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Difference Approximations

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Summary of the Course

# Example: Partial Differential Equations

Michael Hanke

School of Engineering Sciences

Program construction in C++ for Scientific Computing



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### Outline

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- **2** Finite Difference Approximations
- 3 Implementation of Differential Operators
- 4 Boundary Conditions
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### What Do We Have

- Two simple classes for structured grids (Domain, Curvebase)
- A simple implementation of a matrix class (Matrix; don't use it for production codes!)

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### What Do We Want

- A class for representing grid functions
- Imposing boundary conditions
- A class for solving PDEs

Our running example will be the heat equation in 2D,

$$\frac{\partial}{\partial t}u = \frac{\partial^2}{\partial x^2}u + \frac{\partial^2}{\partial y^2}u.$$

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### The Domain Class

This is what we have so far:

```
class Domain {
  public:
    Domain(Curvebase&, Curvebase&, Curvebase&);
    void generate_grid(...);
    // more members
  private:
    Curvebase *sides[4];
    // more members
};
```

- We will need additional members for handling grids. Since grids do not allow any algebraic manipulation, using our Matrix class is not appropriate.
- We will use C-style arrays.
- It might be more convenient to use STL containers (e.g., vector).

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### The Domain Class: Enhanced

```
class Domain {
 public:
    Domain(Curvebase&, Curvebase&, Curvebase&,
           Curvebase&) : m(0), n(0), x(nullptr),
           y(nullptr) {}
    void generate_grid(int m_, int n_);
    int xsize() { return m: }
    int ysize() { return n; }
    Point operator()(int i, int j);
    bool grid_valid() { return m != 0; }
    // more members
 private:
    Curvebase *sides[4];
    int m, n;
    double *x, *y;
    // more members
};
```

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### One Dimensional Differences 1

• Consider a grid  $\Omega_h$ ,

$$a = x_0 < x_1 < \cdots < x_{n-1} < x_m = b.$$

• Let  $h_i = x_i - x_{i-1}$ . Then define, for a grid function  $u : \Omega_h \to \mathbb{R}$ ,

$$D_{-}u_{i} = \frac{u_{i} - u_{i-1}}{h_{i}}$$

$$D_{+}u_{i} = \frac{u_{i+1} - u_{i}}{h_{i+1}}$$

- If u is the restriction of a smooth function onto  $\Omega_h$ , these approximations are first order accurate.
- If the grid is equidistant,  $D_+D_-$  is a second order accurate approximation of  $u''(x_i)$  and

$$D_{+}D_{-}u_{i} = \frac{u_{i+1} - 2u_{i} + u_{i-1}}{h^{2}}$$

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### One Dimensional Differences 2

$$Du_i = \frac{u_{i+1} - u_{i-1}}{2h}$$

- First oder approximation to u' on a general grid
- Second order accuracy on a constant stepsize grid

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### **Boundaries**

- The operators introduced above are not applicable at boundaries.
- Possibility 1: One-sided differences

$$Du_0 = \frac{3u_0 - 4u_1 + u_2}{3h}$$

$$Du_m = \frac{u_{m-2} - 4u_{m-1} + 3u_m}{3h}$$

• Possibility 2: Use ghost points

$$Du_0 = \frac{u_1 - u_{-1}}{2h}$$

$$Du_m = \frac{u_{m+1} - u_{m-1}}{2h}$$

How to get values for the ghost points?

