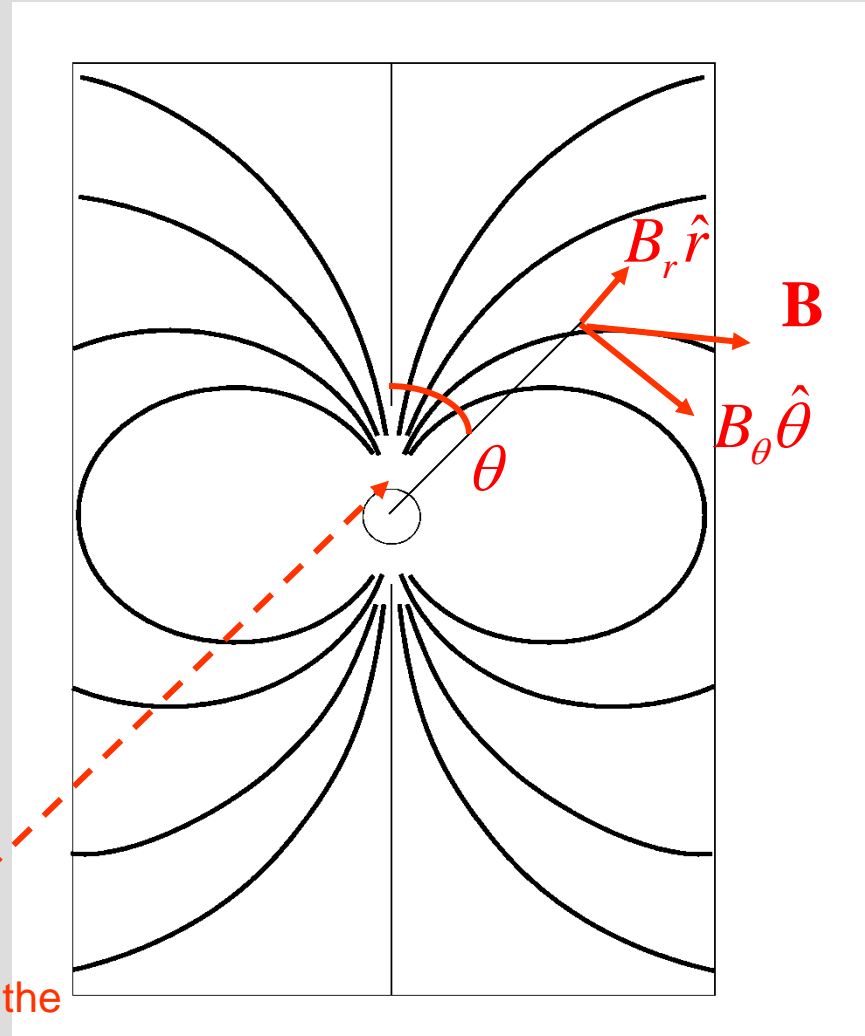
$$B_r = B_p \left(\frac{R_E}{r}\right)^3 \cos \theta$$

$$B_\theta = \frac{B_p}{2} \left(\frac{R_E}{r}\right)^3 \sin \theta$$

$$a = \frac{2\pi R_E^3 B_p}{\mu_0}$$

magnetic dipole moment

Magnetic field at the "north pole"



Geomagnetic field

Alternative formulation of dipole field

$$B_r = B_p \left(\frac{R_E}{r}\right)^3 \cos \theta$$

$$B_\theta = \frac{B_p}{2} \left(\frac{R_E}{r}\right)^3 \sin \theta$$

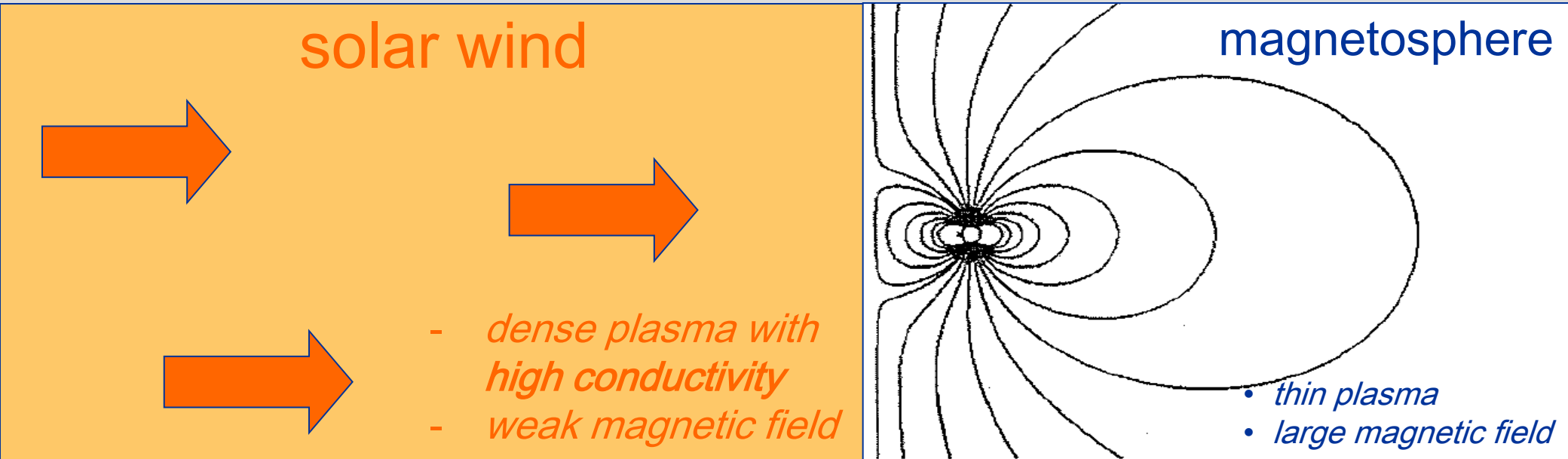
$$B_r = \frac{\mu_0 a}{2\pi} \frac{1}{r^3} \cos \theta$$

$$B_\theta = \frac{\mu_0 a}{2\pi} \cdot \frac{1}{2} \cdot \frac{1}{r^3} \sin \theta$$

$$a = \frac{2\pi R_E^3 B_p}{\mu_0}$$

magnetic dipole moment

Stand-off distance from pressure balance



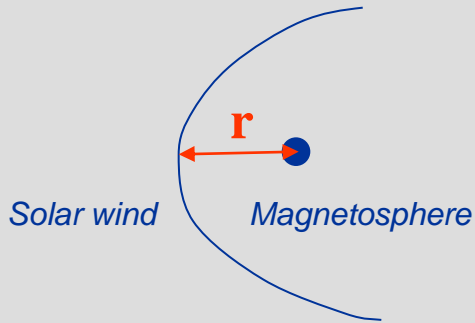
Dynamic pressure:

$$p_d = \rho_{SW} v_{SW}^2$$

Magnetic pressure:

$$p_B = \frac{B^2}{2\mu_0}$$

Magnetopause “stand-off distance”



Dynamic pressure: $p_d = \rho_{SW} v_{SW}^2$

Magnetic pressure: $p_B = \frac{1}{2\mu_0} B^2$

Dipole field strength
(in equatorial plane): $B = \frac{\mu_0 a}{4\pi} \frac{1}{r^3}$

$$p_d = p_B \Rightarrow \rho_{SW} v_{SW}^2 = \left[\frac{\mu_0 a}{4\pi} \frac{1}{r^3} \right]^2 / 2\mu_0 \Rightarrow$$

$$r = \left(\frac{\mu_0 a}{4\pi} \right)^{1/3} \left(2\mu_0 \rho_{SW} v_{SW}^2 \right)^{-1/6}$$

$a = 8 \times 10^{22} \text{ Am}^2$, $v = 500 \text{ km/s}$, $\rho_{SW} = 10^7 \times 1.7 \times 10^{-27} \text{ kg/m}^3$:

$r = 7 R_e$ (1 $R_e = 6378 \text{ km}$)



Adiabatic invariant

DEFINITION:

An **adiabatic invariant** is a property of a physical system which stays constant when changes are made slowly.

By 'slowly' in the context of charged particle motion in magnetic fields, we mean much slower than the gyroperiod.

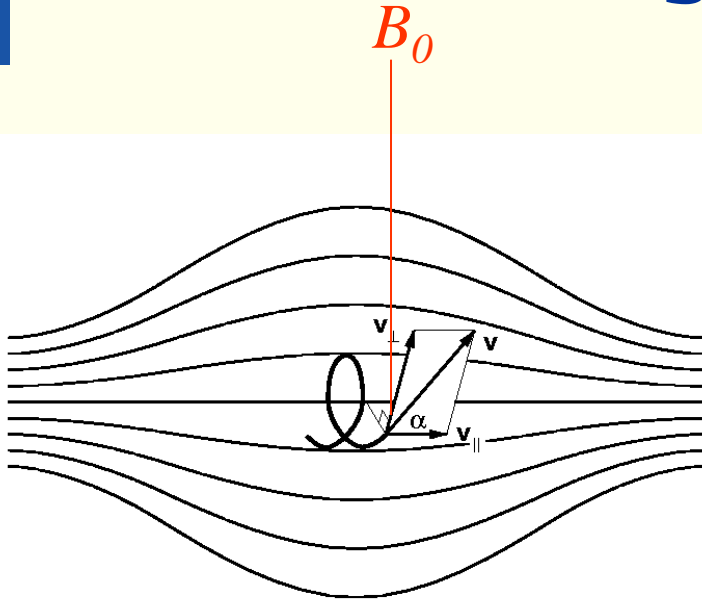
'First adiabatic invariant' of particle drift:

$$\mu = \frac{mv_{\perp}^2}{2B}$$

Magnetic mirror

$mv^2/2$ constant (energy conservation) 

$$\frac{\sin^2 \alpha}{B} = \textit{konst}$$



The magnetic moment μ is an *adiabatic invariant*.

$$\mu = \frac{mv_{\perp}^2}{2B} = \frac{mv^2 \sin^2 \alpha}{2B}$$

Red

α increases

Yellow

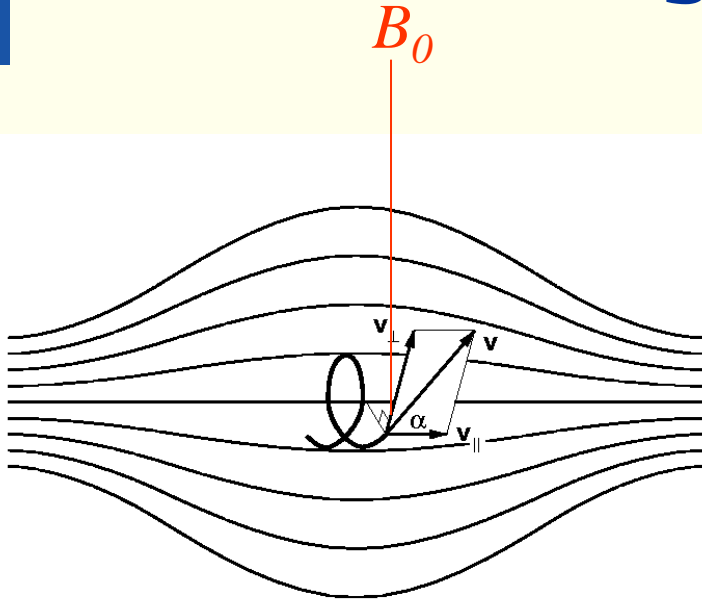
α decreases

What happens with α as the particle moves into the stronger magnetic field?

Magnetic mirror

$mv^2/2$ constant (energy conservation) 

$$\frac{\sin^2 \alpha}{B} = konst$$



The magnetic moment μ is an *adiabatic invariant*.

$$\mu = \frac{mv_{\perp}^2}{2B} = \frac{mv^2 \sin^2 \alpha}{2B}$$

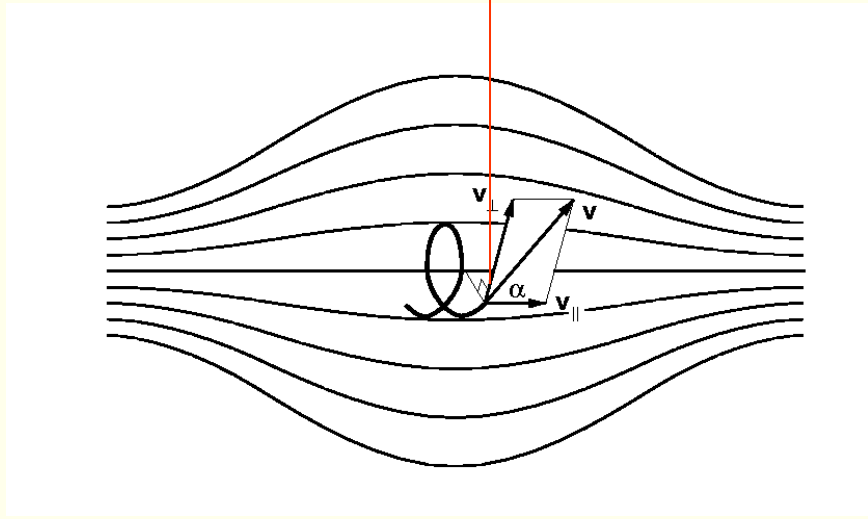
What happens with α as the particle moves into the stronger magnetic field?

$$\sin \alpha = \sqrt{B \cdot konst}$$

Red

α increases

Magnetic mirror



$mv^2/2$ constant (energy conservation) →

$$\frac{\sin^2 \alpha}{B} = \text{konst}$$

particle turns when $\alpha = 90^\circ$ →

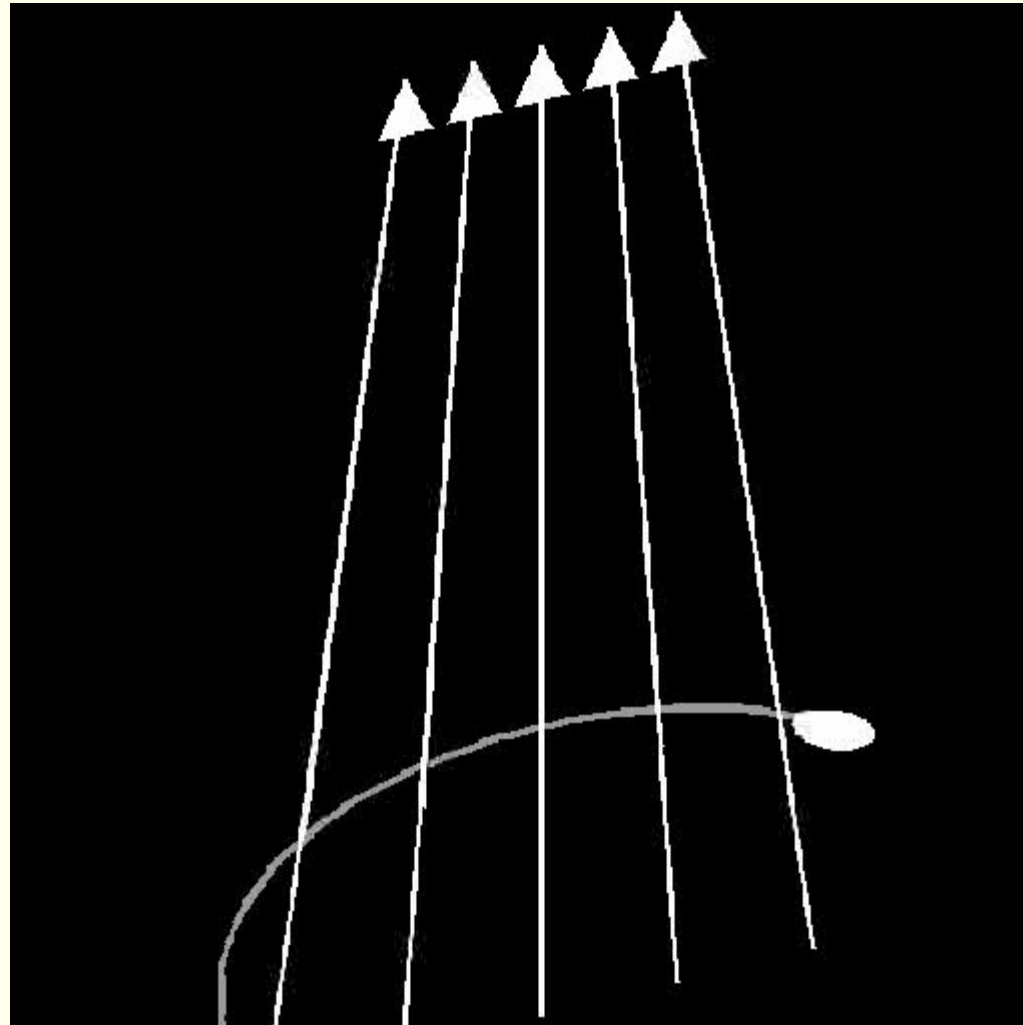
$$\frac{B_{\text{turn}}}{\sin^2 90^\circ} = \frac{B_0}{\sin^2 \alpha} \rightarrow$$

$$B_{\text{turn}} = \frac{B_0}{\sin^2 \alpha}$$

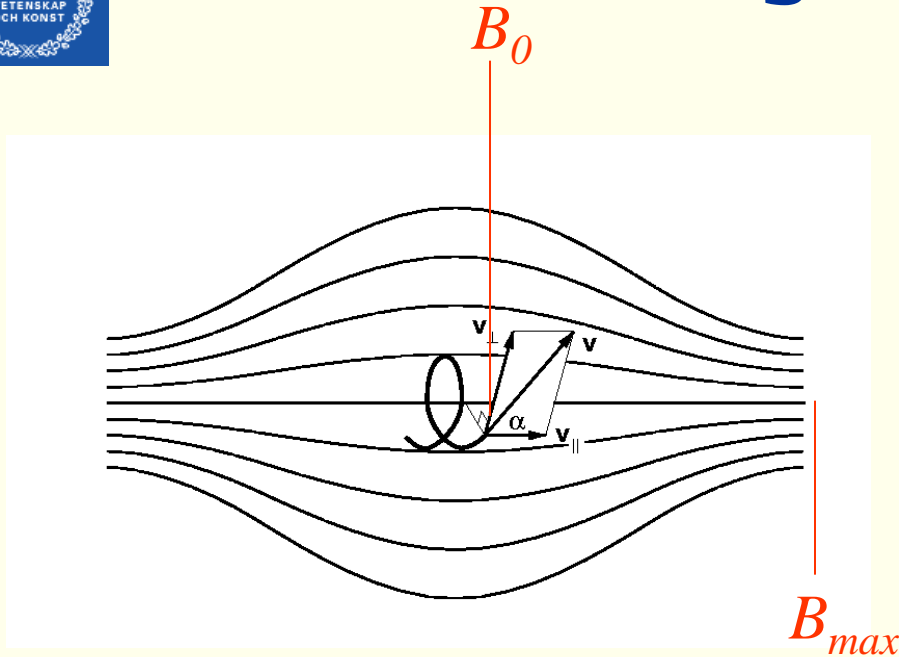
The magnetic moment μ is an *adiabatic invariant*.

$$\mu = \frac{mv_{\perp}^2}{2B} = \frac{mv^2 \sin^2 \alpha}{2B}$$

Magnetic mirror



Magnetic mirror



$mv^2/2$ constant (energy conservation) →

$$\frac{\sin^2 \alpha}{B} = \text{konst}$$

particle turns when $\alpha = 90^\circ$ →

$$B_{\text{turn}} = B_0 / \sin^2 \alpha$$

If maximal B -field is B_{max} a particle with pitch angle α can only be turned around if

$$B_{\text{turn}} = B_0 / \sin^2 \alpha \leq B_{\text{max}} \rightarrow$$

$$\alpha > \alpha_{lc} = \arcsin \sqrt{B_0 / B_{\text{max}}}$$

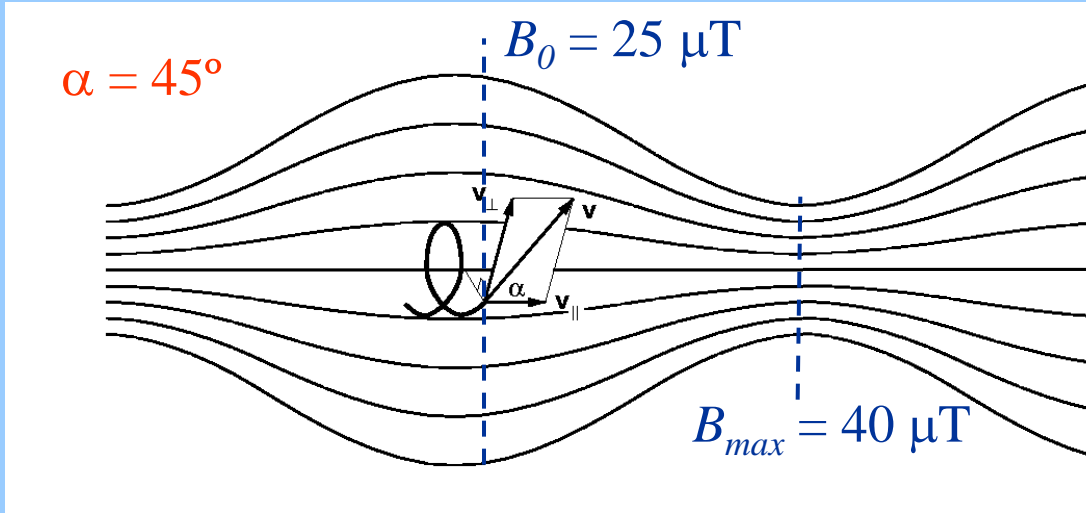
Particles in
loss cone :

$$\alpha < \alpha_{lc}$$

The magnetic moment μ is an *adiabatic invariant*.

$$\mu = \frac{mv_{\perp}^2}{2B} = \frac{mv^2 \sin^2 \alpha}{2B}$$

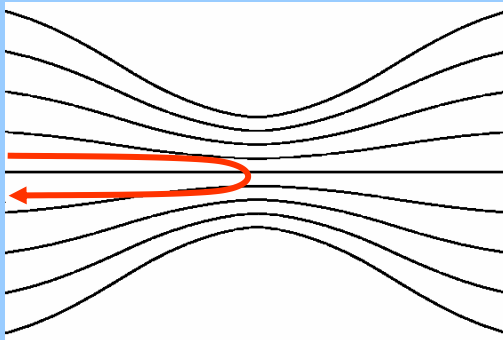
What will happen to the particle?



