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## Optical Physics: Summary Session 4

## Geometrical optics: Principal planes

Thick Lenses: All lenses can not be approximated as thin, i.e. their thickness, $d$, has to be taken into account. A thick lens optical properties in air are described by two principal planes, $H_{1}$ and $H_{2}$, and an effective focal length, $f_{e f f}$ :
$\frac{1}{f_{\text {eff }}}=\left(n_{l}-1\right)\left[\frac{1}{R_{1}}-\frac{1}{R_{2}}+\frac{\left(n_{l}-1\right) d}{n_{l} R_{1} R_{2}}\right]$
Vertex: point where the lens surface intersect the optical axis
The distance from the left vertex to $H_{1}=-\frac{f_{\text {eff }}\left(n_{l}-1\right) d}{R_{2} n_{l}}$ (positive direction from OS towards IS)

The distance from the right vertex to $H_{2}=-\frac{f_{\text {eff }}\left(n_{l}-1\right) d}{R_{1} n_{l}}$ (positive direction from OS towards IS)

If one "disregards" the distance between $H_{1}$ and $H_{2}$, one can perform calculations and ray constructions according to the same formulas as for thin lenses. One thus pretends that $H_{1}$ is the left vertex of a thin lens and $H_{2}$ the right vertex. $f_{\text {eff }}$ is the focal length of the thin lens, measured as positive from $H_{1}$ towards OS and from $H_{2}$ towards IS, exactly as for a thin lens with positive focal length.

System of thin lenses: A system of thin lenses can be described by principal planes in the same way as a thick lens. It has a front principal plane, $H_{1}$ and a back principal plane $H_{2}$, and an effective focal length, $f_{\text {eff }}$ :
$\frac{1}{f_{e f f}}=\frac{1}{f_{1}}+\frac{1}{f_{2}}-\frac{d}{f_{1} f_{2}}$
The distance from the first lens to $H_{1}: \overline{H_{11} H_{1}}=\frac{f_{\text {eff }} d}{f_{2}}$ (positive direction from OS towards IS)

The distance from the second lens to $H_{2}: \overline{H_{22} H_{2}}=-\frac{f_{\text {eff }} d}{f_{1}}$ (positive direction from OS towards IS)
Observe sign! Error in equation 6.10

## Superposition of waves

Sum of two harmonic waves:
$\overrightarrow{E_{t o t}}=\overrightarrow{E_{1}}+\overrightarrow{E_{2}}=E_{01} e^{i \varphi_{1}}+E_{02} e^{i \varphi_{2}}$ where $\varphi=\vec{k} \cdot \vec{r}-\omega t+\epsilon$
The intensity of the total wave is then:
$I_{t o t}=\left|\overrightarrow{E_{t o t}}\right|^{2}=\overrightarrow{E_{t o t}}{\overrightarrow{E_{t o t}}}^{*}=E_{01}^{2}+E_{02}^{2}+E_{01} E_{02}\left[e^{i\left(\varphi_{1}-\varphi_{2}\right)}+e^{-i\left(\varphi_{1}-\varphi_{2}\right)}\right]$
This can be written as:
$I_{t o t}=\left|\overrightarrow{E_{t o t}}\right|^{2}=\overrightarrow{E_{t o t}}{\overrightarrow{E_{t o t}}}^{*}=E_{01}^{2}+E_{02}^{2}+2 E_{01} E_{02} \cos \left(\varphi_{1}-\varphi_{2}\right)$

If the two harmonic waves have the same frequency, i.e. $\omega_{1}=\omega_{2}=\omega$ (and thus also $k_{1}=k_{2}=k$ ), the total wave is also harmonic and can be written as $\overrightarrow{E_{t o t}}=E_{0} e^{i(\vec{k} \cdot \vec{r}-\omega t+\epsilon)}$. The new phase factor $\epsilon$ is given by:
tan $\epsilon=\frac{E_{01} \sin \epsilon_{1}+E_{02} \sin \epsilon_{2}}{E_{01} \cos \epsilon_{1}+E_{02} \cos \epsilon_{2}}$.
The total amplitude is:
$E_{0}^{2}=E_{01}^{2}+E_{02}^{2}+2 E_{01} E_{02} \cos (\epsilon)$
The last term in the expression of the amplitude is called the interference term and is the one giving rise to destructive and constructive interference.

## Sum of multiple harmonic waves:

When two harmonic waves with different amplitudes and phases are added, it results in a non harmonic wave, but the wave will still be periodic. This is exactly what Fourier's theorem states: $A$ periodic function can always be viewed as a superposition of multiple harmonic waves, i.e. it can always be divided into its frequency components.
Addition of periodic waves in general can therefore be performed by adding the frequency components. The frequency components are found by Fourier transformation of the wave. The Fourier transform of $E(t)$ is given by:
$\widehat{E}(\omega)=\int_{-\infty}^{\infty} E(t) e^{-i \omega t} d t$
Observe that convolution of two functions in time space is a simple multiplication in frequency space and vice versa:
$f * g(t)=\int_{-\infty}^{\infty} f(t-\tau) g(\tau) d \tau=\widehat{f}(\omega) \cdot \widehat{g}(\omega)$
$f(t) \cdot g(t)=\widehat{f} * \widehat{g}(\omega)$

## Laser cavity:

Frequency difference between the longitudinal modes in a laser cavity: $\Delta \nu=\frac{v}{2 L}$ where $v$ is the speed of light in the cavity and $L$ is the cavity length.

## Heisenberg's uncertainty principle:

$\Delta t \Delta E>=\frac{\hbar}{2}$
where $\Delta t$ is the spread in time and $\Delta E$ the spread in energy $(E=h \nu)$.