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### Nonuniform Grids

- Order of approximation is determined using Taylor expansions.
- Ansatz:

$$u'(x_i) \approx a_- u(x_{i-1}) + a_0 u(x_i) + a_+ u(x_{i+1}) =: D_0 u(x_i)$$

Taylor expansion:

$$u(x_{i-1}) = u(x_i) - h_i u'(x_i) + \frac{1}{2} h_i^2 u''(x_i) + O(h^3)$$
  
$$u(x_{i+1}) = u(x_i) + h_{i+1} u'(x_i) + \frac{1}{2} h_{i+1}^2 u''(x_i) + O(h^3)$$

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# Nonuniform Grids (cont)

• Inserting into the expression for  $D_0u$ , we obtain after coefficient comparison

$$a_{-} = \frac{-h_{i+1}}{h_{i}(h_{i} + h_{i+1})}$$

$$a_{0} = \frac{h_{i+1} - h_{i}}{h_{i}h_{i+1}}$$

$$a_{+} = \frac{h_{i}}{h_{i+1}(h_{i} + h_{i+1})}$$

and

$$D_0 u(x_i) - u'(x_i) = \frac{1}{6} h_i h_{i+1} u'''(x_i) + \dots$$

- For an equidistant grid, the coefficients reduce to  $a_- = -1/2h$ ,  $a_0 = 0$ ,  $a_+ = 1/2h$ .
- One sided expressions??

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### An Alternative Idea

• Assume that the grid is created using a mapping  $\phi:[0,1]\to[a,b]$  with  $x_i=\phi(s_i),\ i=0,\ldots,m$  with a uniform grid

$$s_i = i\sigma, \quad \sigma = m^{-1}.$$

• Then,  $du/ds = du/dx \cdot dx/ds$ , and

$$u_{x}(x_{i}) \approx \frac{1}{dx(s_{i})/ds} \frac{u_{i+1} - u_{i-1}}{2\sigma}$$

is a second order approximation.

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### And Another Idea

• If the derivative dx/ds is not known, it can be approximated with second order accuracy by

$$\frac{dx}{ds}(s_i) \approx \frac{x_{i+1} - x_{i-1}}{2\sigma}$$

such that

$$u_{x}(x_{i}) \approx \frac{u_{i+1} - u_{i-1}}{x_{i+1} - x_{i-1}}$$

is second order accurate!

- Needed:  $\phi$  is a smooth mapping!
- Note: We need only two grid points in order to obtain the same order of accuracy as in the approximation in physical domain.

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# Approximation of u''

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Going either way, we have an approximation

$$u'(x_i) \approx D_0 u_i$$
.

A second order approximation to the second derivative can be defined by

$$u''(x_i) \approx D_2 u_i = D_0 D_0 u_i.$$

This approximation evaluates to a five-point stencil!

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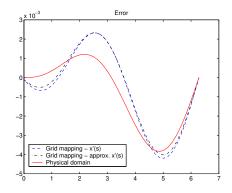
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### Example: Comparison of Accuracy



$$egin{aligned} u(x) &= \sin x \ x(s) &= 2\pi rac{1 + anh(\delta(s-1)/2)}{ anh(\delta/2)}, \quad \delta = 5 \end{aligned}$$

Hyperbolic tangent stretching, 100 gridpoints.

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### Conclusions

- All approximations are 2nd order accurate.
- In this simple example, approximation in physical domain is more accurate.
- The stencil (number of grid points used) is larger in physical domain for obtaining the same order of accuracy.

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## 2D: Physical Domain

Ansatz:

$$u_x(x_{i,j},y_{i,j}) \approx \sum_{k,l} a_{kl} u_{i+k,j+l}$$

Taylor expansion around  $(x_{i,j}, y_{i,j})$ :

$$\begin{split} & \sum_{k,l} a_{k,l} u_{i+k,j+l} \\ & = \sum_{k,l} a_{k,l} \sum_{\nu=0} \frac{1}{\nu!} \left( (x_{i+k,j+l} - x_{i,j}) \frac{\partial}{\partial x} + (y_{i+k,j+l} - y_{i,j}) \frac{\partial}{\partial y} \right)^{\nu} u \\ & = \sum_{\nu=0} \sum_{p=0}^{\nu} \left[ \sum_{k,l} a_{k,l} \frac{1}{\nu!} {\nu \choose p} (x_{i+k,j+l} - x_{i,j})^p (y_{i+k,j+l} - y_{i,j})^{\nu-p} \right] \frac{\partial^p}{\partial x^p} \frac{\partial^{\nu-p}}{\partial y^{\nu-p}} u \end{split}$$