$$\alpha_{lc} = \arcsin \sqrt{B_0 / B_{max}}$$

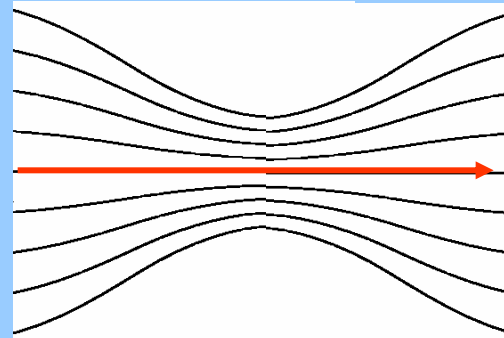
Blue

It will mirror

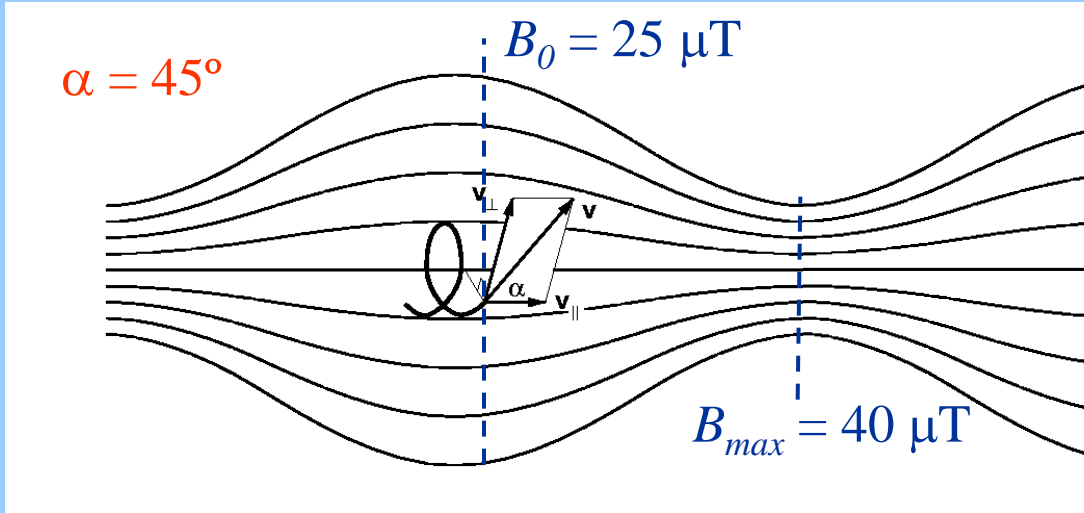


Yellow

It will escape



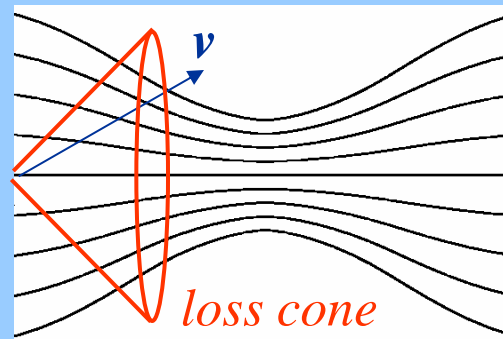
What will happen to the particle?



$$\alpha_{lc} = \arcsin \sqrt{B_0 / B_{max}} =$$

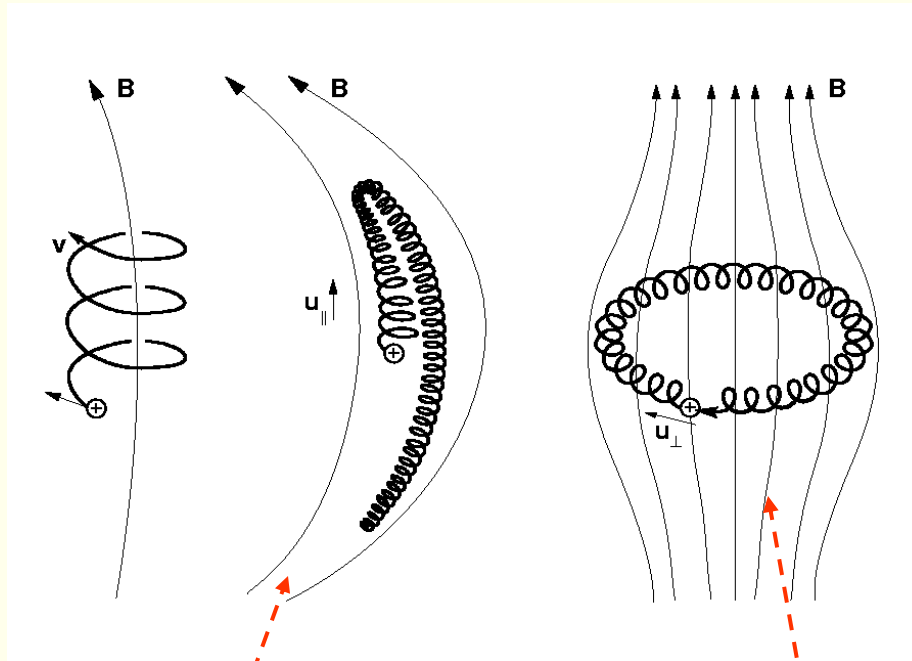
$$\arcsin \sqrt{25 / 40} = 52^\circ$$

Yellow It will escape



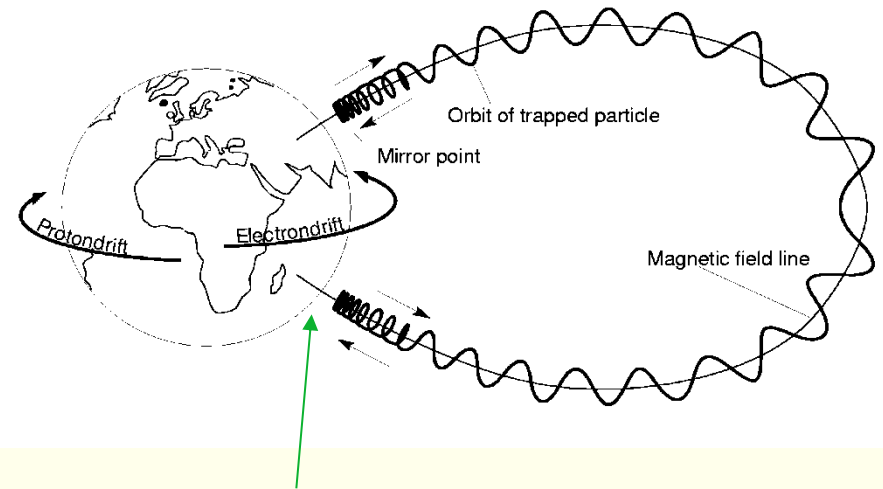
Particle motion in geomagnetic field

longitudinal gyration oscillation azimuthal drift



Magnetic mirror

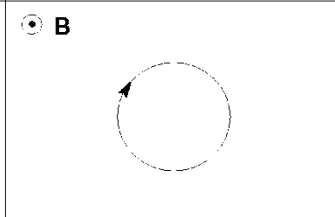
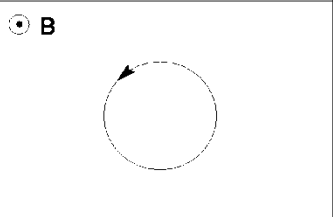
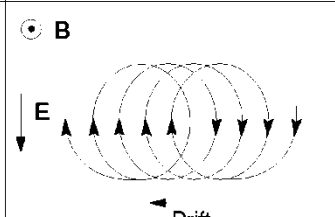
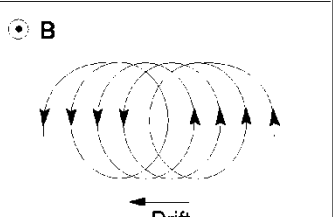
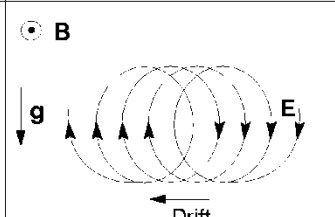
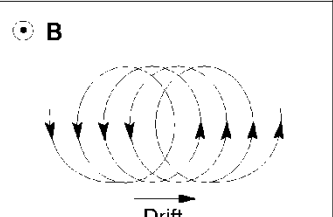
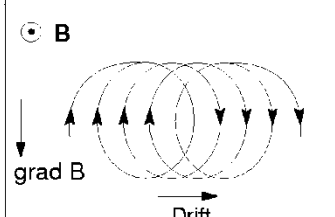
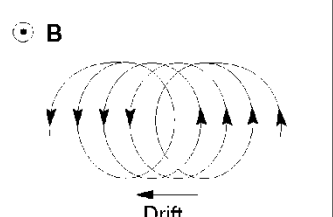
grad B drift



Particles in the loss cone create the aurora!

Drift motion

$$\mathbf{u} = \frac{\mathbf{F} \times \mathbf{B}}{qB^2}$$

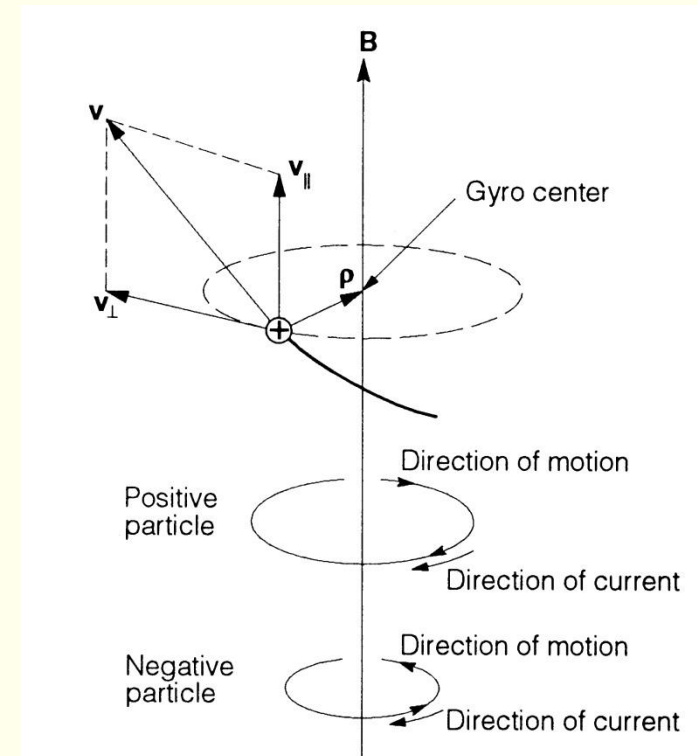
	Positive particles	Negative particles
Homogeneous magnetic field No disturbing force $\mathbf{F} = 0$		
Homogeneous magnetic field Homogeneous electric field $\mathbf{F} = q\mathbf{E}$		
Homogeneous magnetic field Gravitation $\mathbf{F} = m\mathbf{g}$		
Inhomogeneous magnetic field $\mathbf{F} = -\mu \text{ grad } B$		

Force on magnetic dipole

$$\boldsymbol{\mu} \sim -\mathbf{B} \Rightarrow \boldsymbol{\mu} = -\mu \frac{\mathbf{B}}{B}$$

$$\mathbf{F} = \nabla (\boldsymbol{\mu} \cdot \mathbf{B}) = -\mu \nabla \left(\frac{\mathbf{B}}{B} \cdot \mathbf{B} \right) =$$

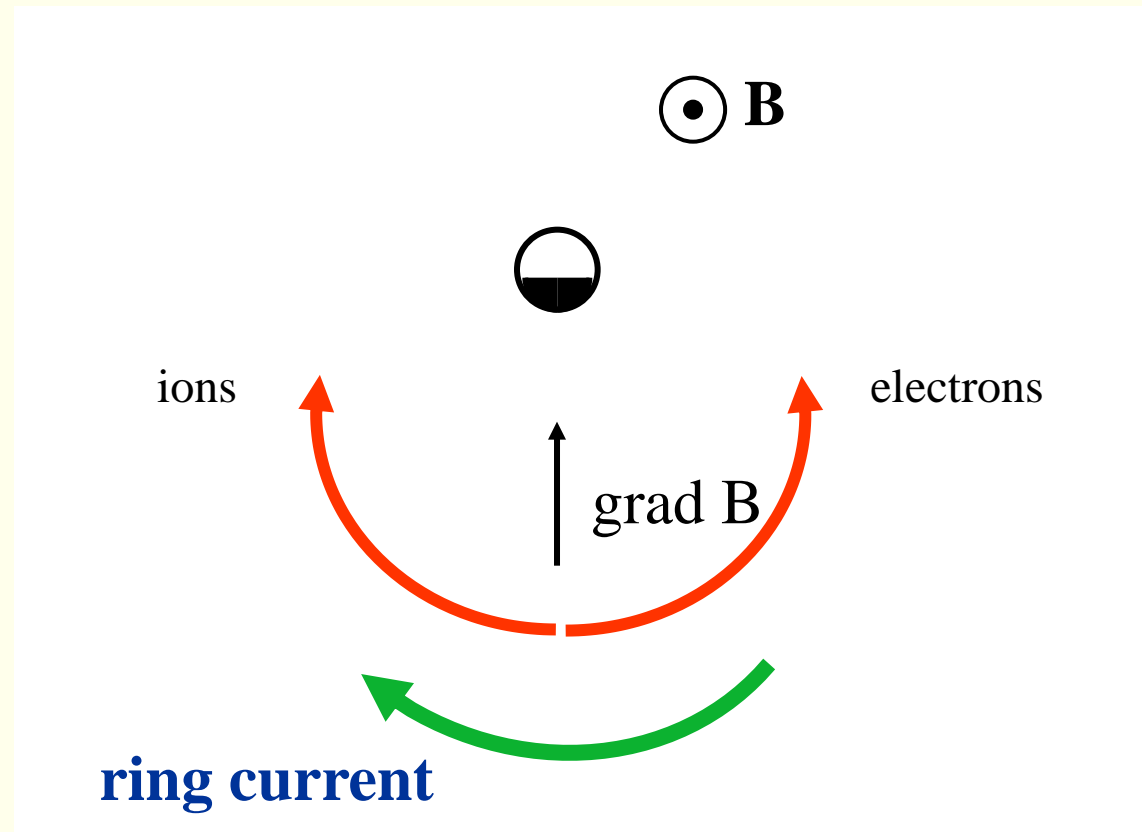
$$= -\mu \nabla \left(\frac{B^2}{B} \right) = -\mu \nabla B$$



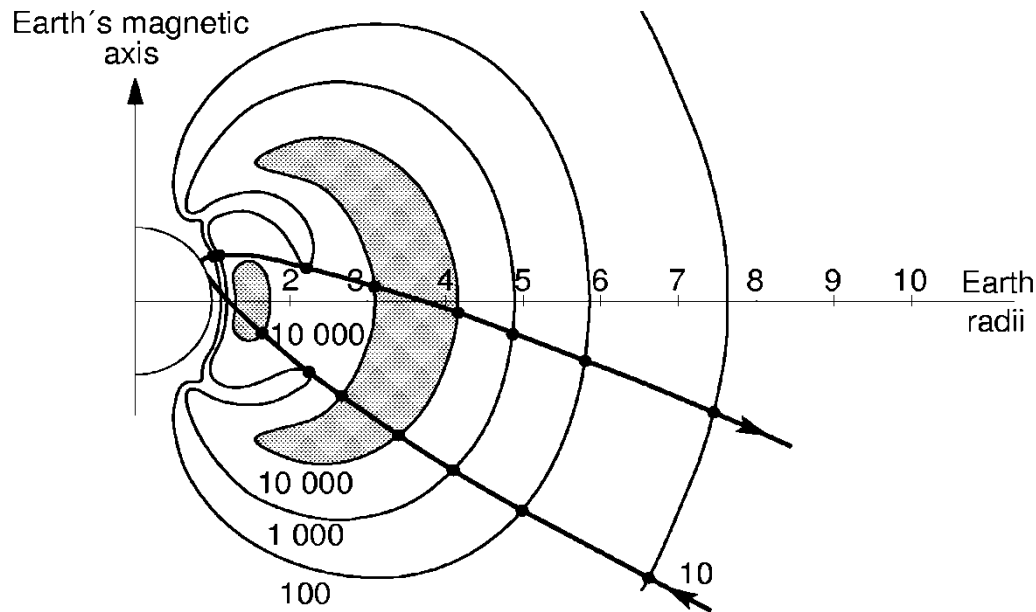
Ring current and particle motion

$$\mathbf{u} = -\frac{\mu \nabla B \times \mathbf{B}}{qB^2}$$

$$\mu = \frac{mv_{\perp}^2}{2B}$$



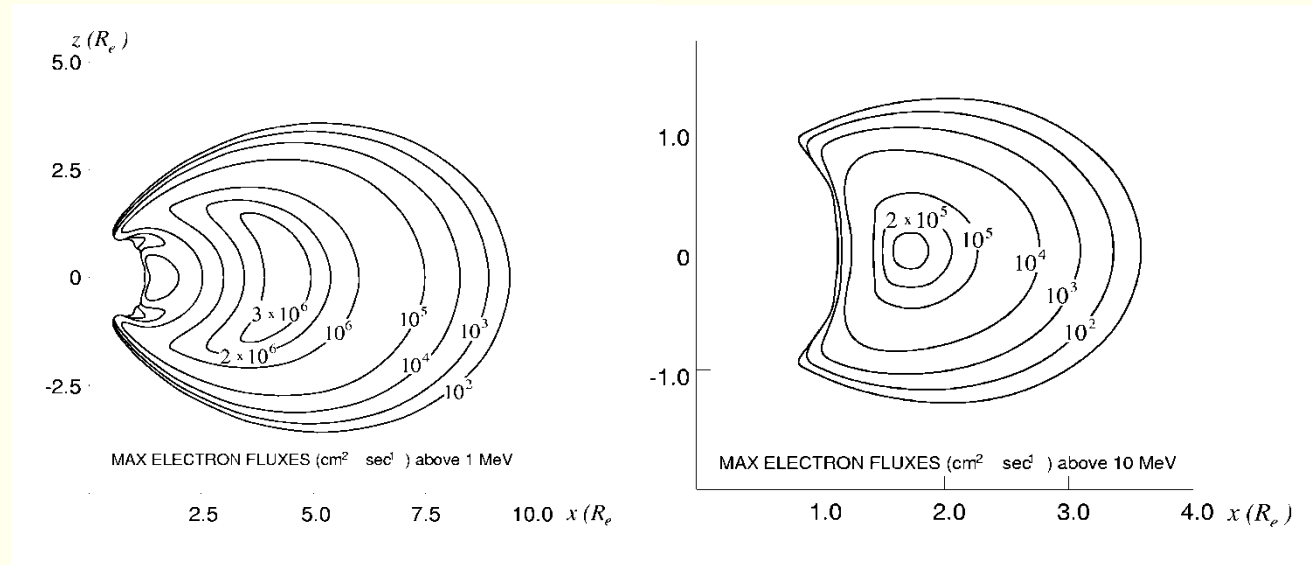
Radiation belts



I. Van Allen belts

- Discovered in the 50s , Explorer 1
- Inner belt contains protons with energies of ~30 MeV
- Outer belt (Explorer IV, Pioneer III): electrons, $W > 1.5$ MeV

Radiation belts



- At lower energies there is a more or less continuous population of energetic particles in the inner magnetosphere. (Inner part of *plasma sheet*)
- source: **CRAND** (Cosmic Ray Albedo Neutron Decay).
- a danger for satellites and astronauts.
- associated with a current (*ring current*) which distorts the inner part of the geomagnetic field.

CRAND (Cosmic Ray Albedo Neutron Decay)

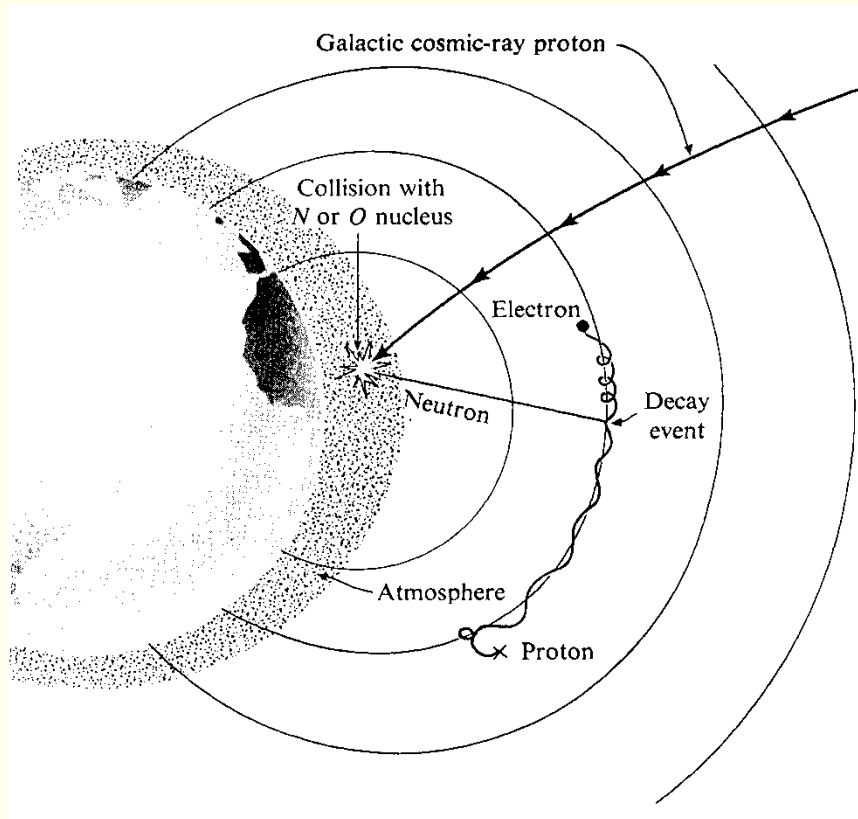
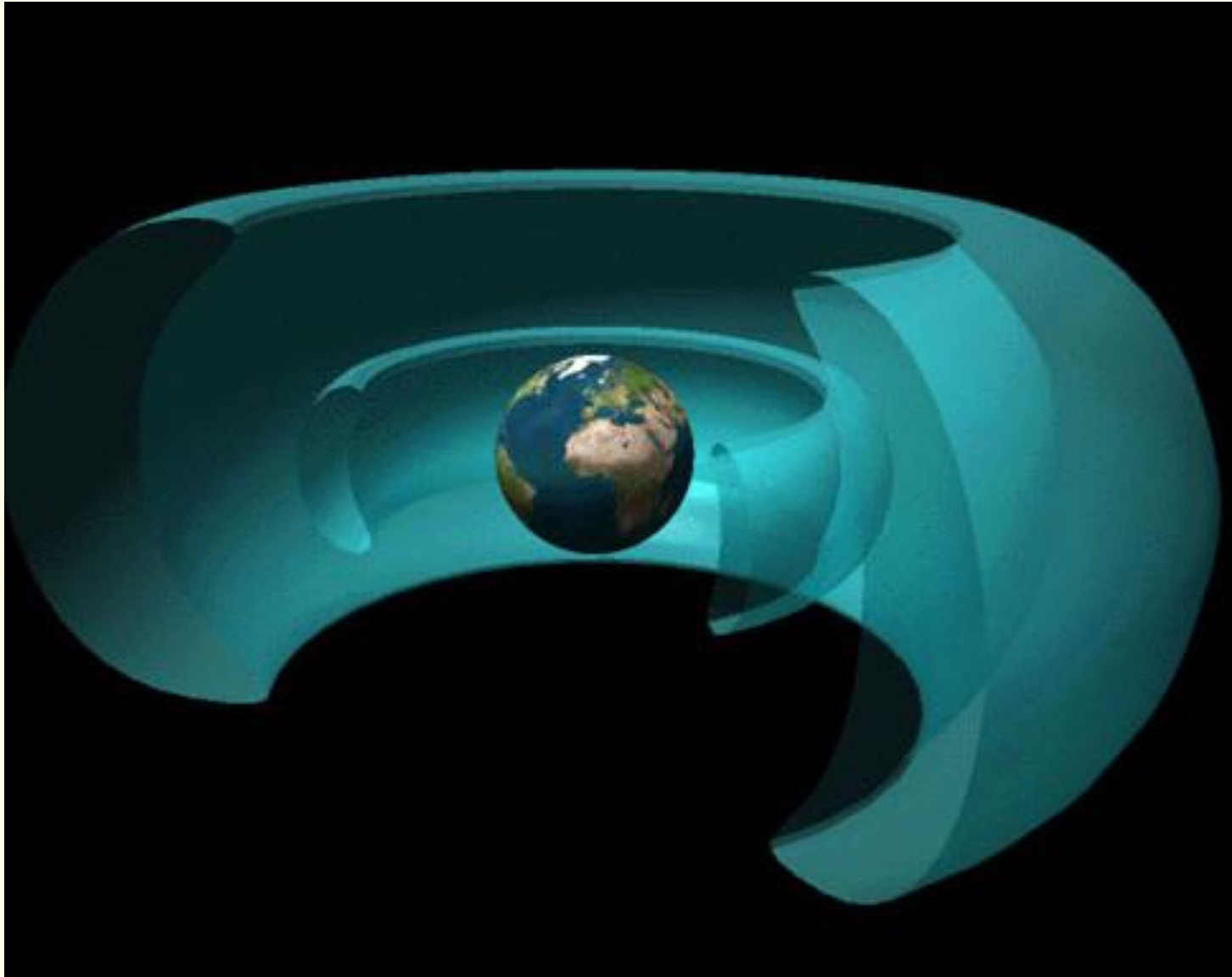


Figure 8. An illustration of the CRAND process for populating the inner radiation belts [Hess, 1968].

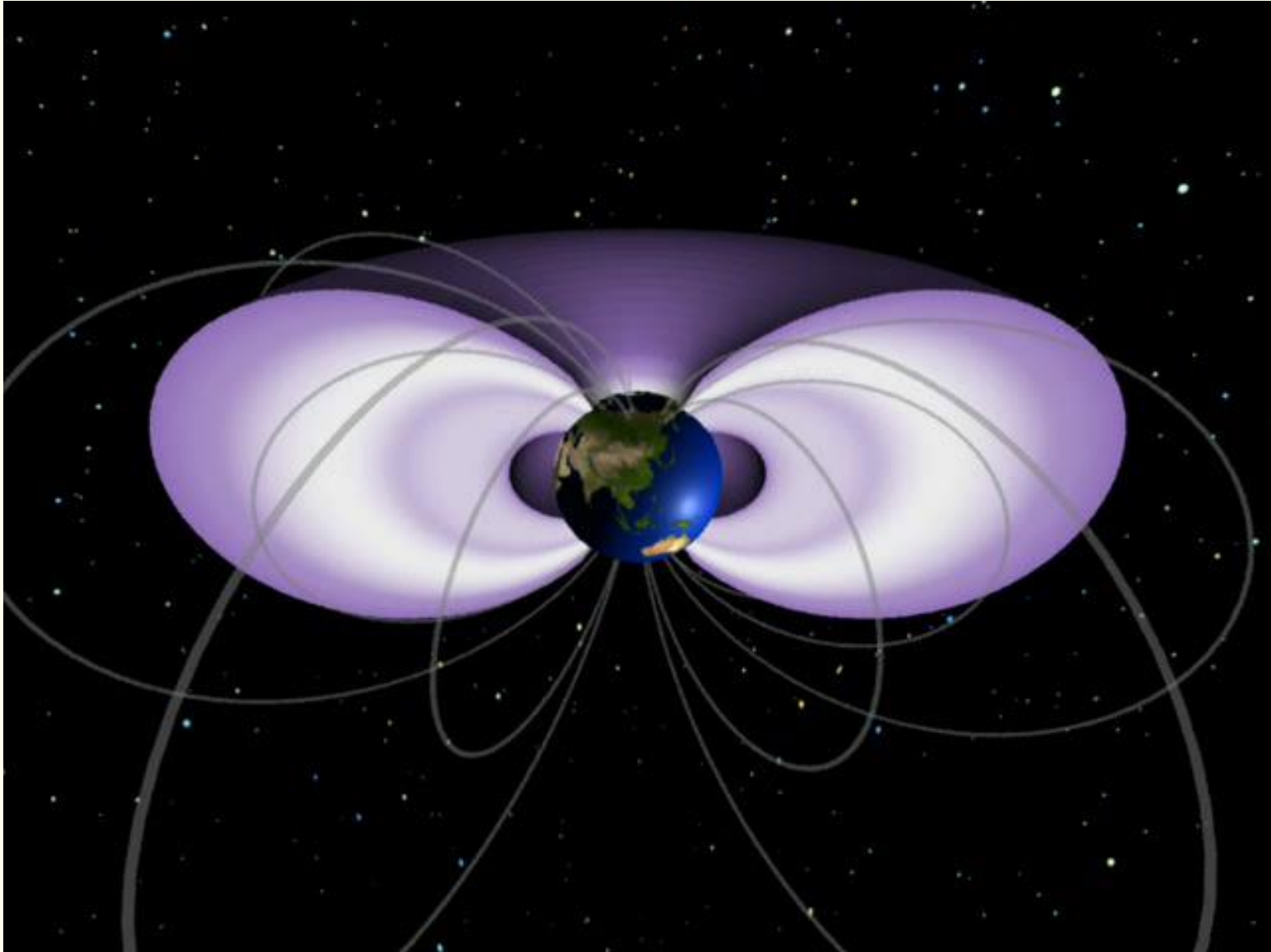
Collisions between cosmic ray particles and the Earth create new particles. Among these are neutrons, that are not affected by the magnetic field. They decay, soon after they happen to be in the radiation belts. The resulting protons and electrons are trapped in the radiation belts.

This contribution to the radiation belts are called the ***neutron albedo***.

Radiation belts



Radiation belts

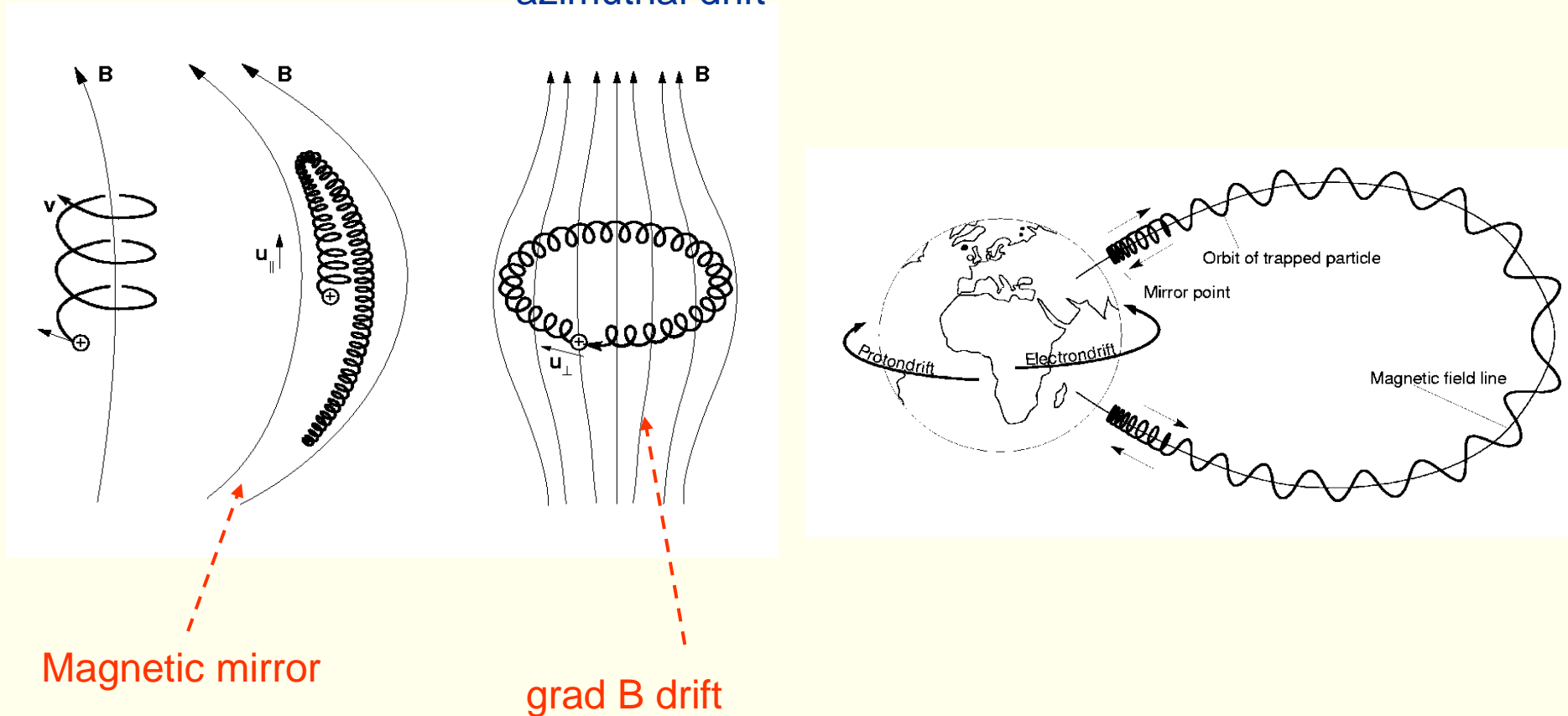


Particle motion in geomagnetic field

longitudinal oscillation

gyration

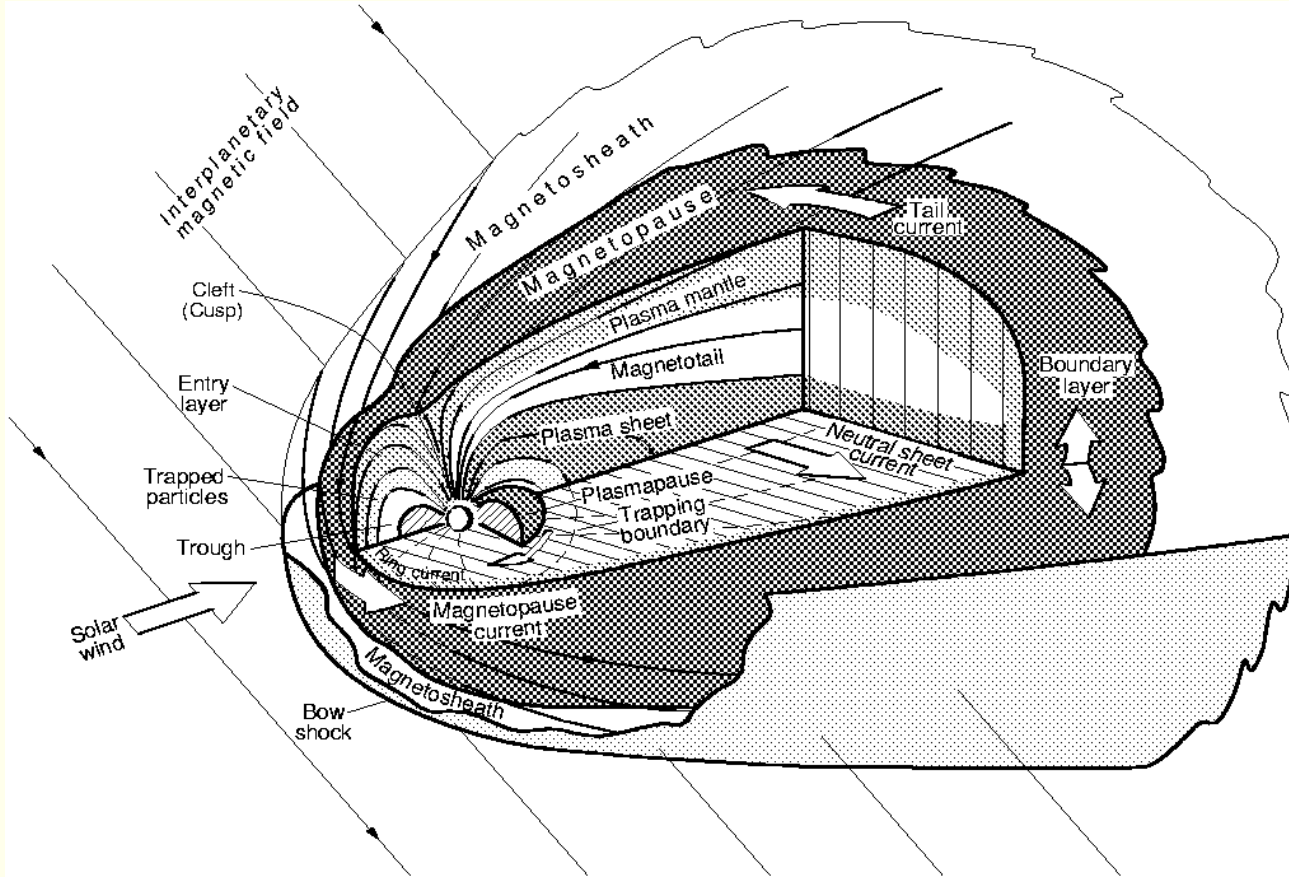
azimuthal drift



Magnetic mirror

grad B drift

Structure of magnetosphere

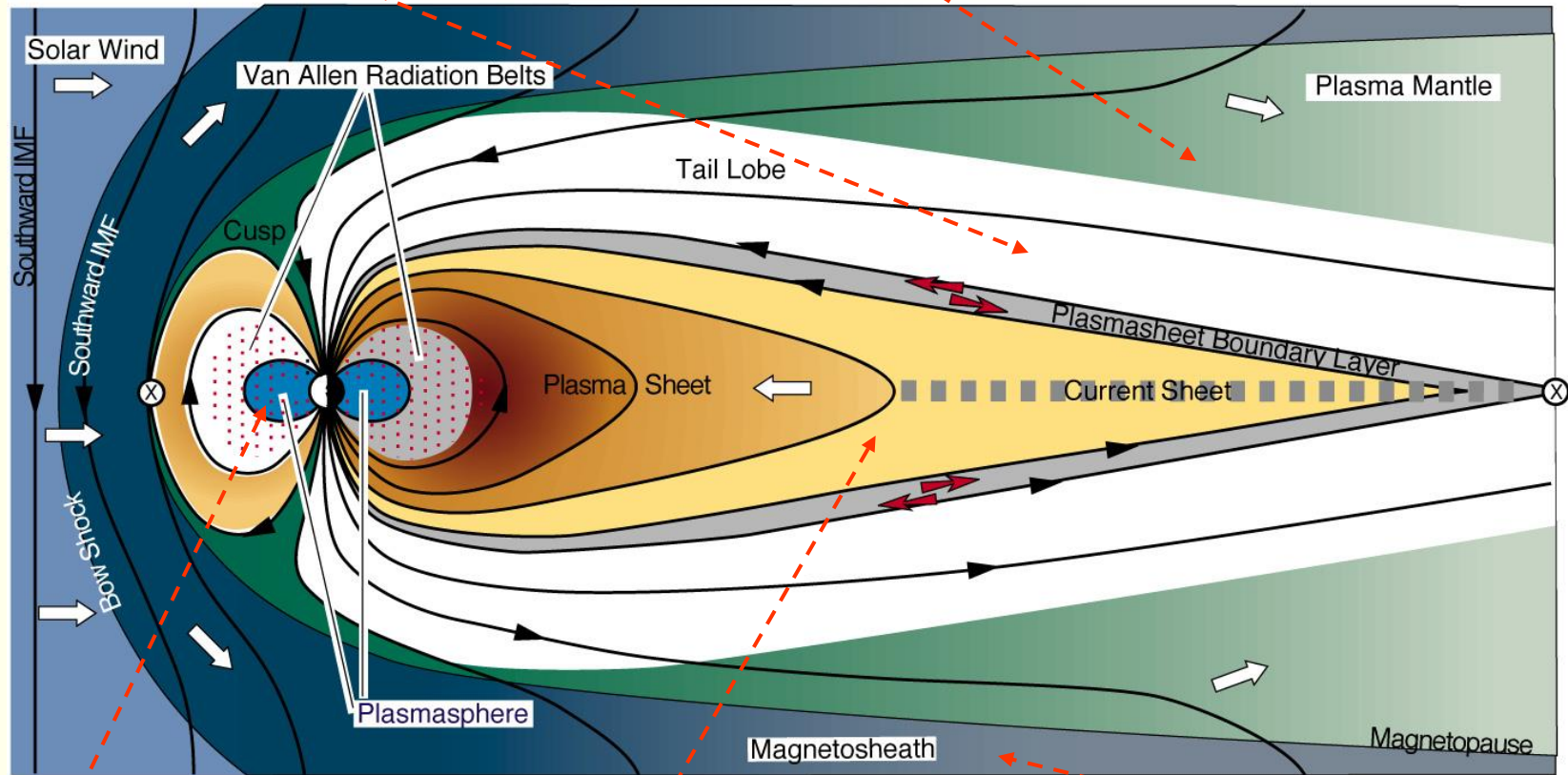


- The plasma in the is made up of approximately equal parts of H^+ and O^+ .
- Plasma populations organized by geomagnetic field.
- Particles will mirror between northern and southern hemispheres on closed field lines

Magnetospheric structure

polar plumes = tail lobe
 $n_e \sim 0,01 \text{ cm}^{-3}$, $T_e \sim 10^6 \text{ K}$

plasma mantle
 $n_e \sim 0,1-1 \text{ cm}^{-3}$, $T_e \sim 10^6 \text{ K}$

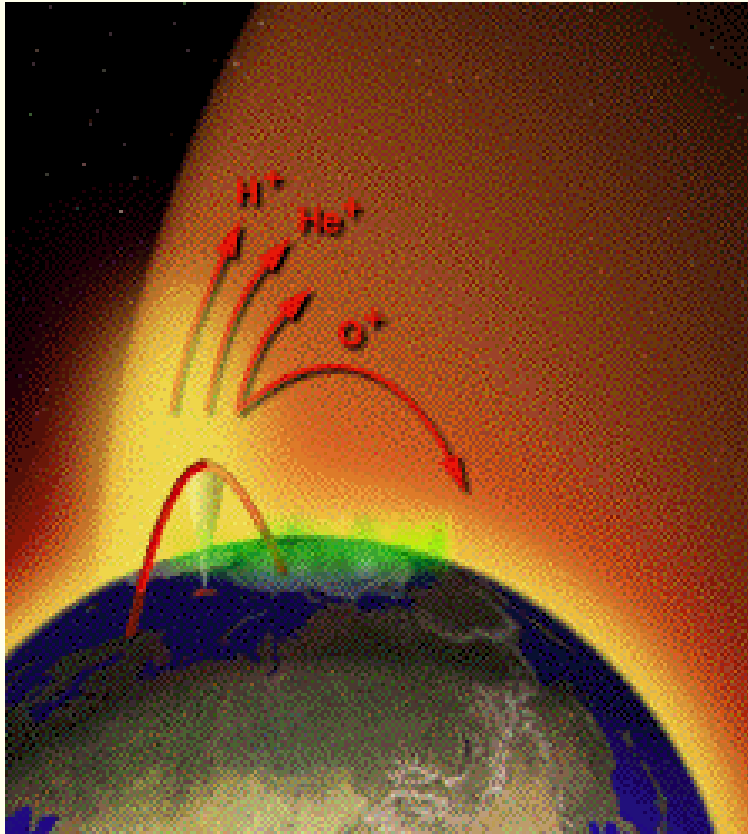


plasmasphere:
 $n_e \sim 10-100 \text{ cm}^{-3}$, $T_e \sim 1000 \text{ K}$

plasma sheet:
 $n_e \sim 1 \text{ cm}^{-3}$, $T_e \sim 10^7 \text{ K}$

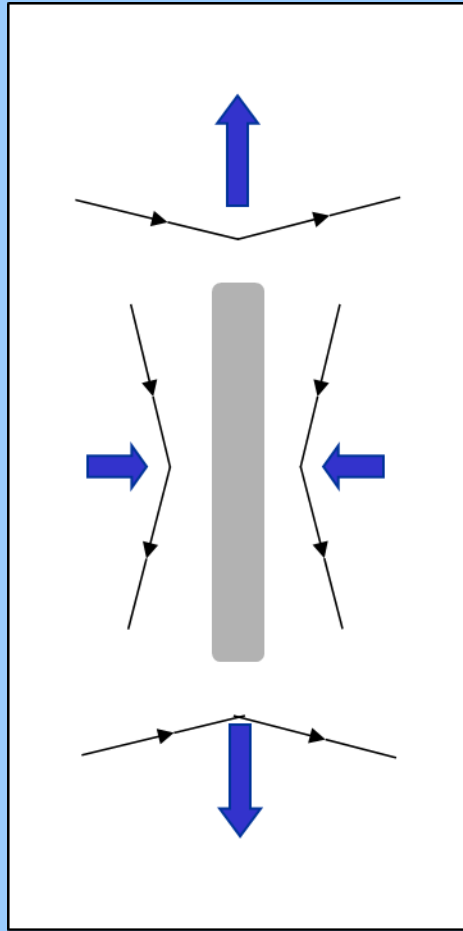
magnetosheath:
 $n_e \sim 5 \text{ cm}^{-3}$, $T_e \sim 10^6 \text{ K}$

Outflow from the ionosphere

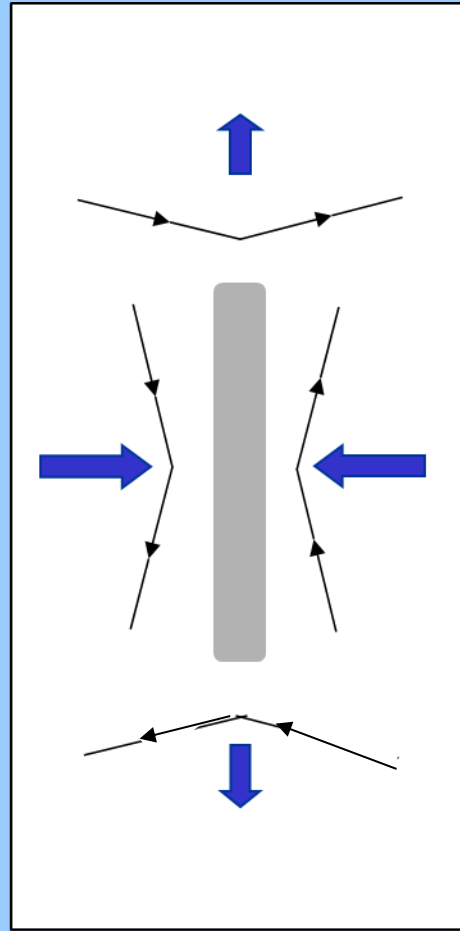


An important source for the magnetospheric plasma.
Research is ongoing.