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# $D_{0,x}$ in Physical Domain

First order:

$$\sum_{k,l} a_{k,l} = 0$$

$$\sum_{k,l} a_{k,l} (x_{i+k,j+l} - x_{i,j}) = 1$$

$$\sum_{k,l}^{k,l} a_{k,l} (y_{i+k,j+l} - y_{i,j}) = 0$$

Second order additionally:

$$\sum_{k,l} a_{k,l} (x_{i+k,j+l} - x_{i,j})^2 = 0$$

$$\sum_{k,l} a_{k,l} (x_{i+k,j+l} - x_{i,j}) (y_{i+k,j+l} - y_{i,j}) = 0$$

$$\sum_{k,l} a_{k,l} (y_{i+k,j+l} - y_{i,j})^2 = 0$$

So we expect 6 gridpoints necessary for second order accuracy!

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### Stencil in Reference Coordinates

### Remember:

- Let  $\Phi$  to a (smooth) one-to-one mapping  $\Phi: [0,1]^2 \to \Omega$ .
- For given m, n, a uniform grid on  $[0, 1]^2$  can be defined by:

$$\xi_i = ih_1, \quad h_1 = 1/m, \quad i = 0, \dots, m,$$
  
 $\eta_j = jh_2, \quad h_2 = 1/n, \quad j = 0, \dots, n.$ 

• A strucured grid on  $\Omega$  can then simply be obtained via

$$x_{ij} = \Phi_x(\xi_i, \eta_j), \quad y_{ij} = \Phi_y(\xi_i, \eta_j), \quad i = 0, \dots, m, j = 0, \dots, n.$$

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# Reference Coordinates (cont)

Using the chain rule of differentiation, we obtain

$$\frac{\partial u(x,y)}{\partial \xi} = \frac{\partial u}{\partial x} \cdot \frac{\partial \Phi_x}{\partial \xi} + \frac{\partial u}{\partial y} \cdot \frac{\partial \Phi_y}{\partial \xi}$$
$$\frac{\partial u(x,y)}{\partial \eta} = \frac{\partial u}{\partial x} \cdot \frac{\partial \Phi_x}{\partial \eta} + \frac{\partial u}{\partial y} \cdot \frac{\partial \Phi_y}{\partial \eta}$$

Since the transformation  $\Phi_x, \Phi_y$  is known, this is a linear system for the partial derivatives  $\partial u/\partial x$ ,  $\partial u/\partial y$ .

Let

$$J = \begin{pmatrix} \frac{\partial \Phi_{\mathbf{x}}}{\partial \xi} & \frac{\partial \Phi_{\mathbf{y}}}{\partial \xi} \\ \frac{\partial \Phi_{\mathbf{x}}}{\partial \eta} & \frac{\partial \Phi_{\mathbf{y}}}{\partial \eta} \end{pmatrix}$$

Then

$$\frac{\partial u}{\partial x} = \frac{1}{\det J} \left( \frac{\partial u}{\partial \xi} \cdot \frac{\partial \Phi_x}{\partial \xi} - \frac{\partial u}{\partial \eta} \cdot \frac{\partial \Phi_y}{\partial \xi} \right)$$
$$\frac{\partial u}{\partial x} = \frac{1}{\det J} \left( \frac{\partial u}{\partial \eta} \cdot \frac{\partial \Phi_y}{\partial \eta} - \frac{\partial u}{\partial \xi} \cdot \frac{\partial \Phi_x}{\partial \eta} \right)$$

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### Referens Coordinates (cont)

- The derivatives with respect to reference coordinates can be approximated by standard stencils (4-point stencil).
- Once all partial derivatives w r t  $\xi$  have been evaluated, the necessary partial derivatives w r t x, y can be computed.