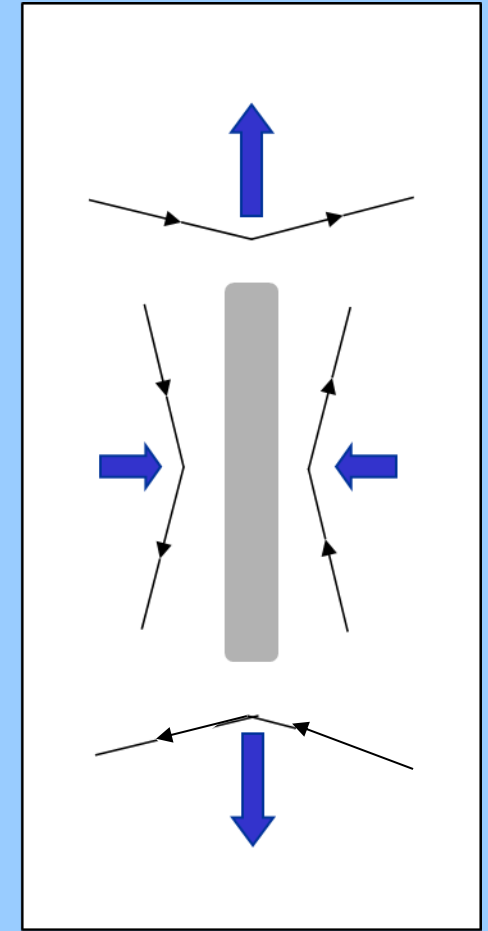
Magnetic reconnection



Green

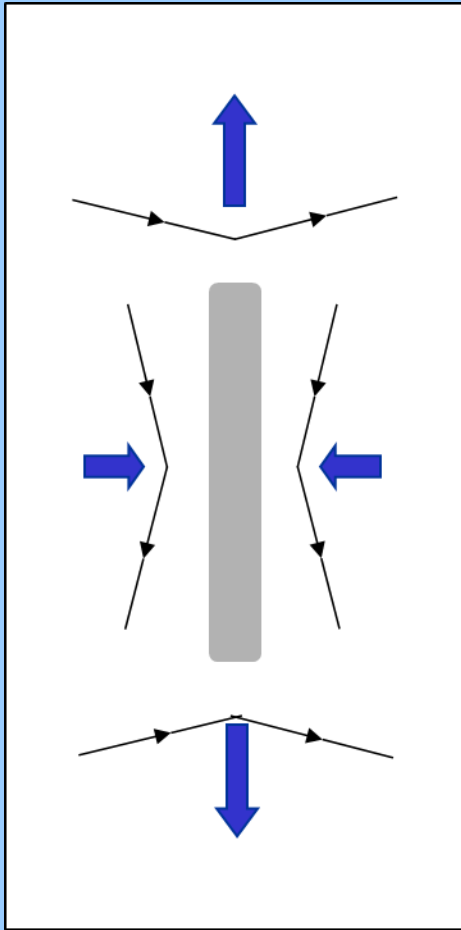


Yellow

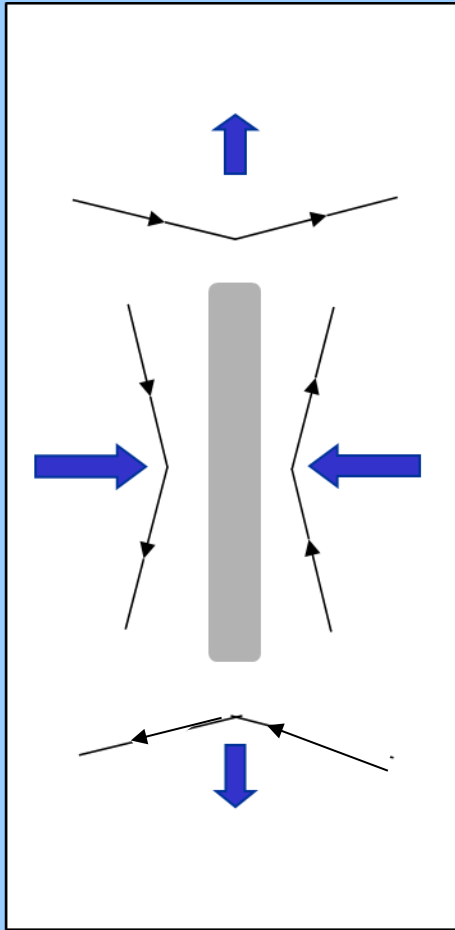


Red

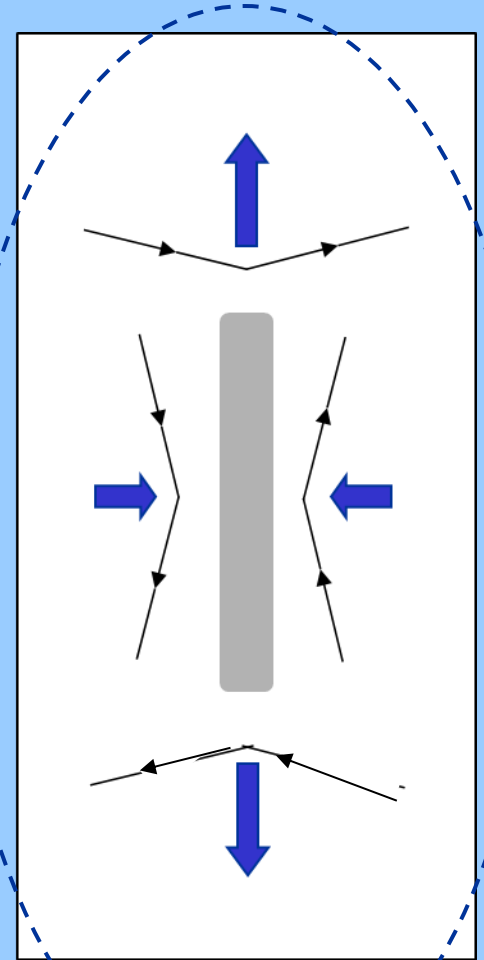
Magnetic reconnection



Green

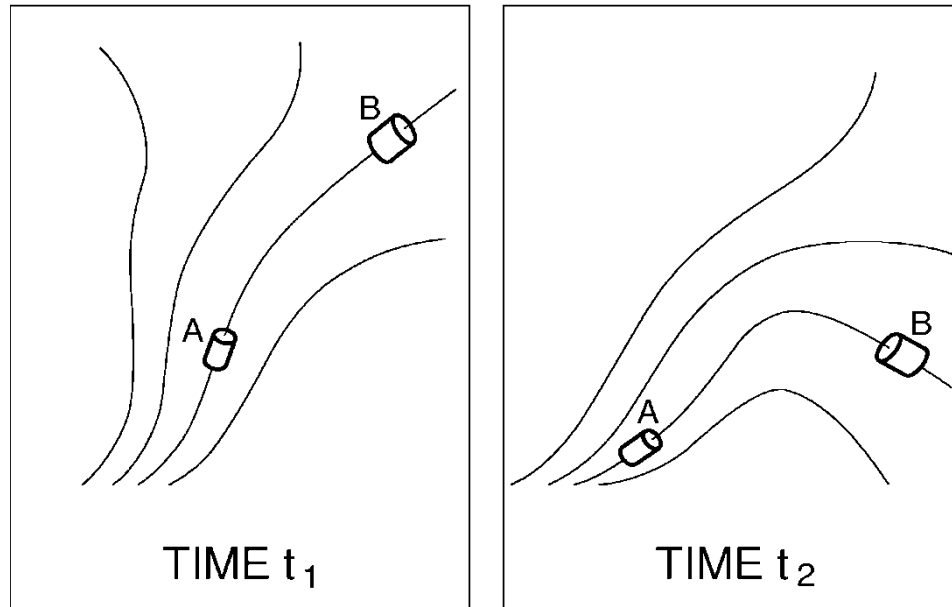


Yellow



Red

Frozen in magnetic field lines

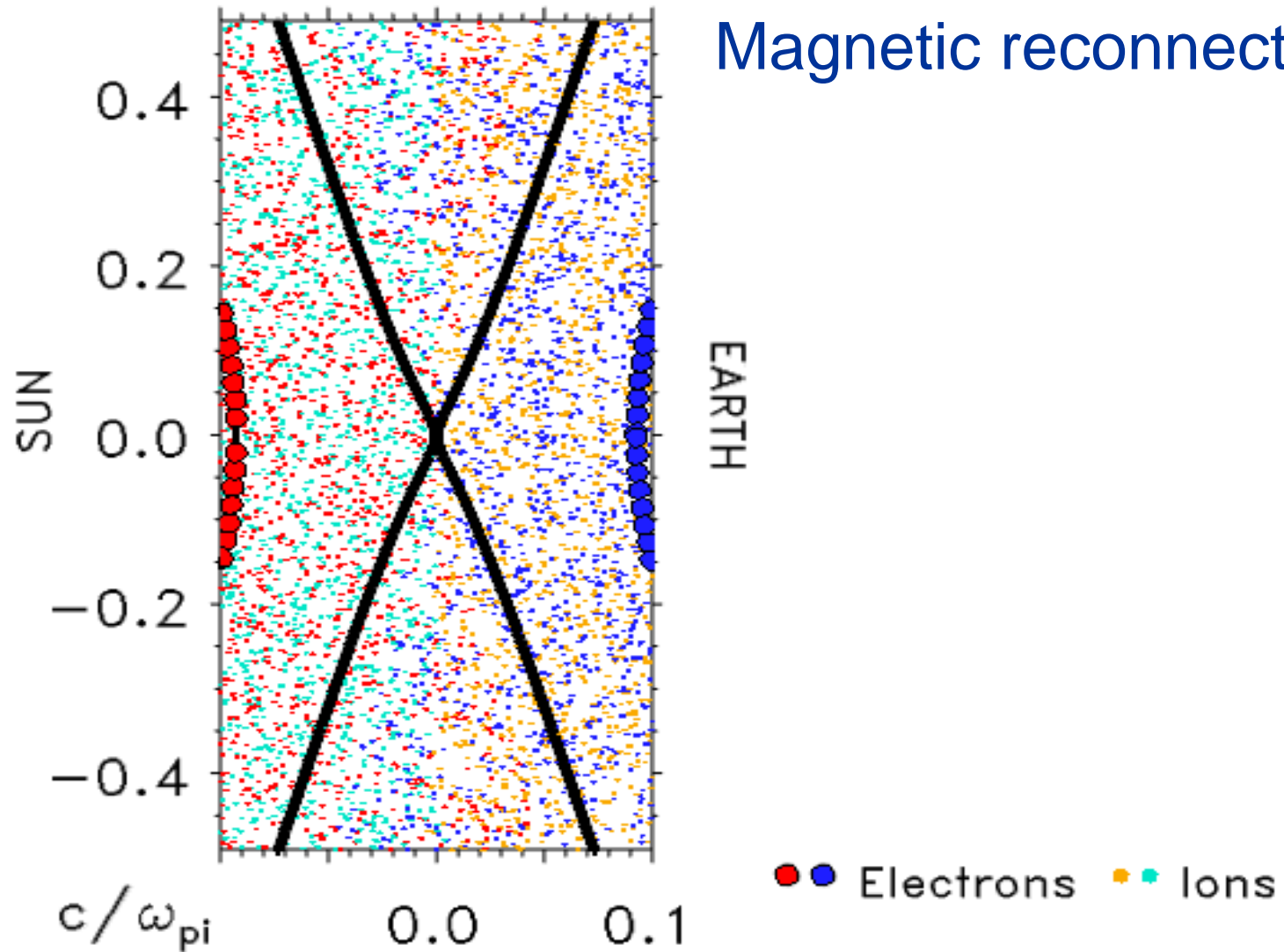


In fluid description of plasma two plasma elements that are connected by a common magnetic field line at time t_1 will be so at any other time t_2 .

This applies if the magnetic Reynolds number is large:

$$R_m = \mu_0 \sigma l_c v_c \gg 1$$

An example of the collective behaviour of plasmas.



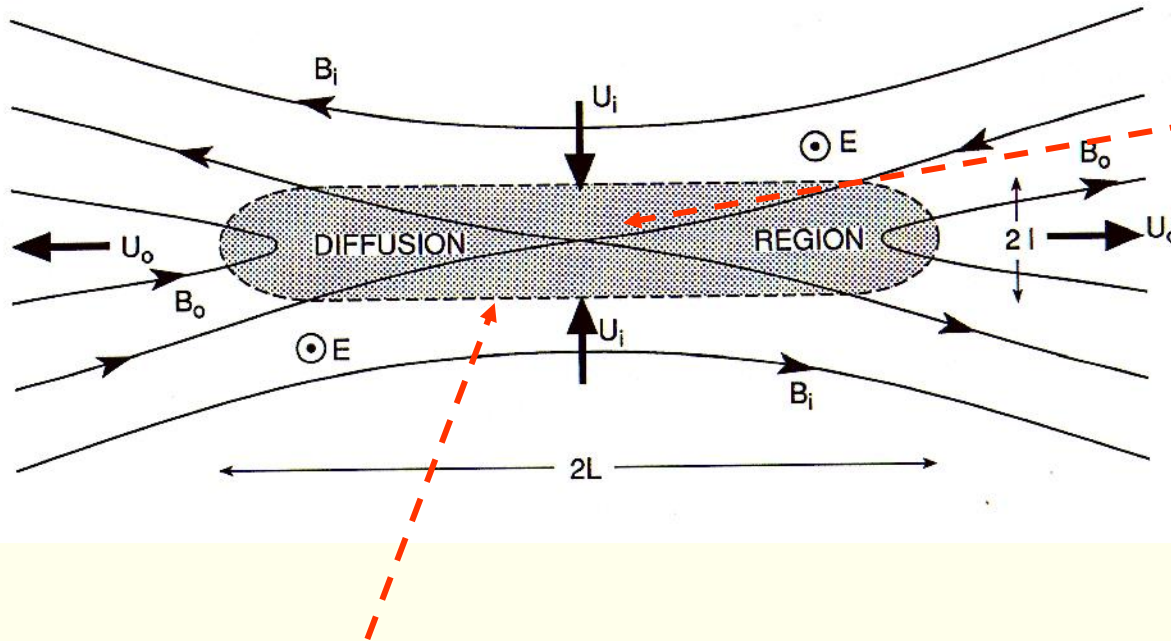
Reconnection

In 'diffusion region':

$$R_m = \mu_0 \sigma l v \sim 1$$

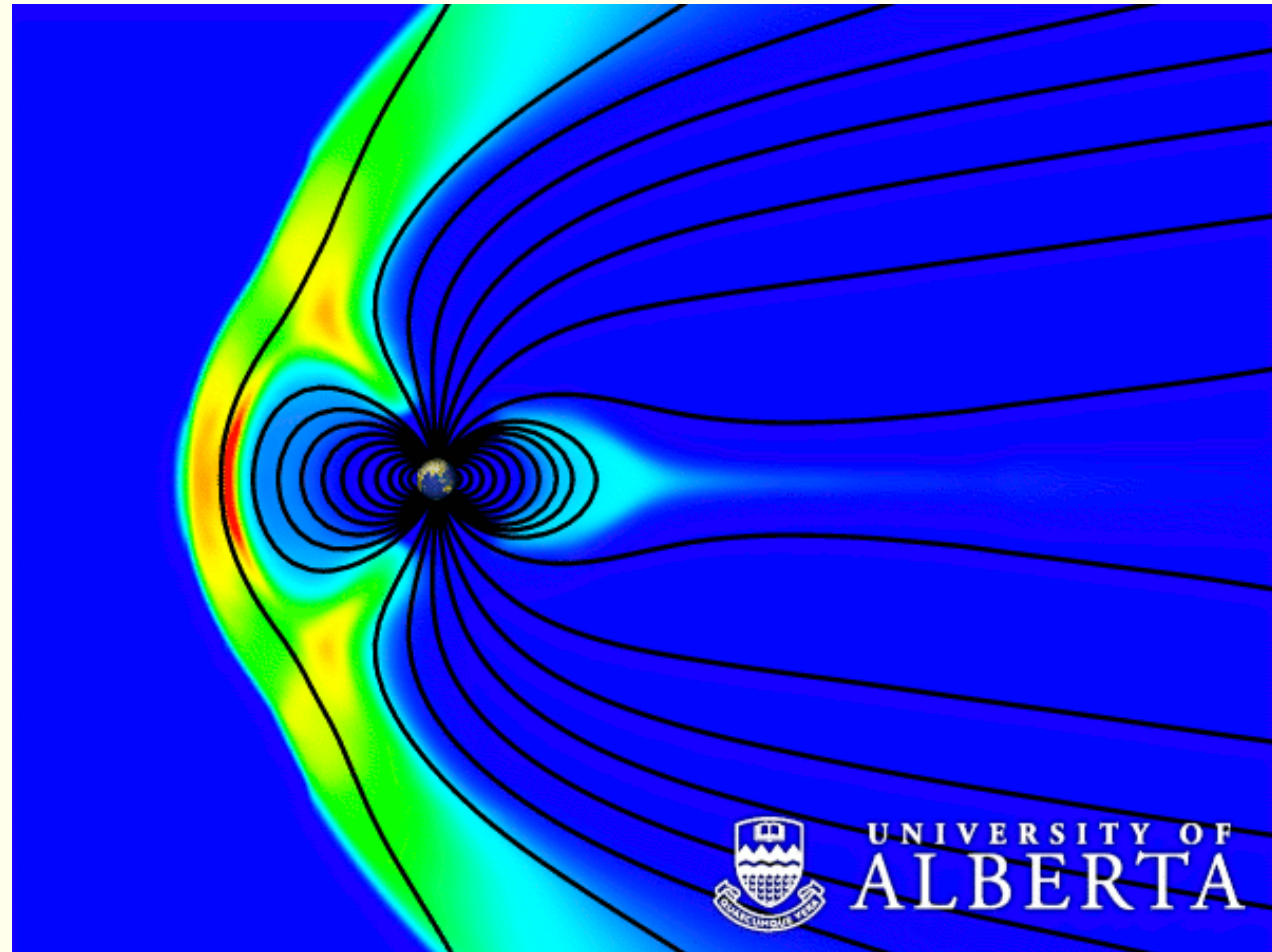
Thus: **condition** for frozen-in magnetic field breaks down.

A second **condition** is that there are two regions of magnetic field pointing in *opposite* direction:



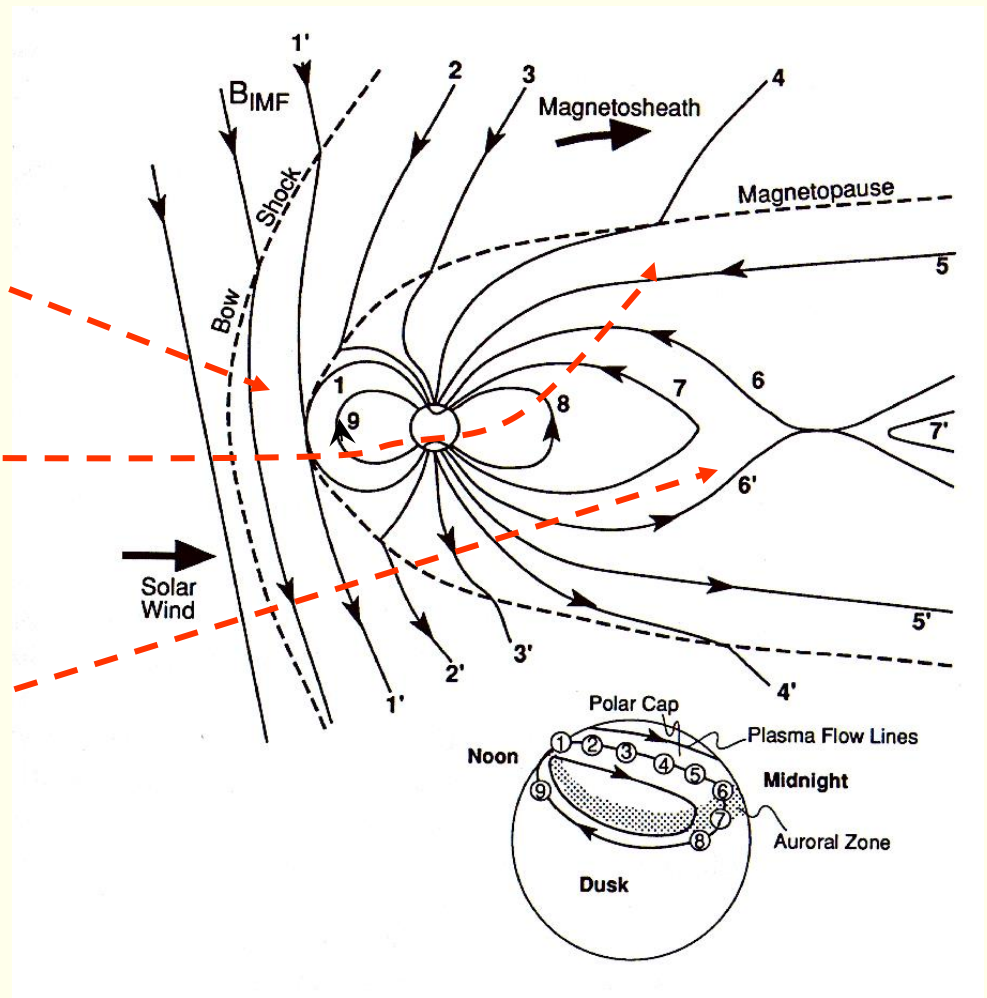
- Field lines are “cut” and can be re-connected to other field lines
- **Magnetic energy is transformed into kinetic energy ($U_o \gg U_i$)**
- **Plasma from different field lines can mix**

Reconnection and plasma convection

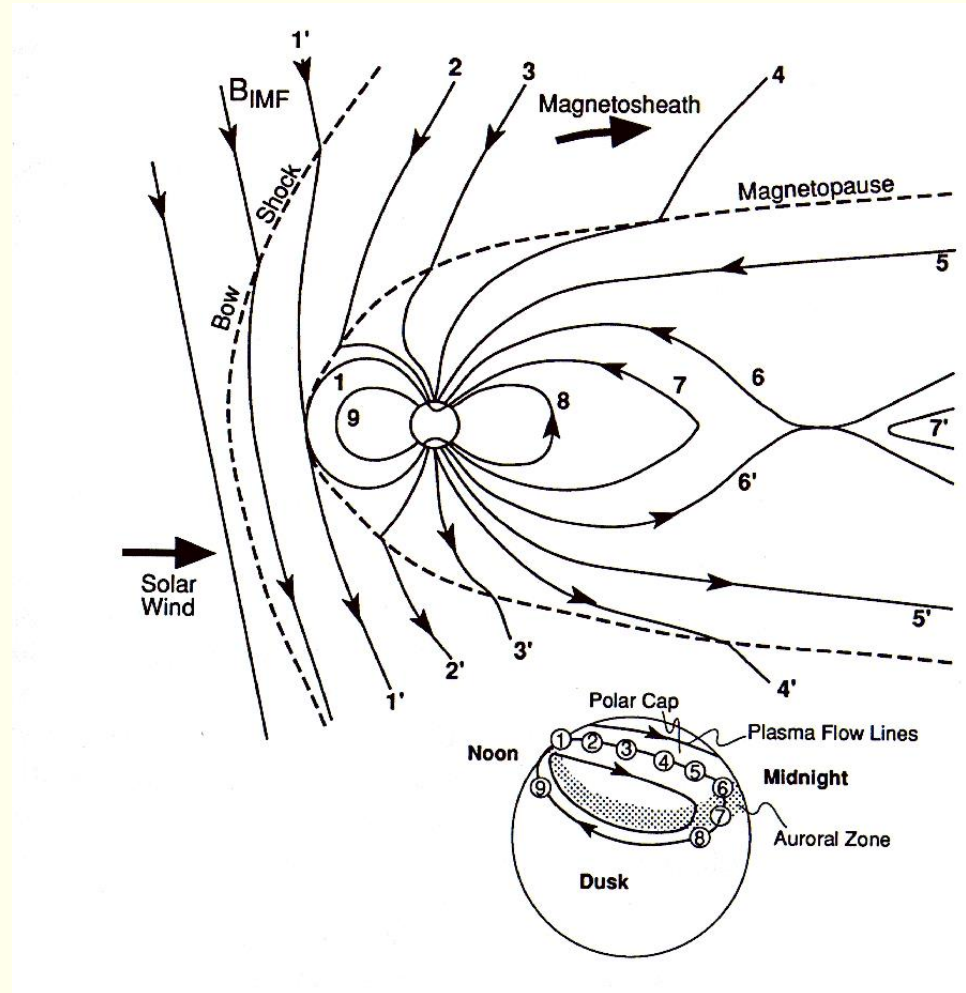


Reconnection och plasma convection

- Reconnection on the dayside “re-connects” the solar wind magnetic field and the geomagnetic field
- In this way the plasma convection in the outer magnetosphere is driven
- In the night side a second reconnection region drives the convection in the inner magnetosphere. The reconnection also heats the plasmashet plasma.

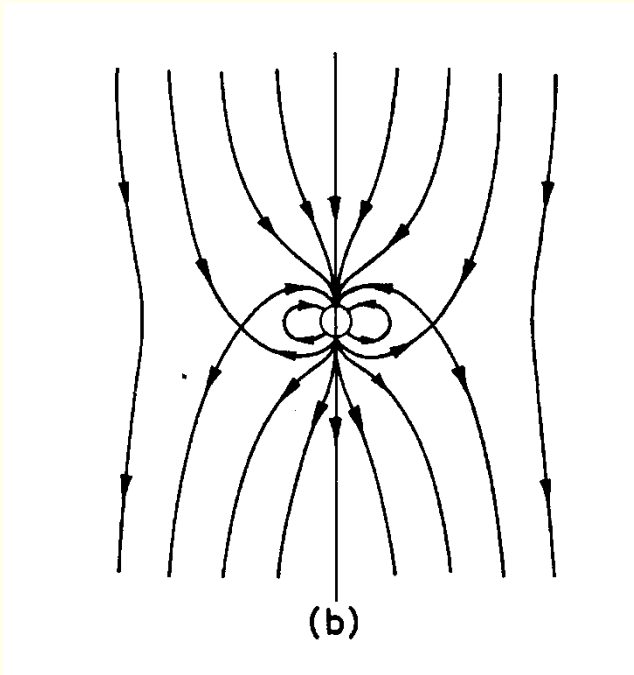


What happens if IMF is northward instead?

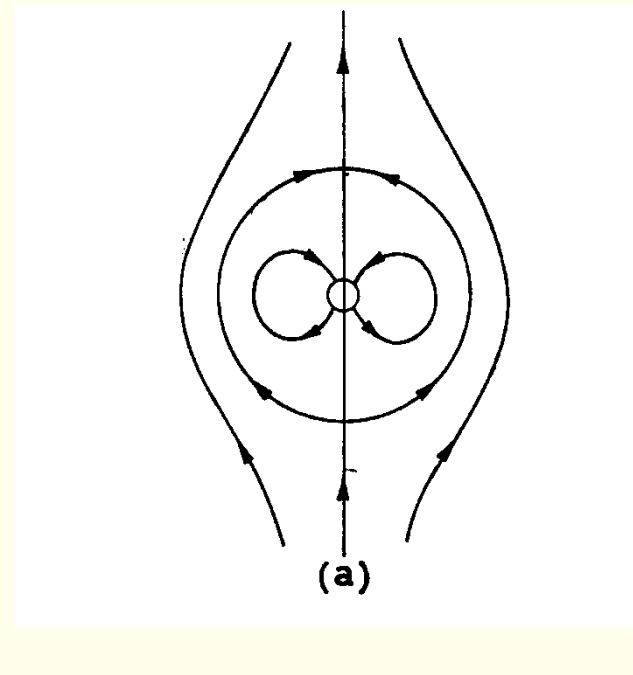


Magnetospheric dynamics

open magnetosphere



closed magnetosphere



southward 

**Interplanetary
magnetic field (IMF)**

 **northward**



What do the magnetospheres of the other planets look like?

Planetary magnetospheres

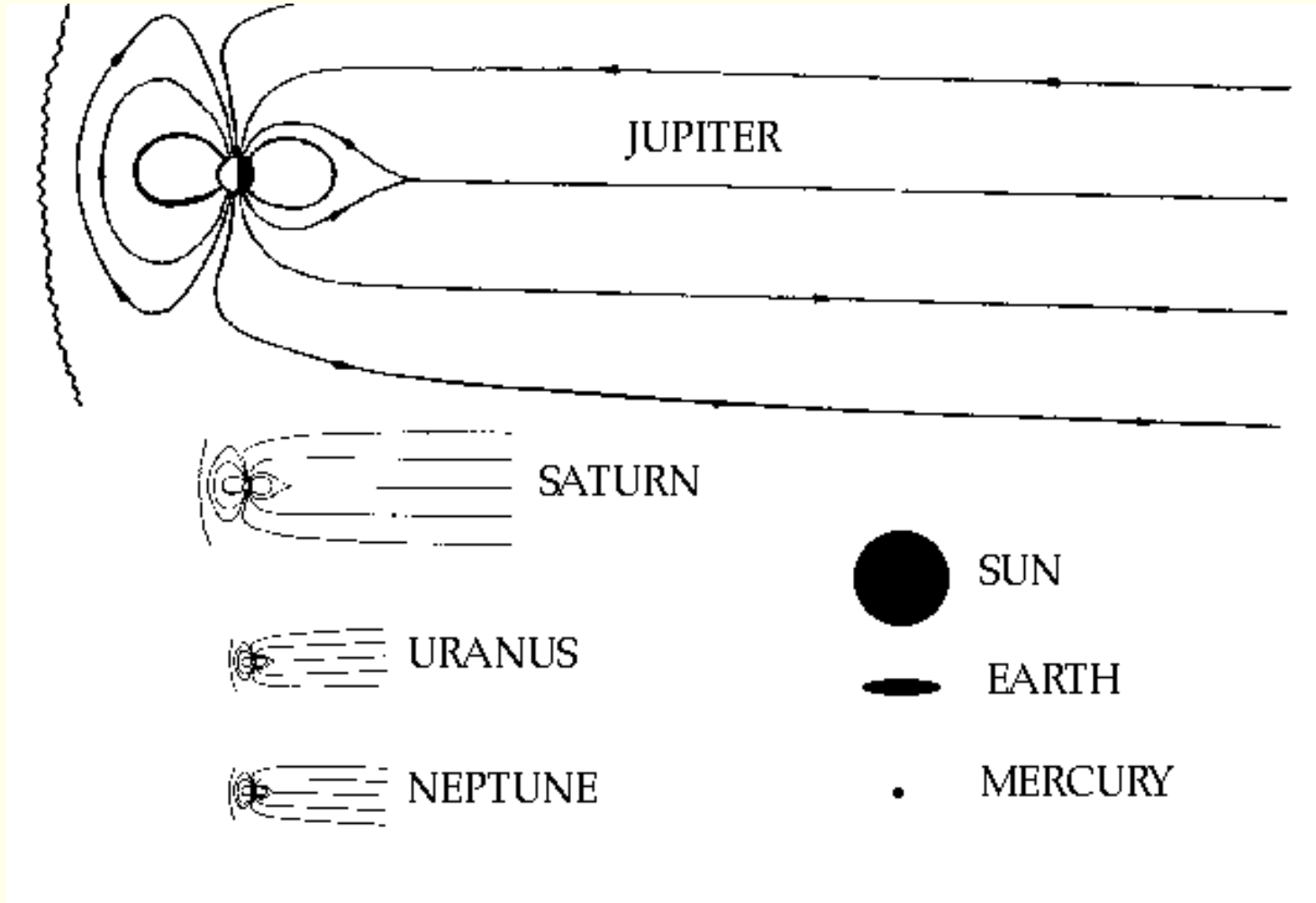
	Radius Earth radii	Spin period (days)	Equatorial field strength (μT)	Magnetic axis direction relative to spin axis	Polarity relative to Earth's	Typical magnetopause distance (planetary radii)
Mercury	0.38	58.6	0.35	10°	Same	1.1
Venus	0.95	243	< 0.03	-	-	1.1
Earth	1.0	1	31	11.5°	Same	10
Mars	0.53	1.02	0.065	-	Opposite	?
Jupiter	11.18	0.41	410	10°	Opposite	60-100
Saturn	9.42	0.44	40	$<1^\circ$	Opposite	20-25
Uranus	3.84	0.72	23	60°	Opposite	18-25
Neptune	3.93	0.74	20-150 ^{*)}	47°	Opposite	26 ^{**)}

*) The magnetic field differs greatly from a dipole field. The numbers represent maximum and minimum strength at the planetary surface

***) Based on single passage

Very weak magnetic fields

Relative size of the magnetospheres



Comparative magnetospheres

In situ observations

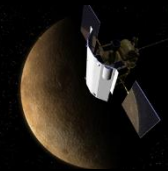
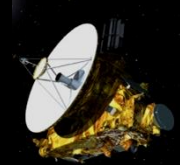
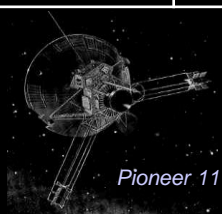
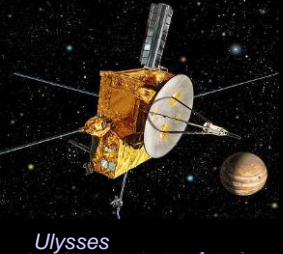
Space probe

Celestial body

Observations

<i>Mariner 10</i>	Mercury	1974 – 1975
<i>Messenger</i> *	Mercury	2008 – 2015
<i>Pioneer 10,11</i>	Jupiter, Saturn	1973 – 1979
<i>Voyager 1,2</i>	Jupiter, Saturn, Uranus, Neptune	1977 – 1989
<i>Ulysses</i>	Jupiter	1992
<i>Galileo</i> *	Jupiter	1995 – 2003
<i>Cassini</i> *	Jupiter, Saturn	2004 –
<i>New Horizons</i>	Jupiter	2007
<i>Rosetta</i>	Churyumov-Gerasimenko	2014 - 2016

* *Orbiters*



Mercury

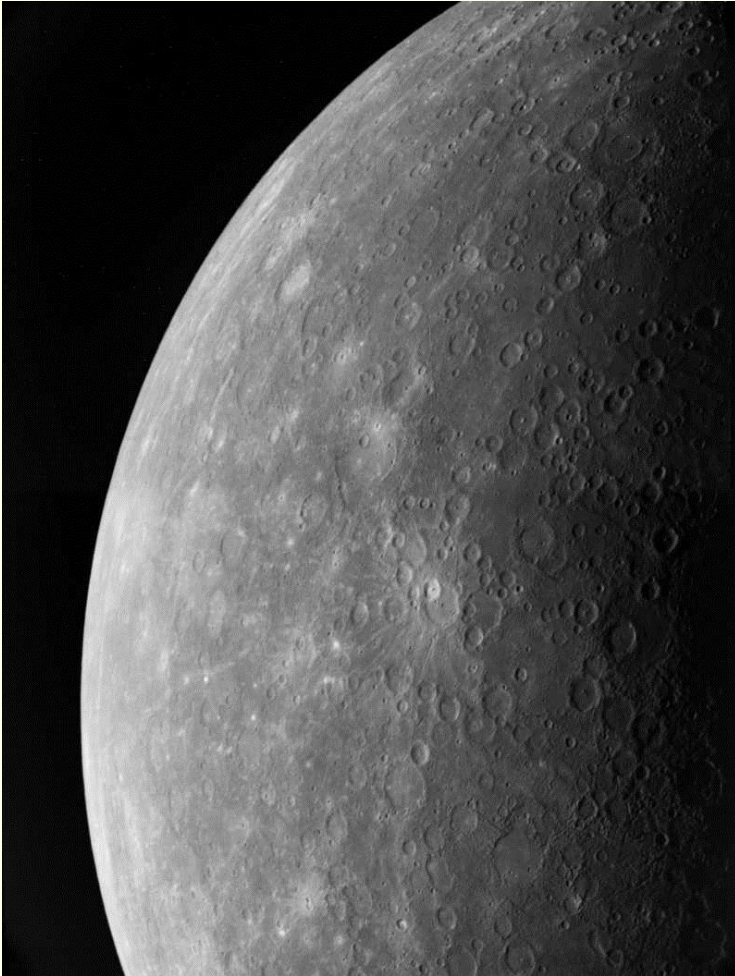
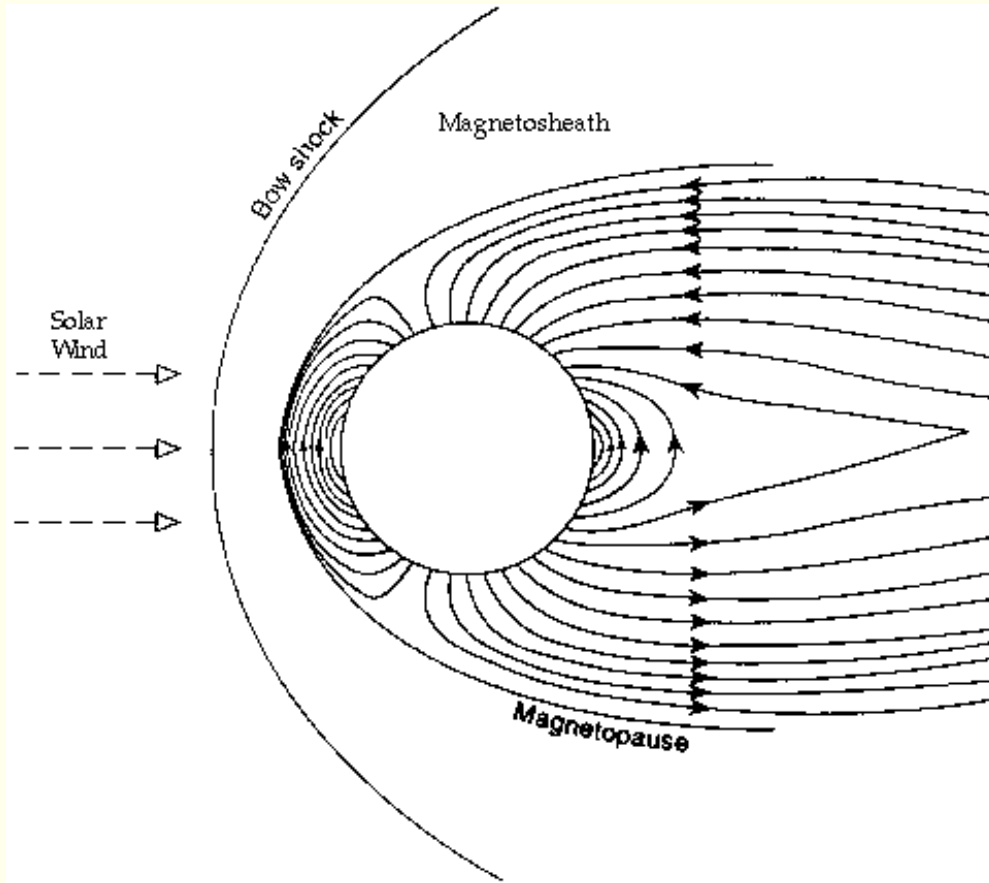


Photo from Mariner 10

- $r_M = 0,38 r_E$
- $m_M = 0,06 m_E$
- distance from sun: 0,39 AU
- no or extremely thin atmosphere

Mercury's magnetosphere



- rather large magnetic field
- the high solar wind density close to the sun makes the magnetosphere very small
- no ionosphere