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# Class for Grid Functions: Requirements

- (Scalar) grid functions are defined on grids.
- We are using structured grids as represented in the class Domain.
- Operations allowed with grid functions:
  - Addition, multiplication by a scalar (they form a vector space)
  - Pointwise multiplication (together, this becomes a commutative algebra)
  - Differentiation (e.g., by finite differences)
  - Computation of norms
  - Integration (? maybe)

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### Further Considerations

- In the two-dimensional case, many of these operations are already implemented in the Matrix class!
- However, some operations are not meaningful for grid functions, e.g., matrix-matrix multiplication.
- A grid functions lives only on a specific grid:
  - Shall the grid be part of an object?
  - Many grid functions share the same grid!
  - Algebraic manipulations are only defined for grid functions living on the same grid

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### Remember: The Matrix Class

```
class Matrix {
  int m, n; // should be size_t
  double *A:
public:
  Matrix(int m_{-} = 0, int n_{-} = 0) : m(m_{-}), n(n_{-}),
             A(nullptr) {
    if (m*n > 0) {
      A = \text{new double}[m*n];
       std::fill(A,A+M*n,0.0);
// etc
```

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### Implementation of Grid Functions

```
class GFkt {
 private:
    Matrix u:
    Domain *grid;
 public:
    GFkt(Domain *grid_) : u(grid_->xsize()+1,
                   grid_->ysize()+1), grid(grid_) {}
    GFkt(const GFkt& U) : u(U.u), grid(U.grid) {}
    GFkt& opearator=(const GFkt& U);
    GFkt operator+(const GFkt& U) const;
    GFkt operator*(const GFkt& U) const;
// etc
};
```

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### A Sample Implementation

```
GFkt GFkt::operator+(const GFkt& U) const {
  if (grid == U.grid) { // defined on the same grid?
     GFkt tmp(grid);
     tmp.u = u+U.u; // Matrix operator+()
     return tmp;
  else error();
GFkt GFkt::operator*(const GFkt& U) const {
  if (grid == U.grid) { // defined on the same grid?
    GFkt tmp(grid);
    for (int j = 0; j <= grid.ysize(); j++)</pre>
      for (int i = 0; i <= grid.xsize(); i++)
        tmp.u(i,j) = u(i,j)*U.u(i,j);
    return tmp;
  else error():
```

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### A Problem And Its Solution

- The grid is handled by the caller.
- In the above implementation, the caller may delete the grid such that all objects referring to it have a dangling pointer!
- In C++ 11 there is a solution: smart pointers
- Smart pointers belong to the C++ library, include file: memory

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### **Smart Pointers**

- There are two types of them: shared\_ptr and unique\_ptr.
- Both classes are in fact template classes: The template argument is a typename.
- shared\_ptr uses a reference count: As soon as the reference count reaches 0, the dynamic object will be destroyed. But not earlier!
- This way, all resources will be freed (including dynamic memory).
- C-type pointers and smart pointers cannot be mixed! There is always an explicit type cast necessary! Recommendation: Avoid mixing.

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# Smart Pointers (cont)

 Create a smart pointer, initialize it to 0 (nullptr): shared\_ptr<class> p1;

• The equivalent of new:

```
shared_ptr<class> p2 = make_shared<class>(args);
The following statement is in error:
```

The following statement is in error

```
shared_ptr<class> p3 = new class(args); // Error!
But this works:
    shared_ptr<class> p3 =
```

shared\_ptr<class>(new class(args));

• There is no equivalent of delete needed.

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### A Better Implementation of GFkt

### Notes:

 We assume silently that, once a grid has been generated, it will never be changed!