Venus

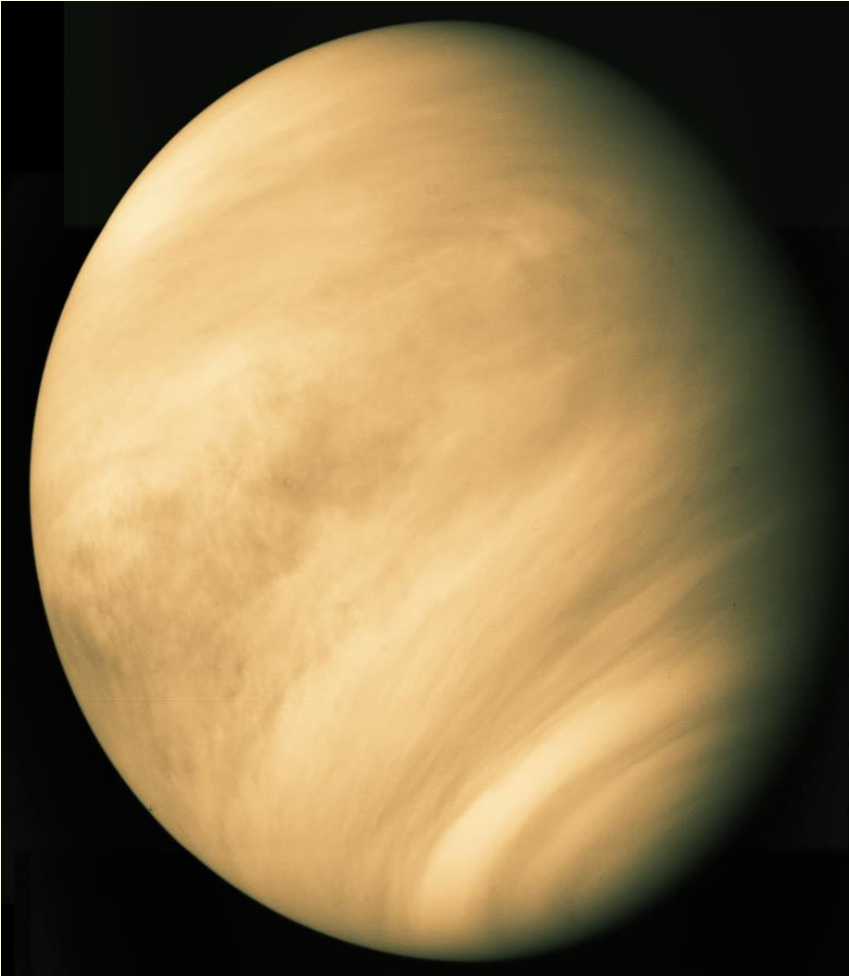
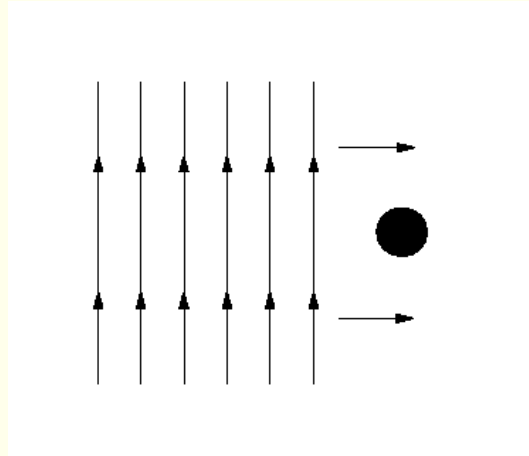


Photo from Galileo

- $r_V = 0,95 r_E$
- $m_V = 0,82 m_E$
- distance from sun : 0,72 AU
- very dense atmosphere
 - ~ 90 atm
 - 96% CO_2
- very weak magnetic field

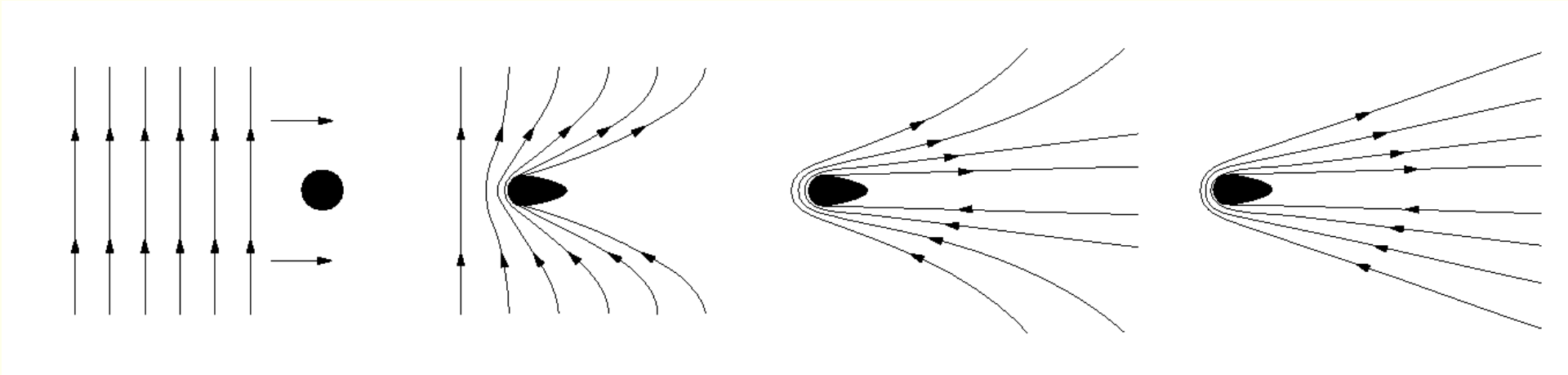
Comets, induced magnetotail



- The *coma* of the comet is ionized when the comet gets close to sun

What will happen when the interplanetary magnetic field (which is frozen into the solar wind) hits the coma?

Comets, induced magnetotail



- The **coma** of the comet is ionized when the comet gets close to sun
- This plasma stops the field lines from passing the comet, due to the frozen-in magnetic field
- Venus and Mars have similar induced, weak "magnetospheres"



Photo from Hubble Space Telescope

Mars

- $r_M = 0,53 r_E$
- $m_M = 0,11 m_E$
- distance from sun : 1,52 AU
- very thin atmosphere
 - ~ 0.01 atm.
 - 95% CO_2
- very weak magnetic field

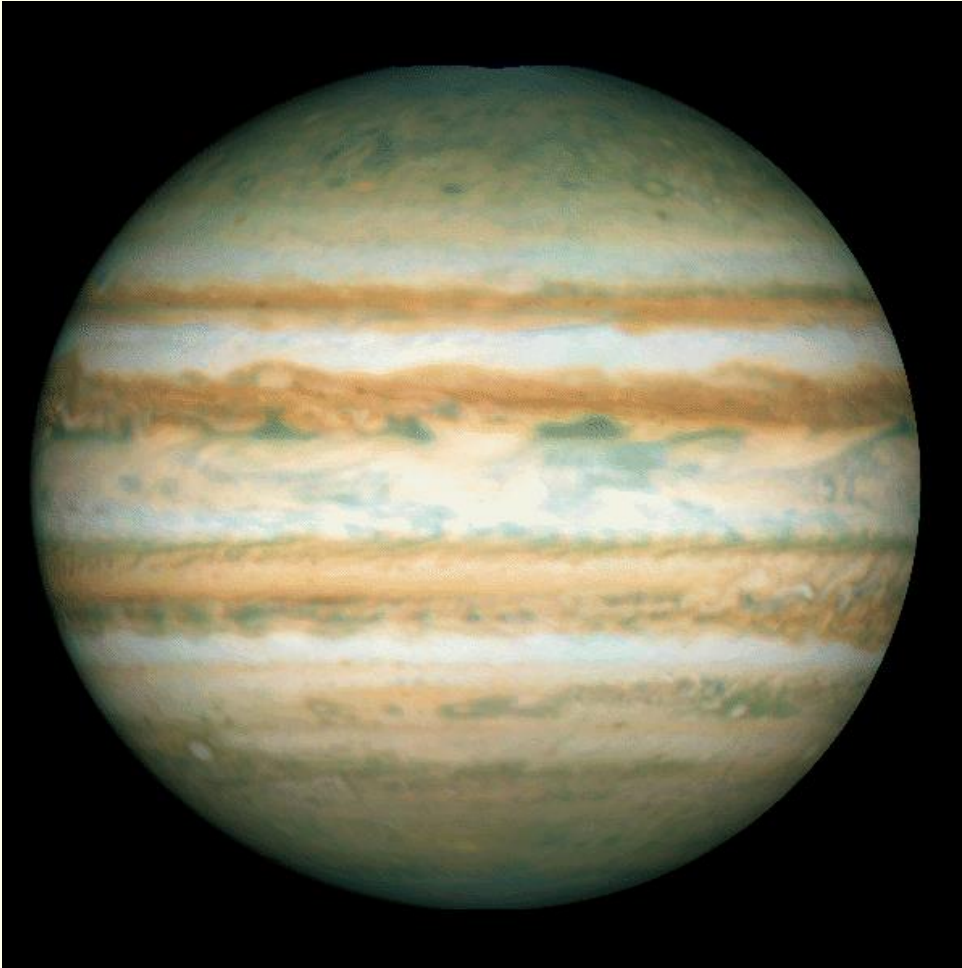
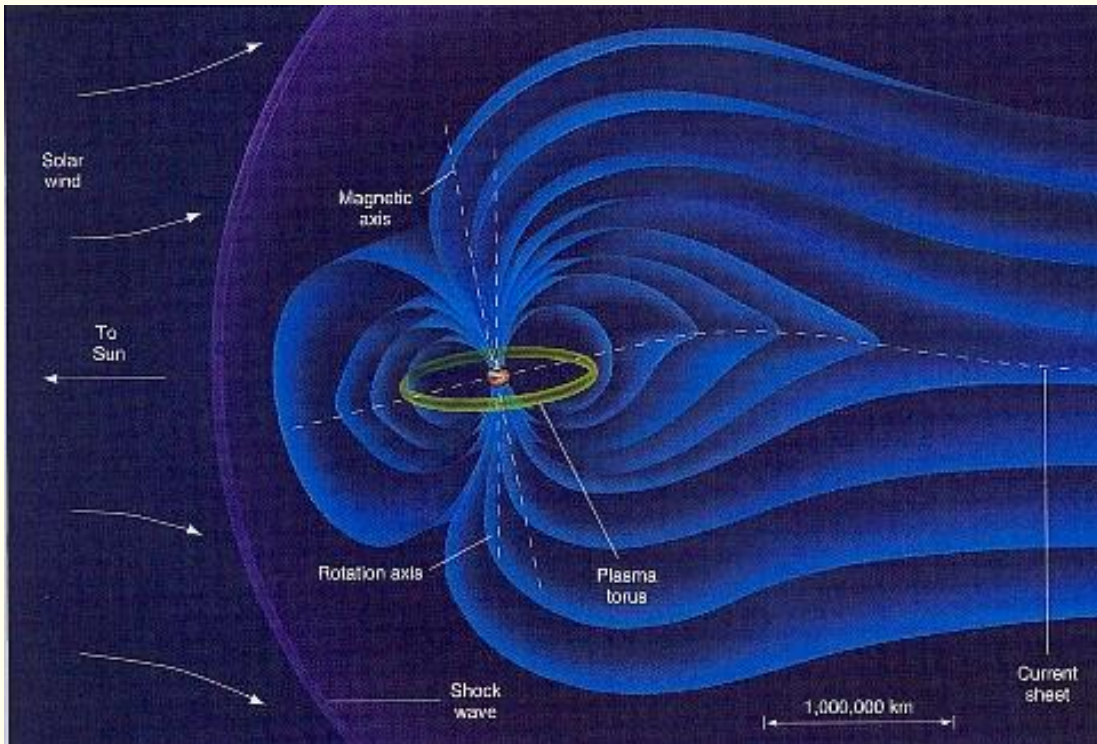


Photo from Hubble Space Telescope

Jupiter

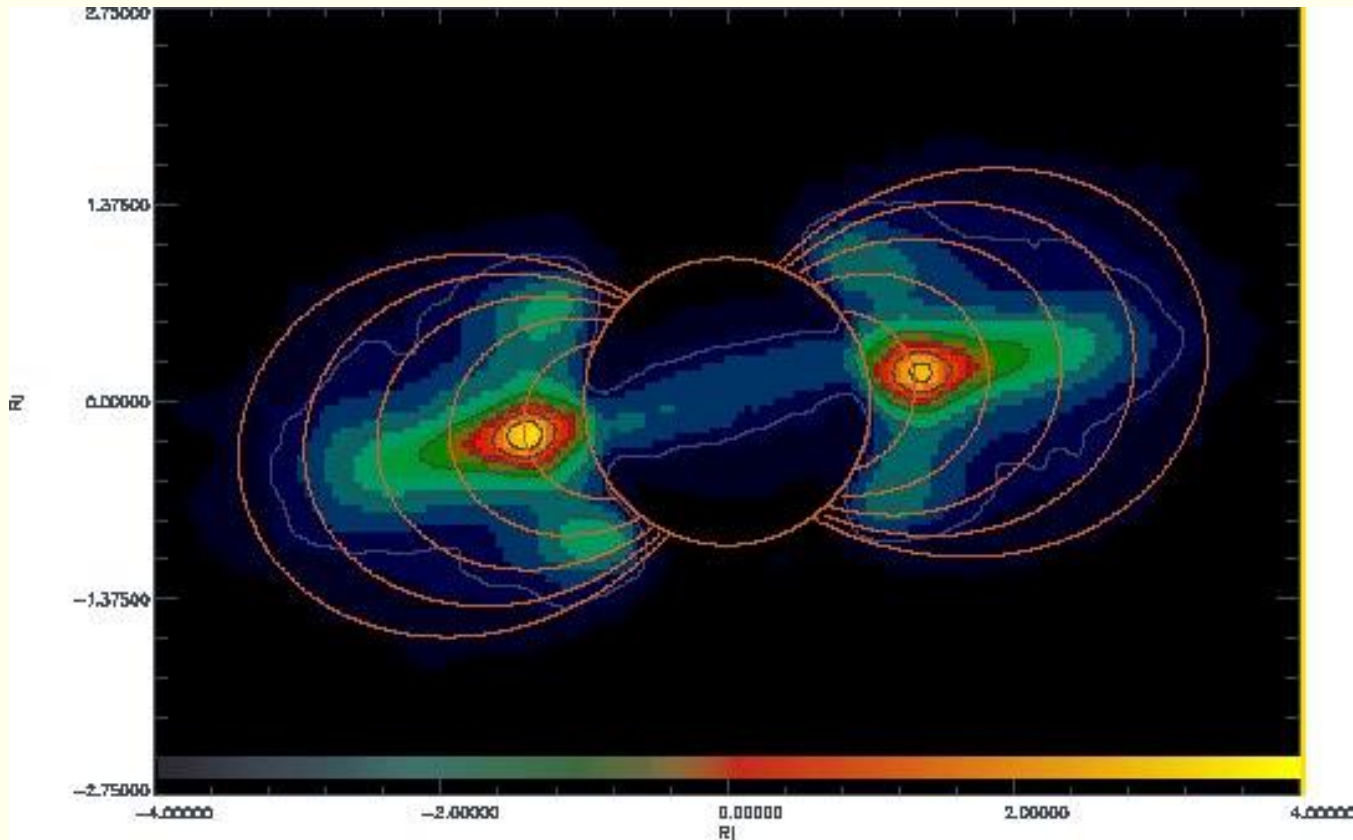
- $r_J = 11,2 r_E$
- $m_J = 318 m_E$
- distance from sun: 5,20 AU
- gas giant with very dense atmosphere containing hydrogen, helium, ammonium, methane, etc.
- ~ 60 moons (+ weak ring system)

Jupiter's magnetosphere



- high plasma density, Io is an important source
- plasma pressure thus becomes an important factor for balancing the solar wind pressure
- the plasma co-rotates with Jupiter, which gives the magnetosphere a flattened look

Synchrotron radiation from Jupiter's radiation belts



- Gyrating electrons emit "synchrotron radiation" with frequencies $\sim f_{ce} = eB/(2\pi m_e)$
- The emitted power is proportional to the electron temperature:
$$P = CT_e$$
- In this way you can get a picture of the radiation belts

Galilean satellites



Volcanic activity, source for plasma.

Oceans under the ice?

Has its own magnetosphere the size of Mercury's

Weak magnetic field

Saturn

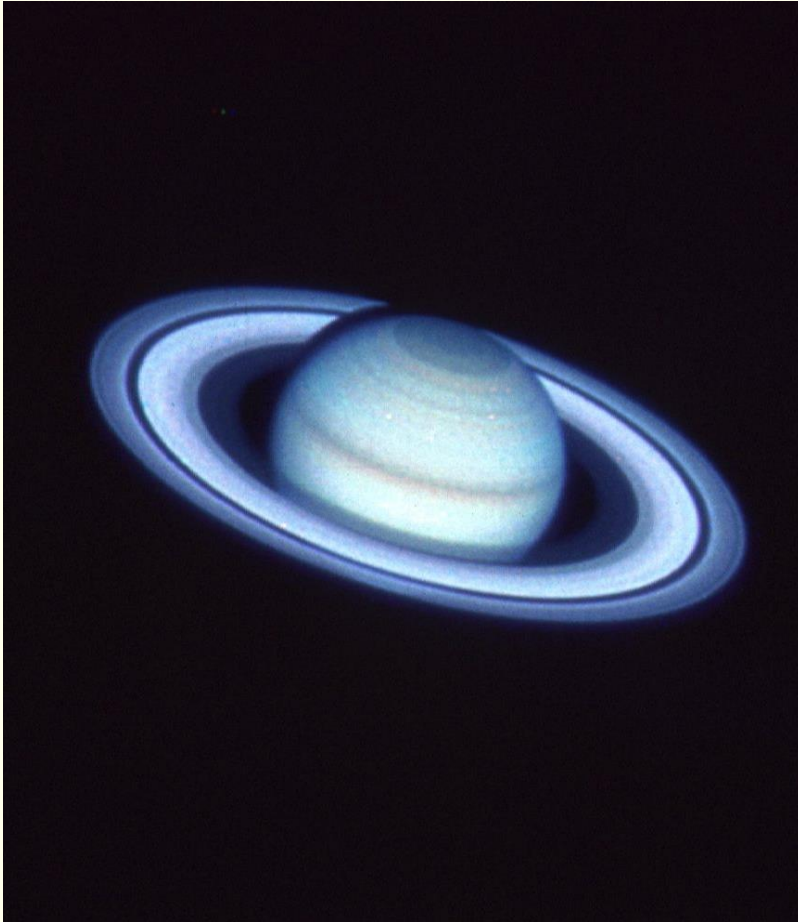
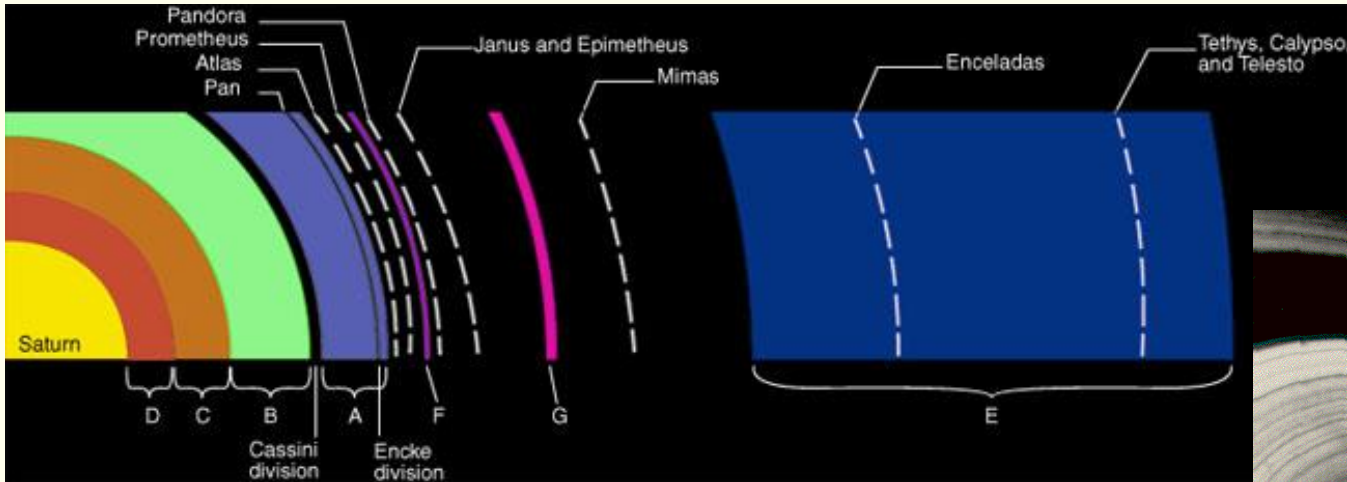


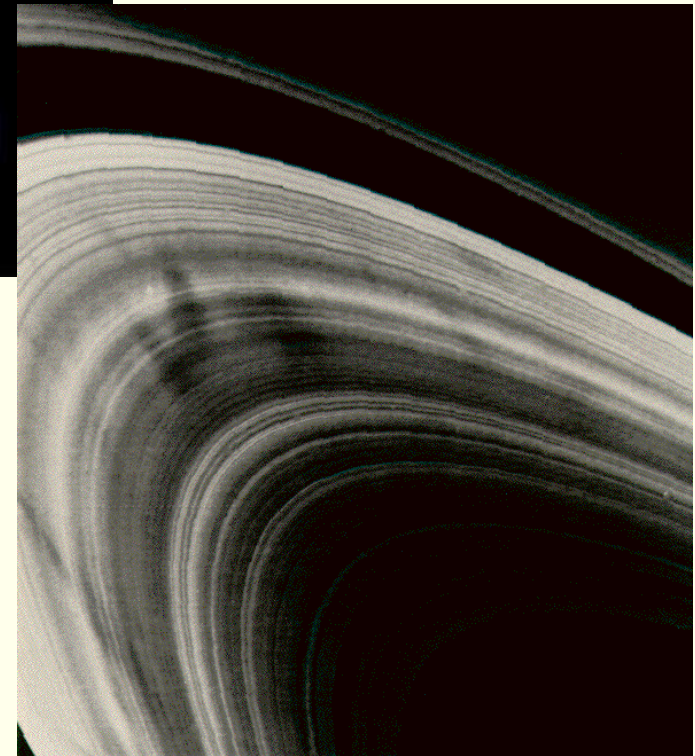
Photo from Hubble Space Telescope

- $r_s = 9,42 r_E$
- $m_s = 95 m_E$
- distance from sun: 9.53 AU
- gas giant with very dense atmosphere containing sulphur, hydrogen, helium, ammonium, methane, etc.
- ~ 31 moons + ring system

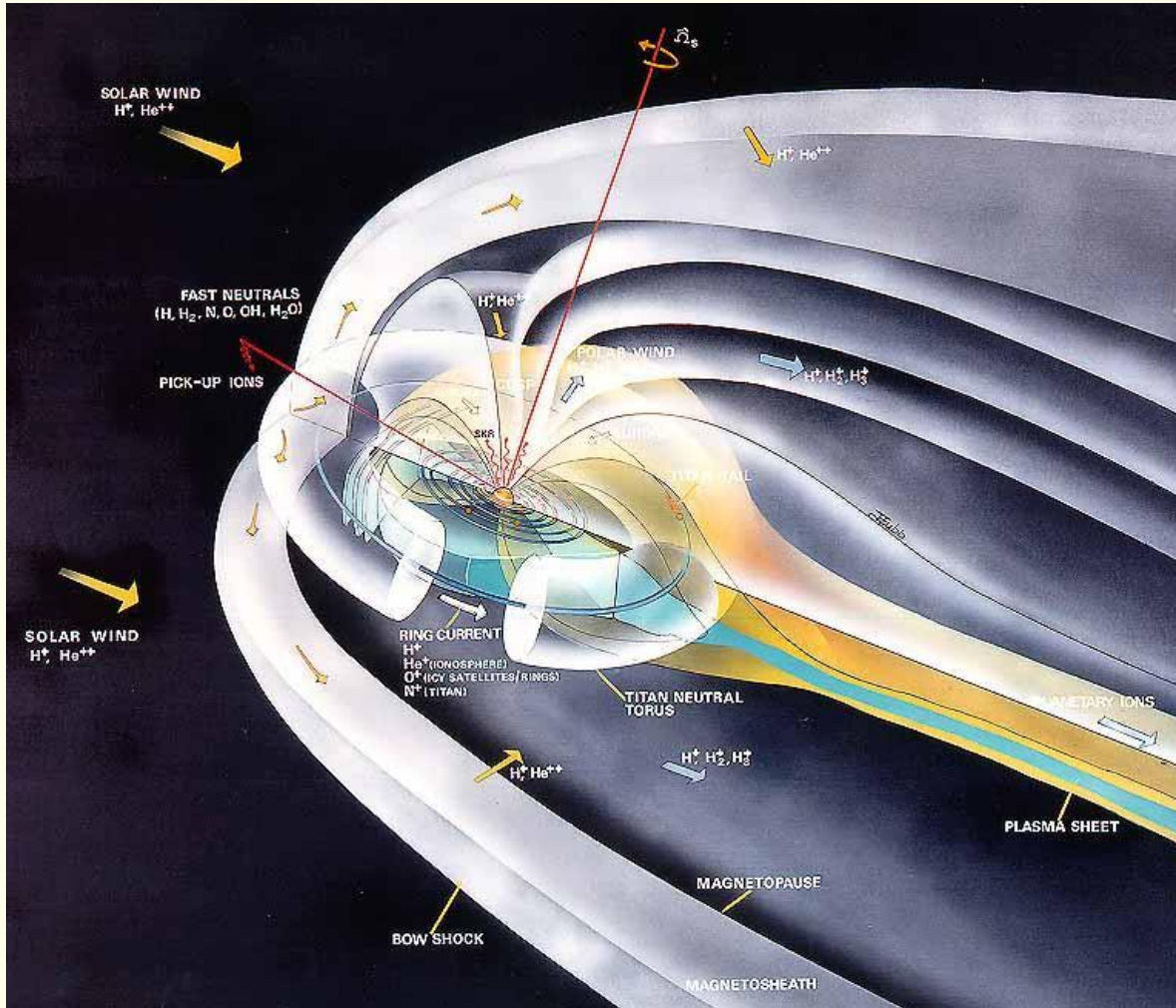
Saturn's rings



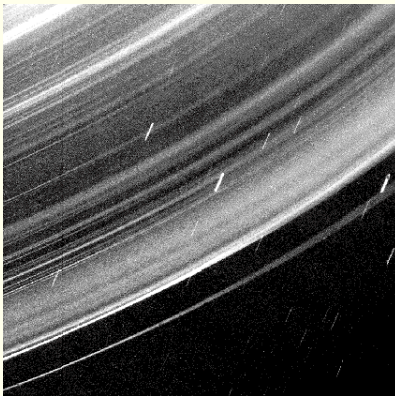
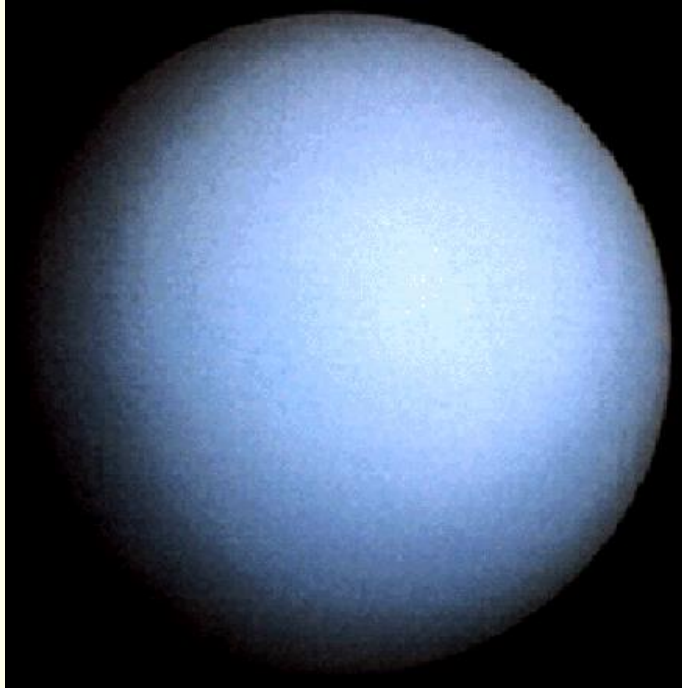
- ring system is made up of ice and mineral particles from ~ 1 cm to ~ 1 km
- rings are only 1.5 km thick



Saturn's magnetosphere



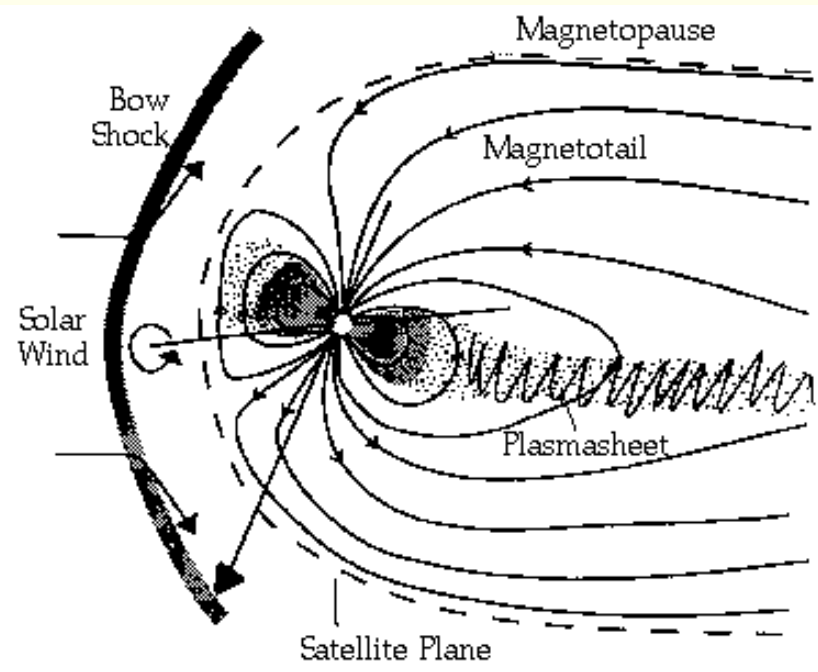
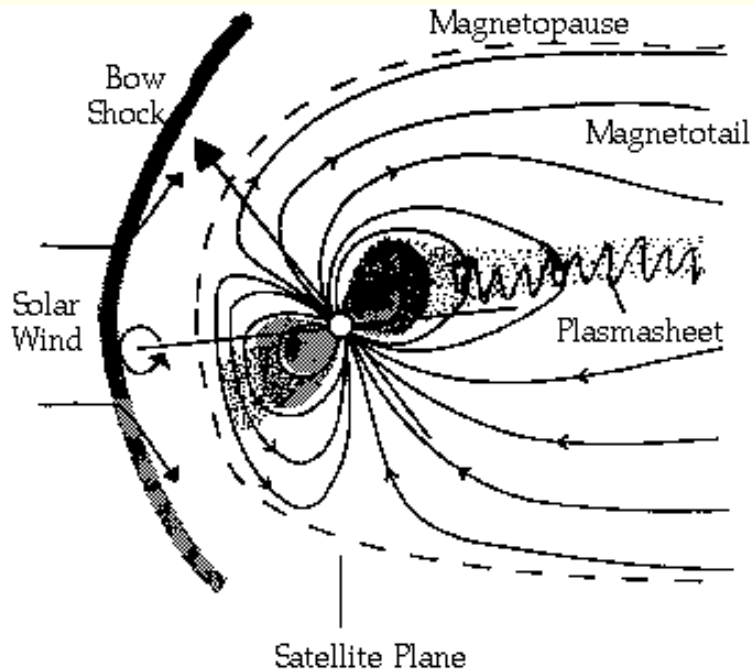
- ring system is both source and sink for plasma and limits the size of the plasmasphere
- plasma pressure important for balancing solar wind pressure, just as for Jupiter
- Intense radiation belt with 50 MeV protons



Uranus

- $r_U = 3,84 r_E$
- $m_U = 14,5 m_E$
- distance from sun: 19,2 AU
- gas giant with very dense atmosphere containing mainly methane
- 20 moons + ring system

Uranus



- Uranus' rotational axis almost in ecliptic plane
- magnetic field axis makes an angle of around 60° with rotational axis
- this gives enormous daily variations of the structure of the magnetosphere

Neptune

- $r_N = 3,93 r_E$
- $m_N = 17,2 m_E$
- distance from sun : 30.1 AU
- gas giant with very dense atmosphere containing mainly methane
- 11 moons + ring system

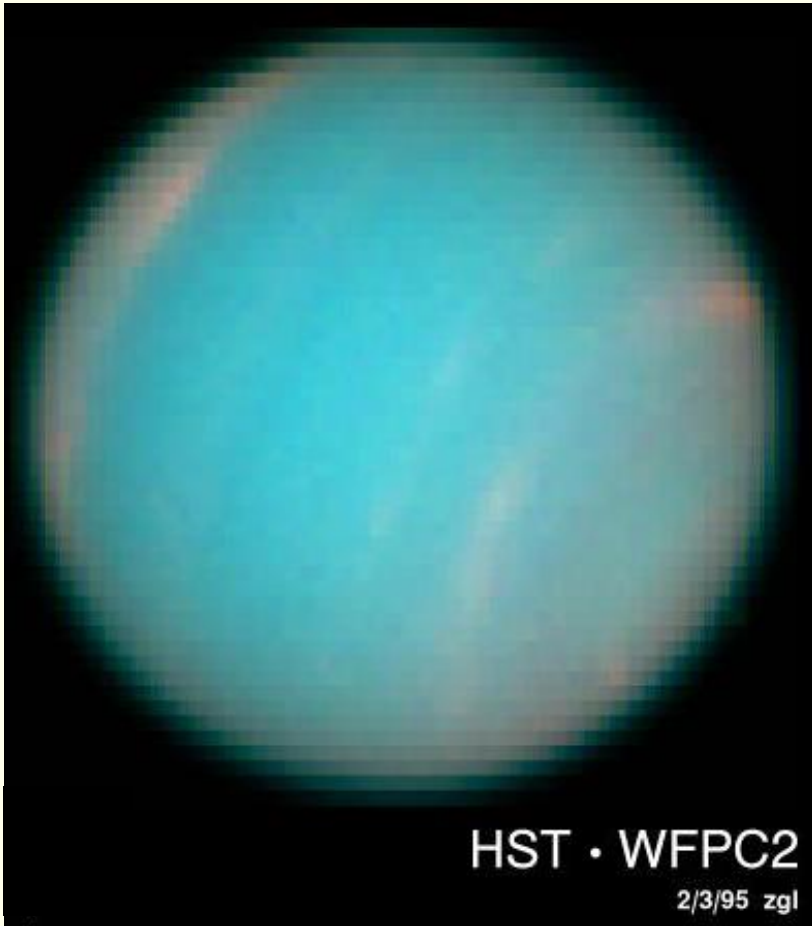
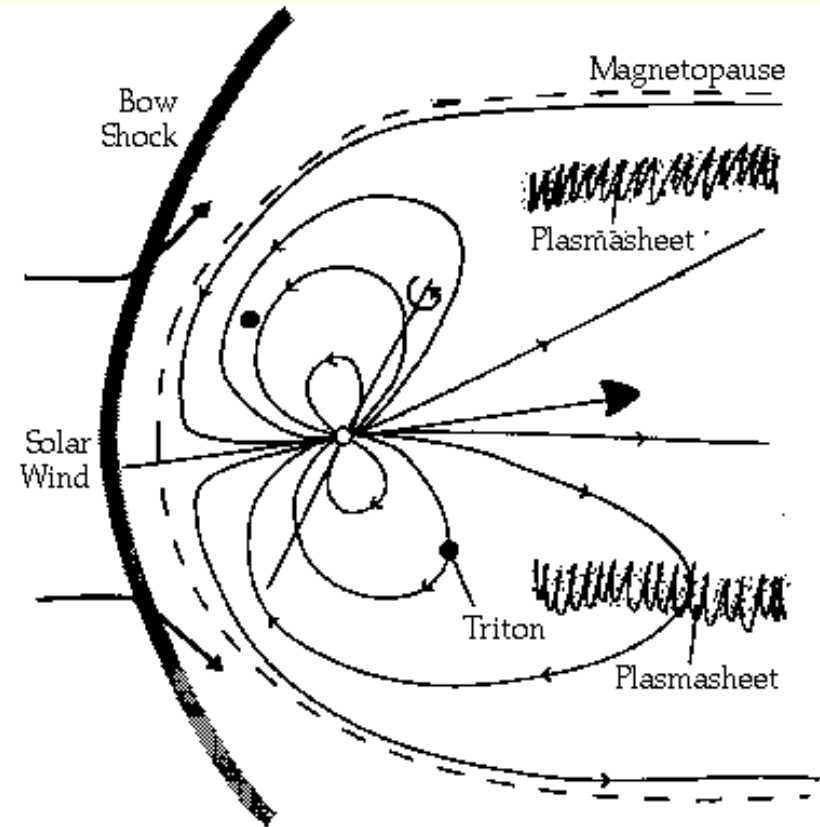
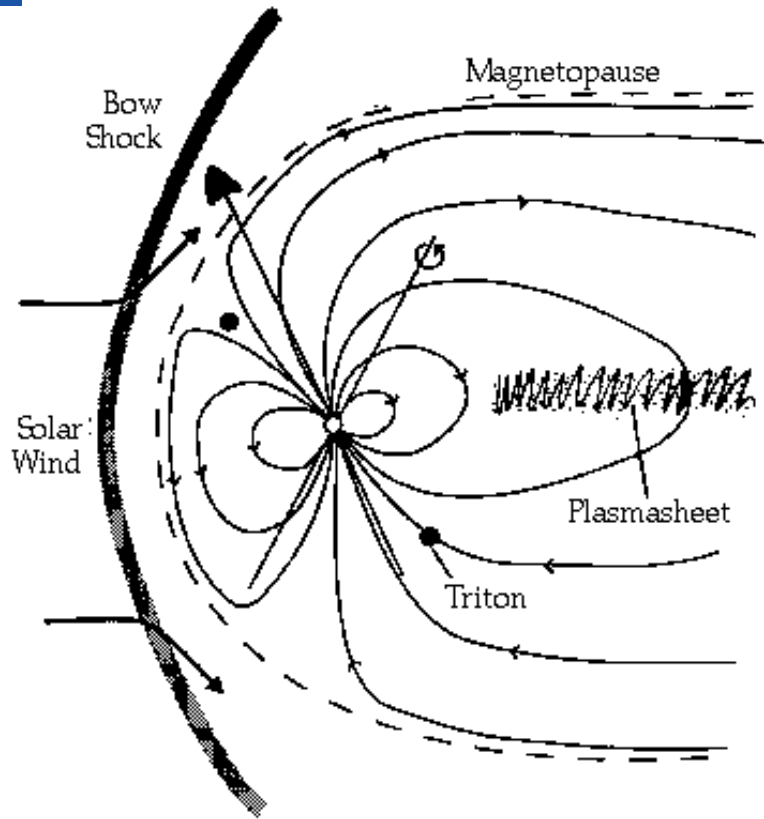


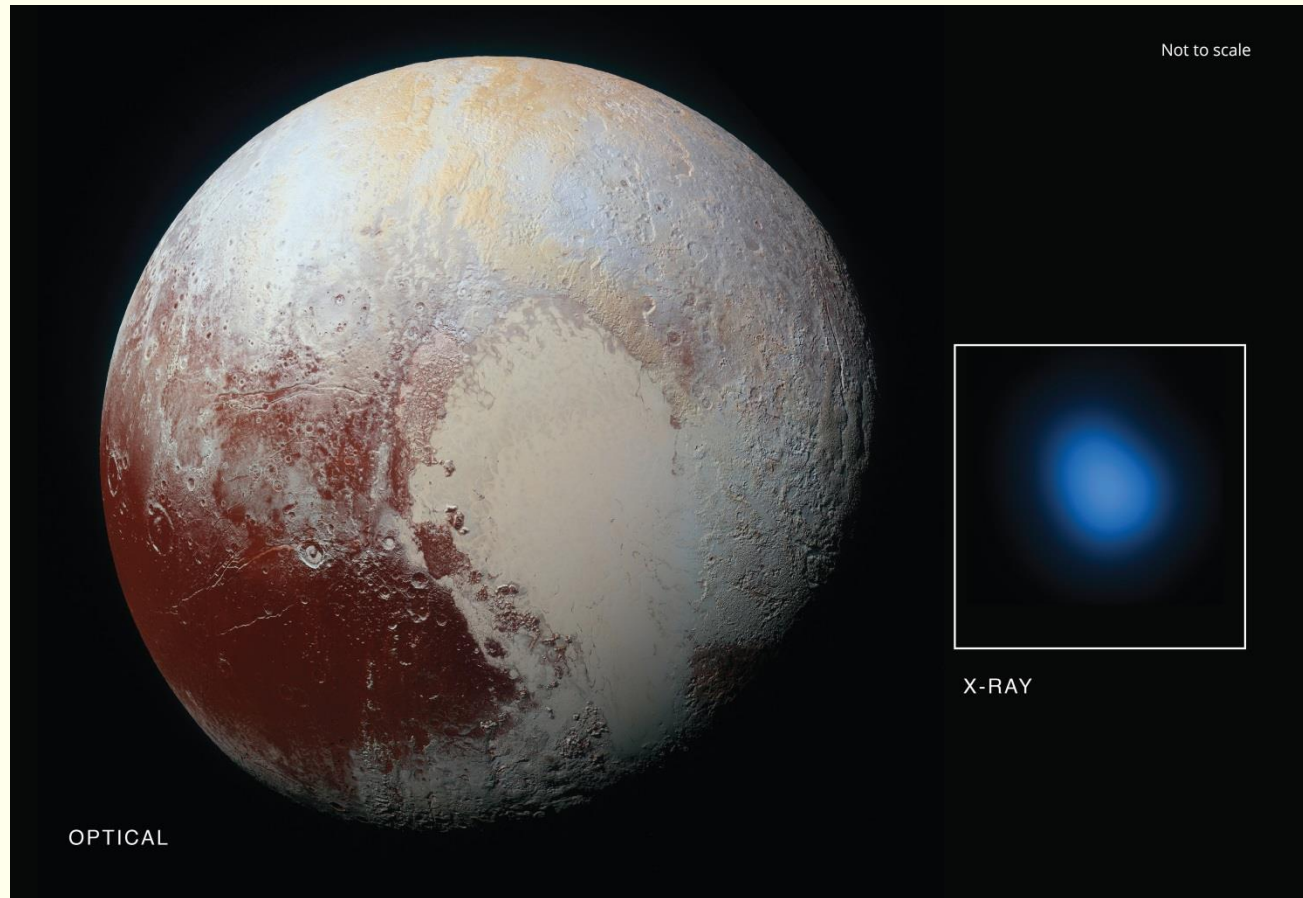
Photo from Hubble Space Telescope

Neptune



- magnetic field axis makes an angle of around 43° with rotational axis
- this gives enormous daily variations of the structure of the magnetosphere also for Neptune

Pluto

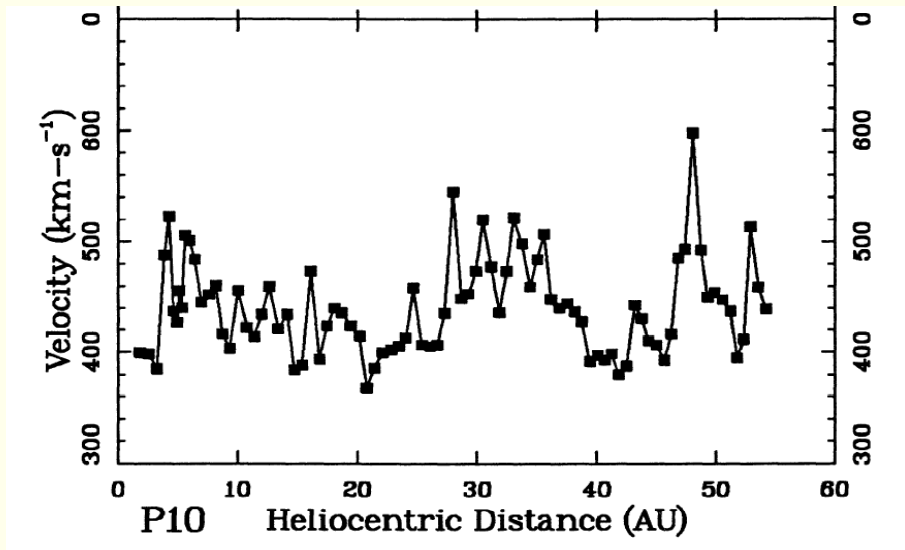


Scientists using NASA's Chandra X-ray Observatory have made the first detections of X-rays from Pluto. These observations offer new insight into the space environment surrounding the largest and best-known object in the solar system's outermost regions.