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 It is most probably a good idea to use shared pointers in Domain, too:
 shared\_ptr<Curvebase> sides[4];

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# Implementation of $D_{0,x}$

```
GFkt GFkt::D0x() const {
    GFkt tmp(grid);
    if (grid->grid_valid()) {
        // generate derivative in tmp
        // according to one of the possibilities above
    }
    return tmp;
}
```

- The function DOy can be implemented similarly.
- In order to reduce overhead, it might be a good idea to implement even

```
void GFkt::D0xy(GFkt *dx, GFkt *dy) const;
```

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## **Boundary Conditions**

Name	Prescribed	Interpretation
Dirichlet	и	Fixed temperature
Neumann	$\partial u/\partial n$	Energy flow
Robin (mixed)	$\partial u/\partial n + f(u)$	Temperature dependent flow
Periodic	, , ,	

Boundary conditions have a crucial impact on the solution.

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## What are Boundary Conditions?

1 The mathematician's point of view:

domain + differential equation + boundary conditions

2 The physicist's point of view:

 $\begin{array}{ccc} \text{differential equation} & \longrightarrow & \text{physics} \\ \text{domain} & \longrightarrow & \text{space} \end{array}$ 

boundary conditions  $\longrightarrow$  influence of outer world

**3** The software engineer's point of view:

 $\begin{array}{ccc} \text{differential equation} & \longrightarrow & \text{expression of} \\ & & & \text{differentials} \end{array}$ 

boundary conditions  $\longrightarrow$  what??

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### Object-Oriented Representation

- As part of the PDE
  - mathematical interpretation
  - requires high-level representation of equation and discretization
  - · difficult to obtain efficiency
- As part of the grid function
  - mathematically correct
  - no class for PDEs needed
  - convenient for exlicit time-stepping
- As part of the operator (e.g., D<sub>0</sub>)
  - · convenient for implicit and explicit methods
  - can be difficult to implement
  - may encounter mathematical contradictions if used wronly

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## A First Attempt

Associate boundary conditions with grid functions:

```
class Solution {
  public:
    Solution(Domain *D) : sol(D) {}
    ~Solution();
    void timesteps(double dt, int nsteps);
    void init(); // Set initial condition
    void print();
  private:
    GFkt sol;
    void impose_bc();
};
```

impose\_bc() will be called in timesteps() for imposing the boundary conditions.

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### Discussion

- The proposed implementation is questionable because the boundary conditions and timestepping are "hardwired".
- It is better to have a *class* for boundary conditions:

```
class BCtype {
  public:
    BCtype(GFkt& u, int boundary_id);
    virtual void impose(GFkt& u) = 0;
};
```

- The actual definition of the boundary condition takes place in derived classes.
- This way, several boundaries can share the same condition (e.g., homogeneous Dirichlet conditions).
- Classes can be derived for Dirichlet, Neumann, Robin boundary conditions.

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### **Example Implementation**

### Assumptions:

- The grid has four distinct edges (as ours in the previous Domain class).
- Each edge is associated with one boundary condition, only.

### Then:

```
class Solution {
  public:
    Solution(Domain *D) : sol(D) {}
    ~Solution();
    void print();
  private:
    GFkt sol;
    shared_ptr<BCtype> bcs[4];
    virtual void init() = 0;
    virtual void bc() = 0;
};
```

We have separated: the grid, the equation, the initial conditions, and the boundary conditions.

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# Time Stepping

For the heat equation in 2D, we can implement the explicit Euler method now:

```
Solution u(&d);
u.init();
for (int step=0; step < maxsteps; step++) {
  u += dt*(u.D2x()+u.D2y());
  t += dt;
  u.bc();
}</pre>
```

(Provided the missing functions are implemented along the lines provided before)

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### Summary

- Finite difference approximations on structured grids.
- Smart pointers
- Implementation strategies for differential operators, boundary conditions, and time steppers.

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## Course Summary

### C++

- Basic elements of C++
- Abstract data types, C++ classes
- Constructors, destructors, memory management, copy, move
- Operator overloading
- Inheritance, abstract classes
- Templates, STL
- I/O

### **Scientific Computing**

- Structured grids, differential operators, boundary conditions
- Implementation strategies and their C++ tools
- Efficient programming
- Scientific libraries

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