Comparative magnetospheres

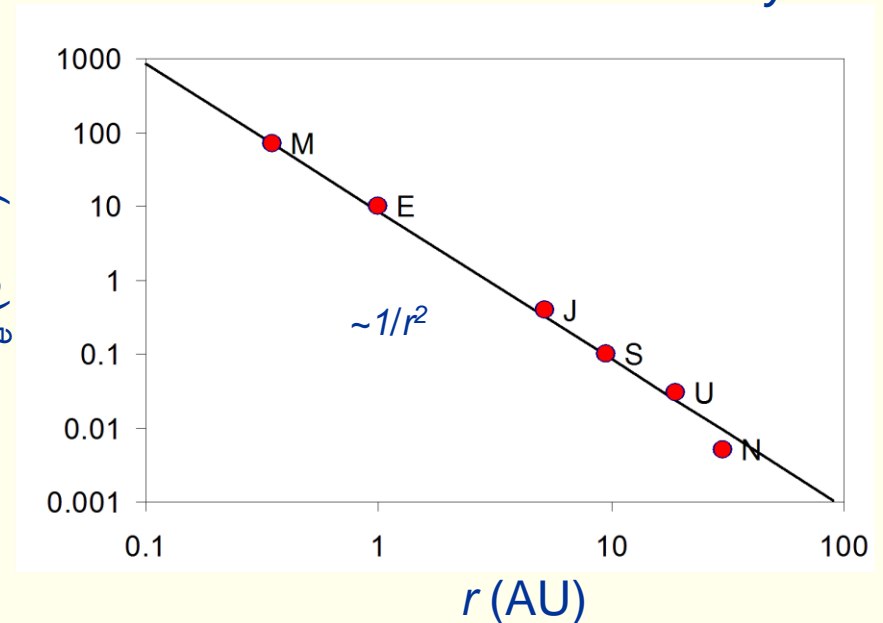
Solar wind properties

Solar wind velocity

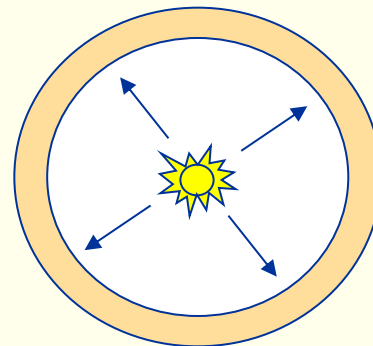


*Pioneer 10, measurements
[Grazin et al., 1994]*

Solar wind electron density



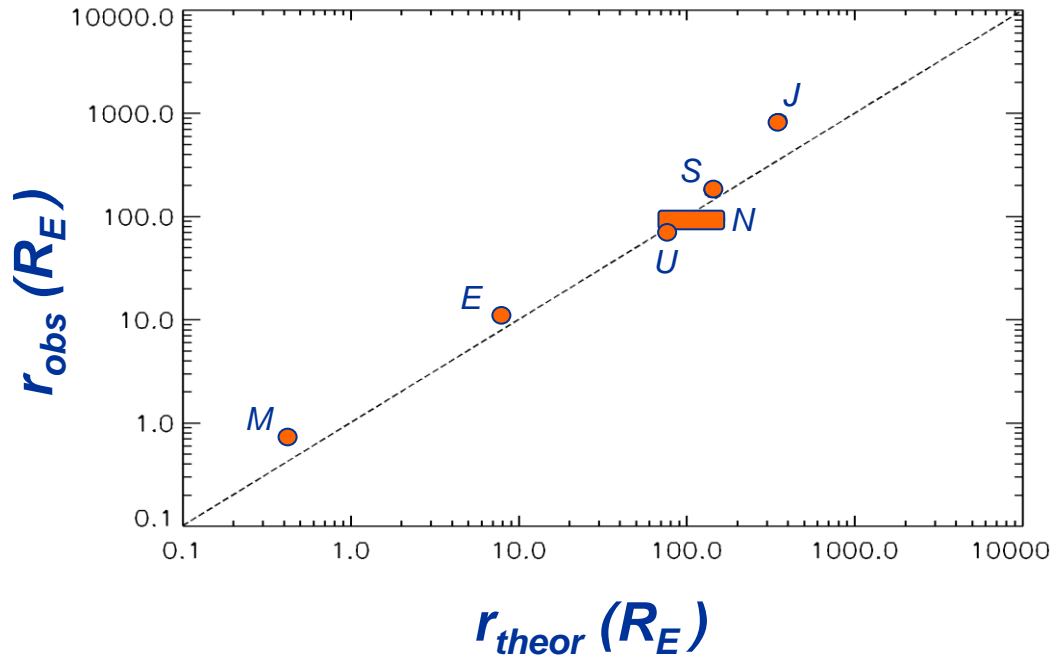
[Blanc et al., 2005]



$$dV = 4\pi r^2 dr$$

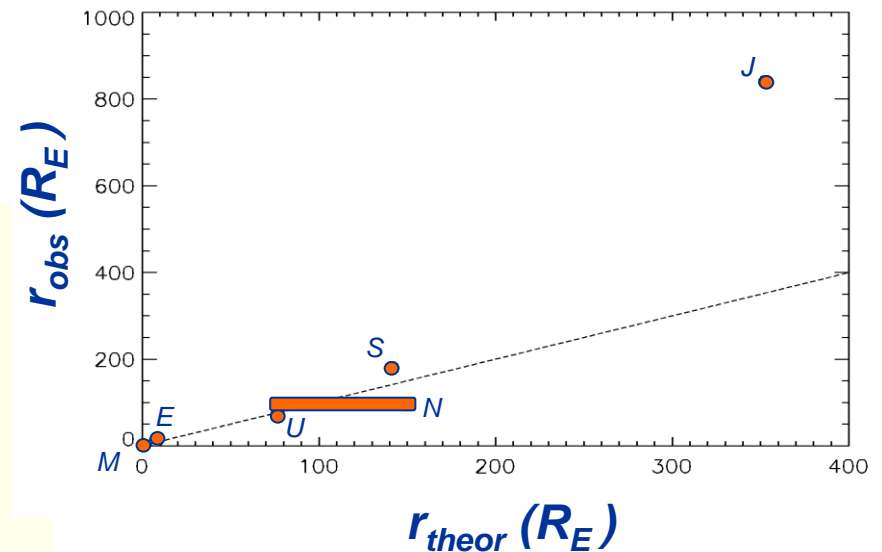
Comparative magnetospheres

Observed vs. theoretical standoff-distance



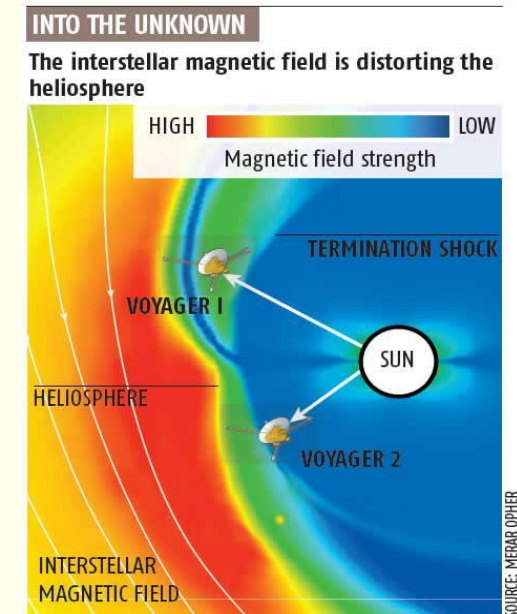
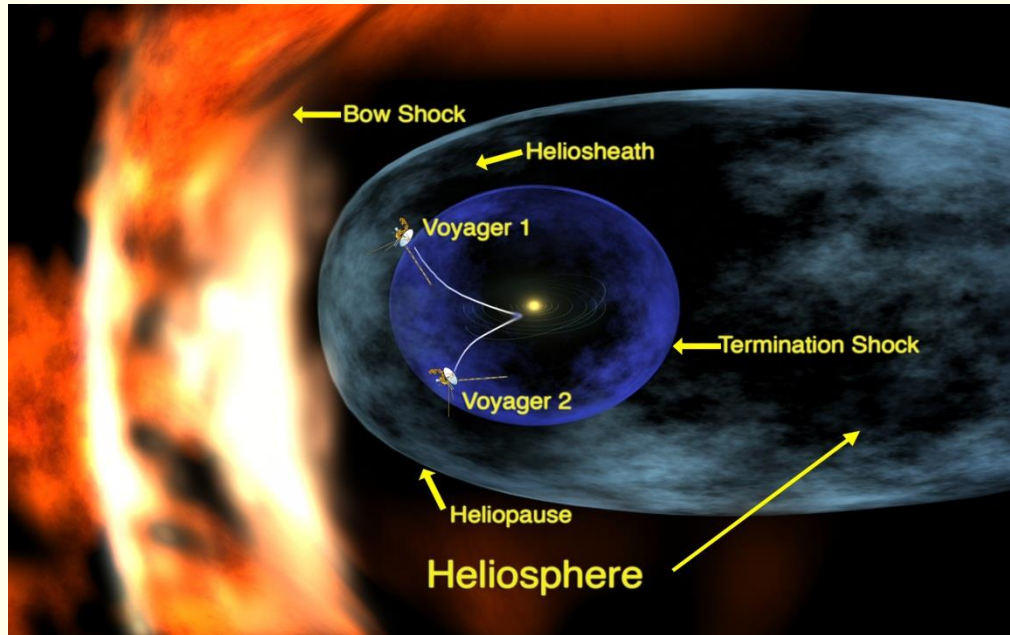
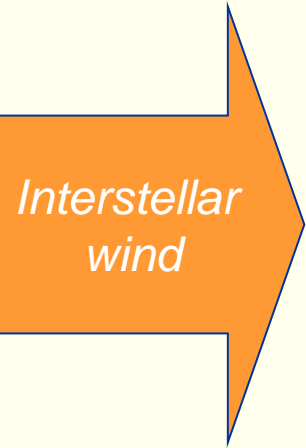
$$r_{theor} = \left(\frac{\mu_0 a}{4\pi} \right)^{1/3} \left(2\mu_0 \rho_{SW} v_{SW}^2 \right)^{-1/6}$$

- Model reasonably valid over three orders of magnitude
- Size of Jupiter's (and maybe Saturn's) magnetosphere underestimated



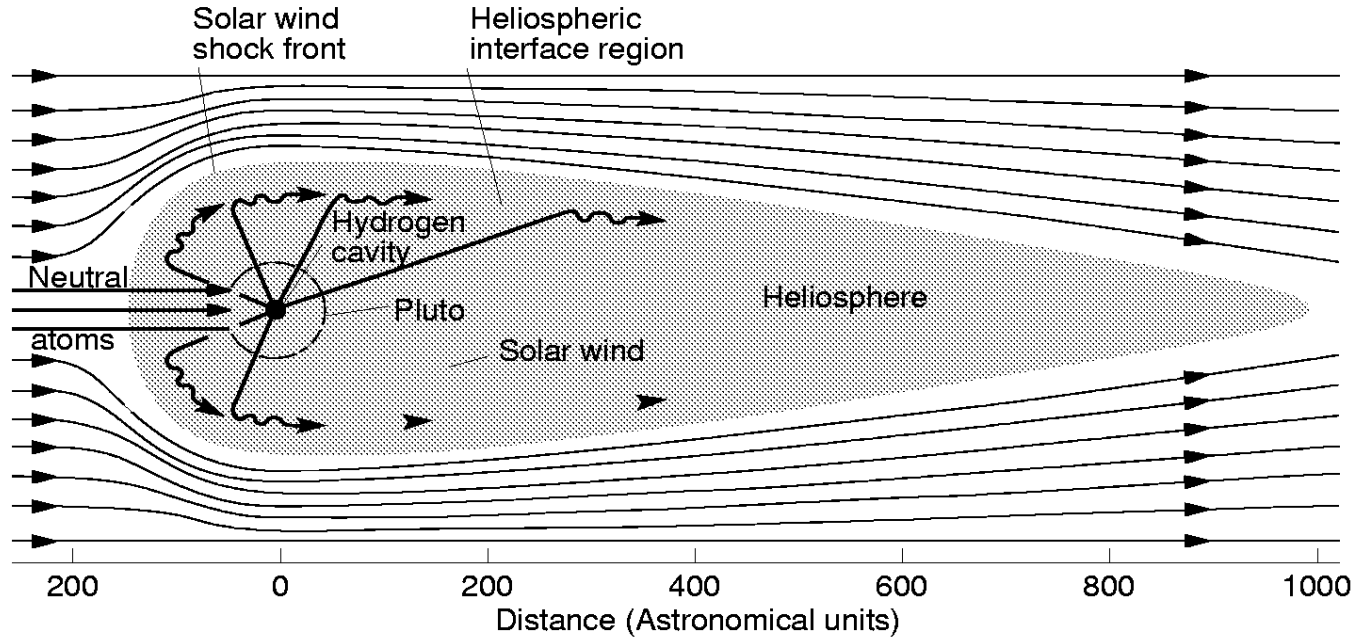
Other other magnetospheres

Heliosphere



[Opher, 2007]

Heliosphere



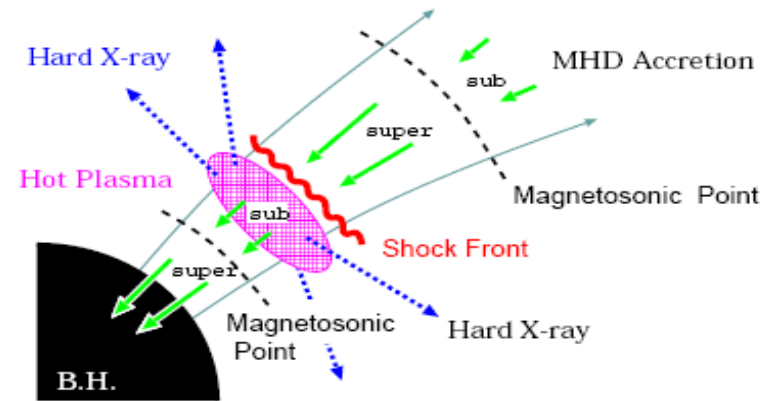
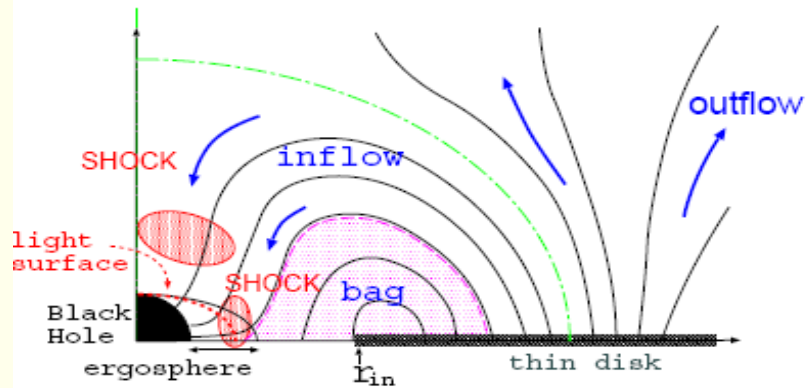
- Reaches approximately 100 AU into space ($=1.5 \times 10^{13}$ m)
- Voyager sonds are approaching/encountering the heliopause right now

Other other magnetospheres

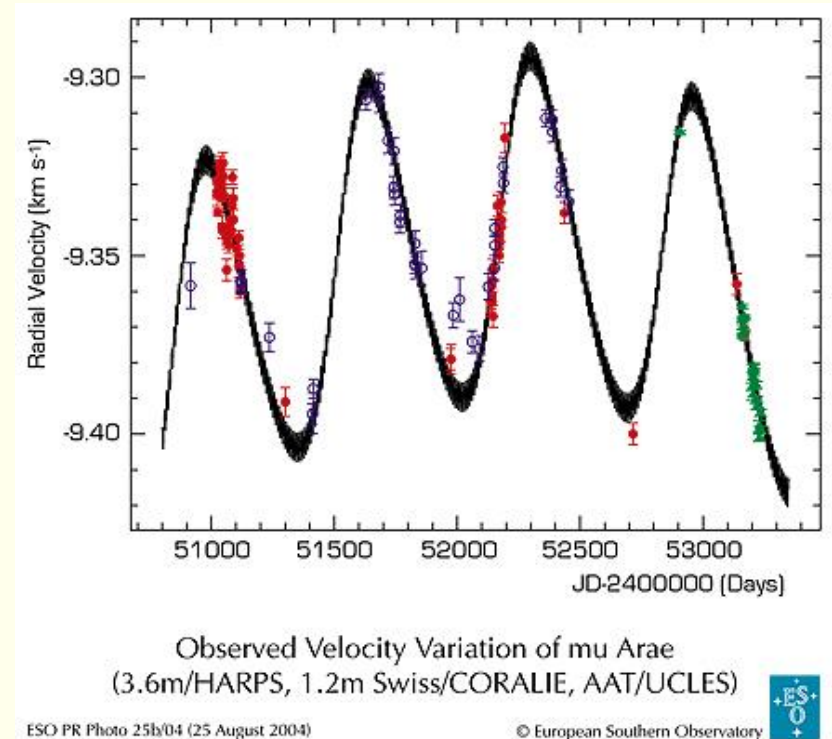
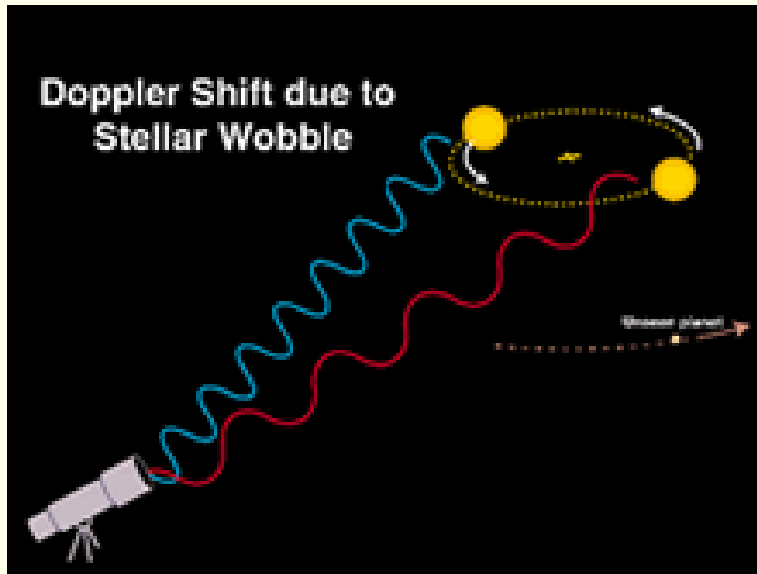
Exotic magnetospheres

“Black Hole Magnetosphere” =

$$\text{Loop (bag)} + \text{Disk-BH (inflow)} + \text{BH (Disk)} - r_{\infty} \text{ (outflow)}$$

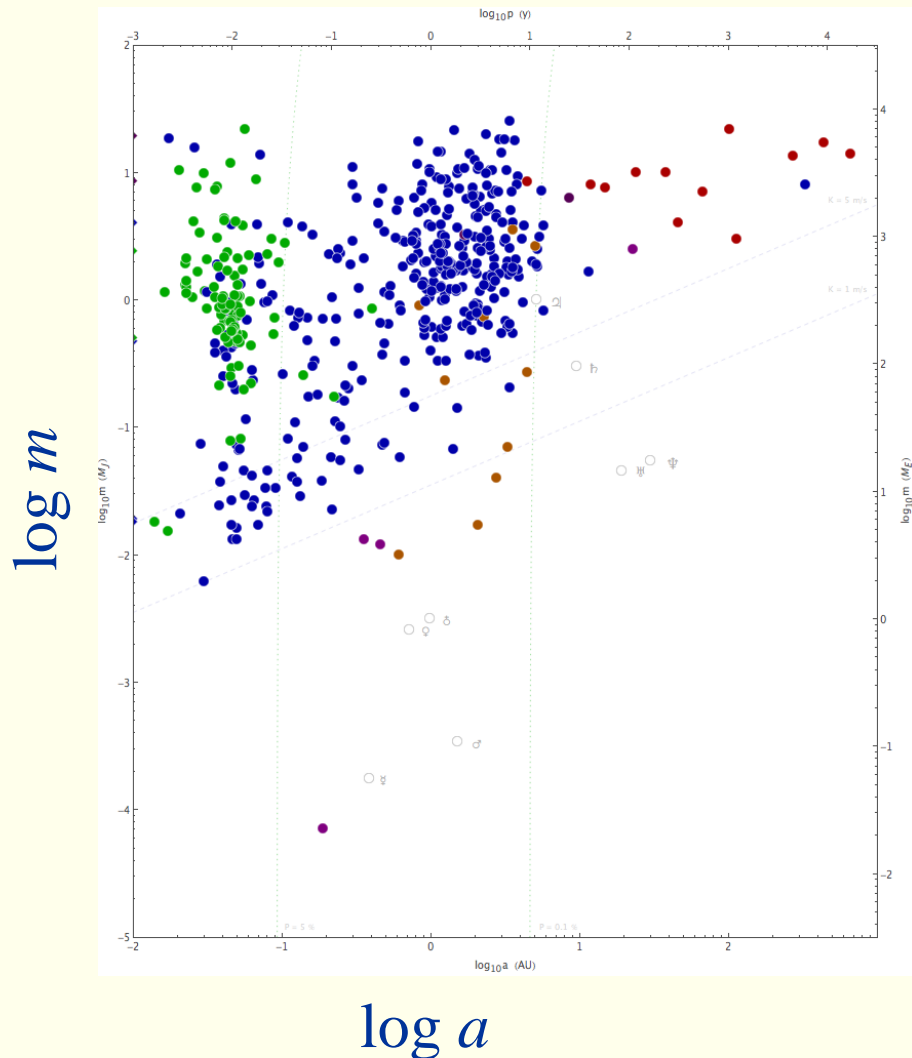


Exoplanets



Mostly detected by doppler shifts in starlight, or by dimming of the emitted starlight during transit of planet.

Magnetospheres of exoplanets



- 728 exoplanets found up to January 23, 2012
- Nothing is known about their atmospheres, ionospheres or magnetospheres
- Can possibly be detected by
 - radio emissions from auroral activity or from radiation belts
 - spectroscopy when the magnetosphere is in front of its sun